

Heavy Ion Collisions

A. Marin (GSI)

Spanish High Energy Physics School
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Hydrodynamics and flow

Hydrodynamics

Framework to describe the space-time evolution of matter under local thermal equilibrium

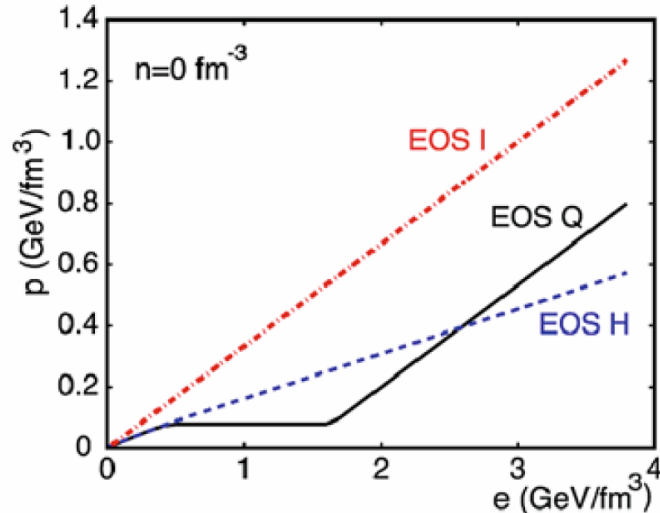
Inputs to hydrodynamical models:

- Equation of state (Ideal hydro):

$P=P(e,n)$ P : pressure, e : energy density, n : baryon density

viscous hydro: needs shear viscosity, bulk viscosity, heat conductivity

P.F. Kolb and U.W. Heinz, nucl-th/0305084



EOS I: Ideal gas of relativistic massless particles

EOS Q: Includes hadron masses and phase transition hadronic matter-QGP

EOS H: Hadron resonance gas model

- Energy-momentum conservation and net-baryon current conservation:

$$\partial_{\mu} T^{\mu\nu} = 0, \quad \partial_{\mu} N_i^{\mu} = 0,$$

- Initial conditions from Glauber model or Color Glass Condensate model

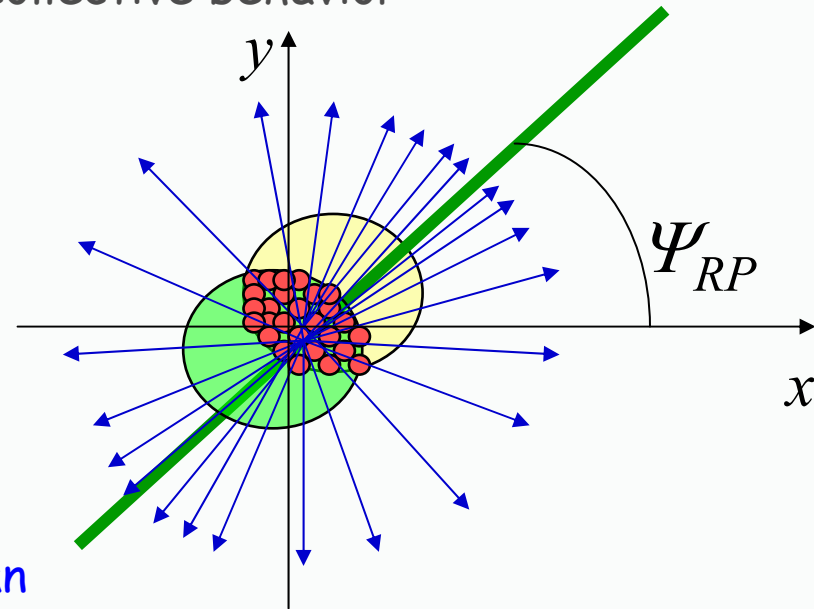
Reaction plane and transverse flow

- Anisotropic transverse flow is a **correlation** between the azimuth [$=\tan^{-1}(p_y/p_x)$] of the produced particles and the reaction plane!

⇒ Unambiguous signature of collective behavior

- The **reaction plane** is the plane defined by the impact parameter and the beam direction

- The orientation of the reaction plane can be estimated from the global event anisotropy

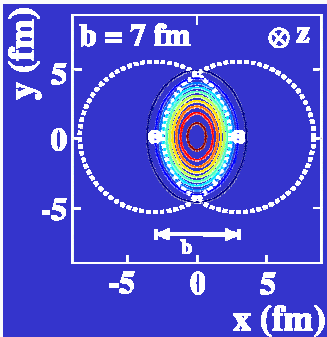
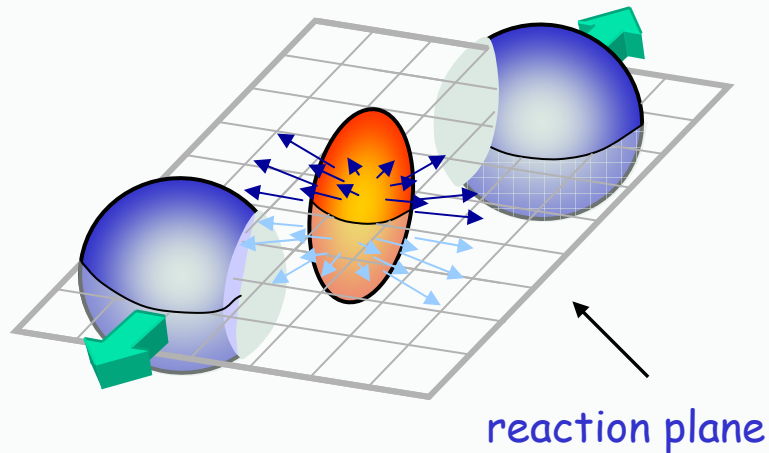


Reaction plane

- The initial spatial anisotropy is transformed into momentum anisotropy.
 - Created and saturates in first femtometers. The observed v_2 is expected to be sensitive to the initial stage of the collision
- Analyze in terms of response of the system to initial pressure
- Deduce (model) thermalizing time, EOS

$$E \frac{d^3 N}{d^3 p} = \frac{1}{2\pi} \frac{d^2 N}{p_t dp_t dy} \left(1 + 2 \sum_{n=1}^{\infty} v_n(p_t, y) \cos(n(\phi - \Psi_r)) \right)$$

$$v_2 = \langle \cos 2(\phi - \Psi_r) \rangle, \quad \phi = \tan^{-1} \left(\frac{p_y}{p_x} \right)$$



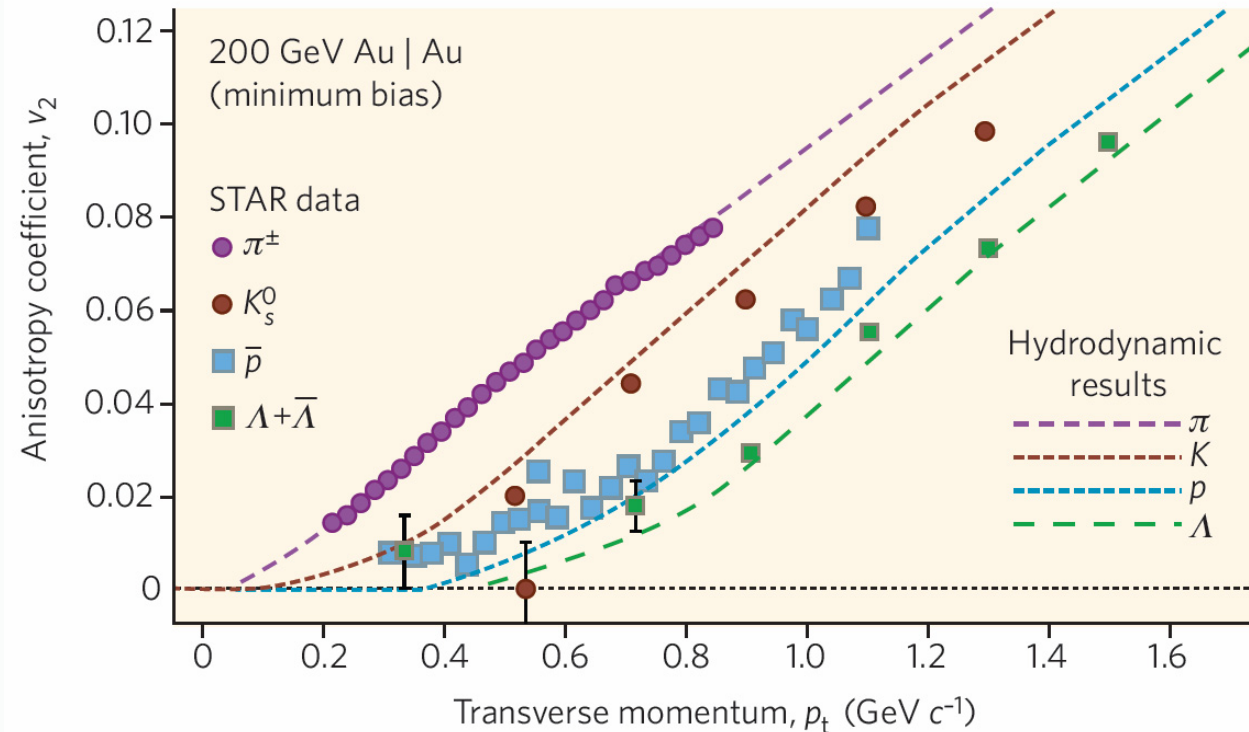
$$\varepsilon \equiv \frac{\langle y^2 - x^2 \rangle}{\langle y^2 + x^2 \rangle}$$

The sine terms in the Fourier expansion vanish because of the reflection symmetry with respect to the reaction plane.

Fourier coefficients: $v_n(p_T; y) = \langle \cos[n(\phi - \Psi_{RP})] \rangle$
 v_1 : Strength of the directed flow (small at midrapidity)
 v_2 : Strength of the elliptic flow

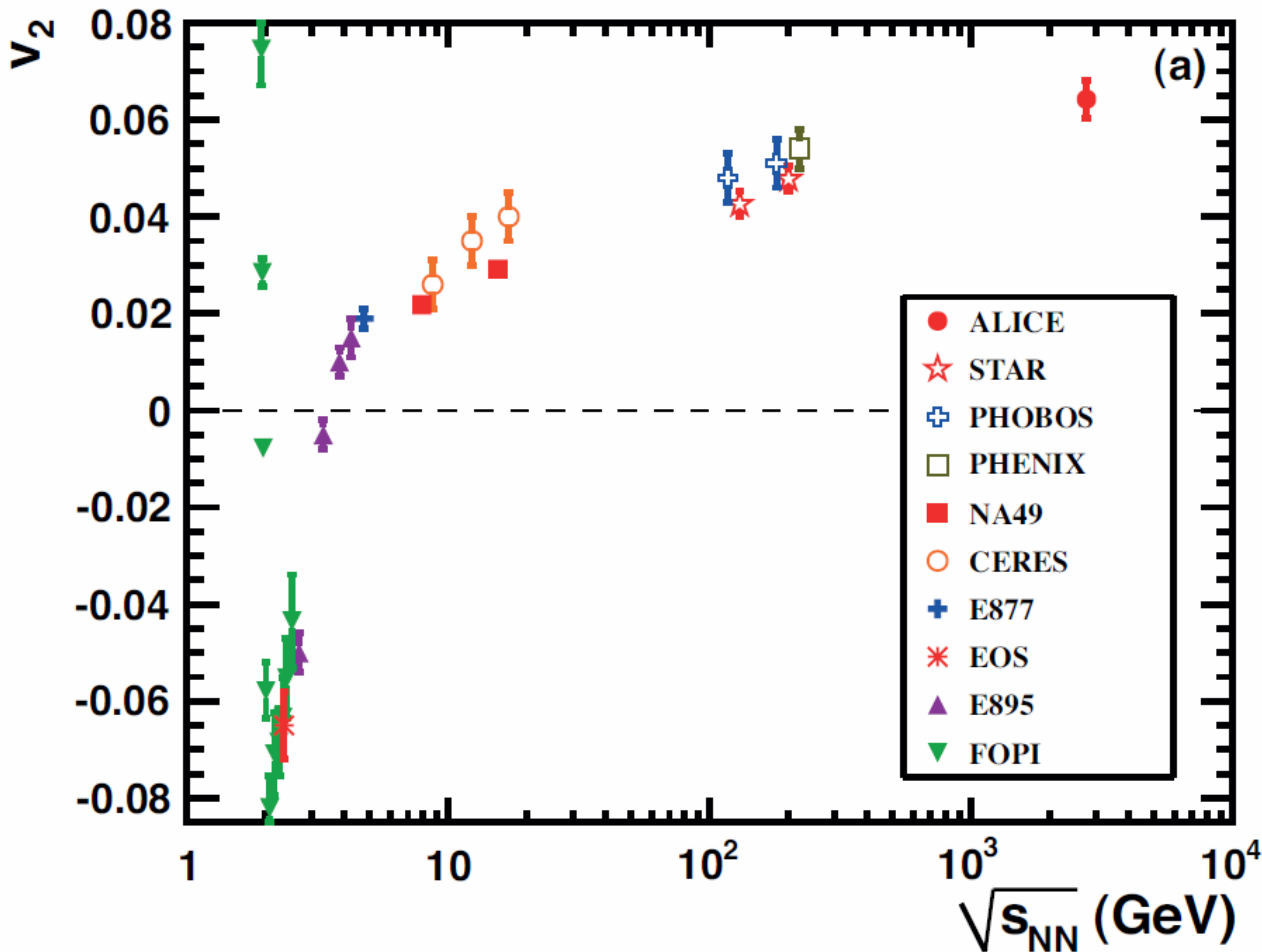
Elliptic flow at RHIC

Phys. Rev. C72, 014904 (2005), Nature 448 (2007) 302



Mass ordering observed. Agrees with hydro for $p_t < 1.5 \text{ GeV}/c$.
It was predicted by ideal hydro.
Perfect liquid created at RHIC.

v_2 excitation function



• $\sqrt{s_{NN}} < 2 \text{ GeV}$:
In-plane, rotation-like emission

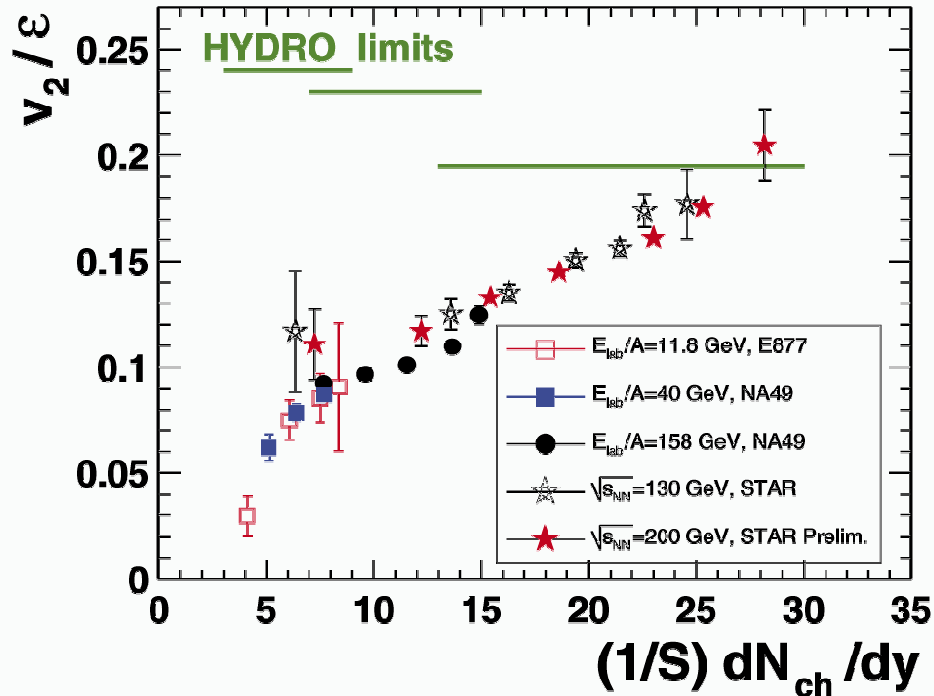
• $2 < \sqrt{s_{NN}} < 4 \text{ GeV}$:
Onset of expansion and presence of spectator matter inhibits in plane particle emission ("squeezeout")

• $\sqrt{s_{NN}} > \sim 4 \text{ GeV}$:
Initial eccentricity leads to pressure gradients that cause positive v_2

New J.Phys.13:055008,2011, R. Snellings, arXiv:1102.3010,
P. Sorensen, arXiv:0905.0174,
Andronic, Phys. Lett. B612 (2005) 173-180 , nucl-ex/0411024

Energy dependence of flow

eccentricity vs. particle multiplicity in overlap region



hydrodynamical limit
reached at RHIC
→ 'ideal fluid'

- clear predictions from hydrodynamics
- sensitive to equation-of-state

v_2 quark number scaling

PHENIX, PRL, 98, 162301

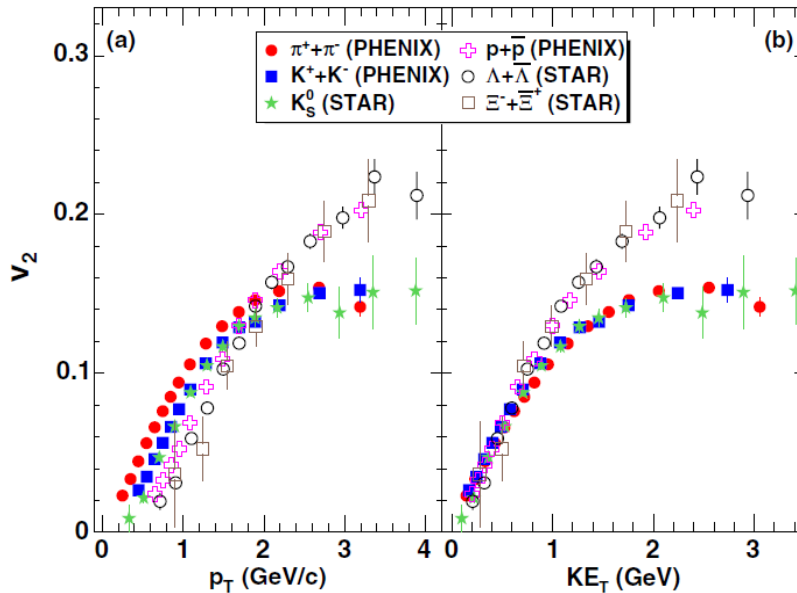


FIG. 2 (color online). (a) v_2 vs p_T and (b) v_2 vs KE_T for identified particle species obtained in minimum-bias Au + Au collisions. The STAR data are from Refs. [24,43].

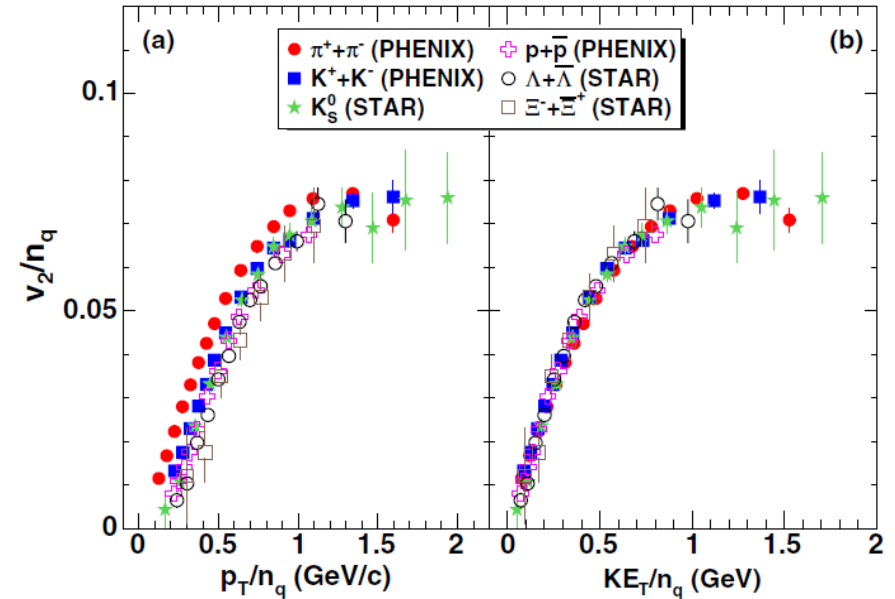
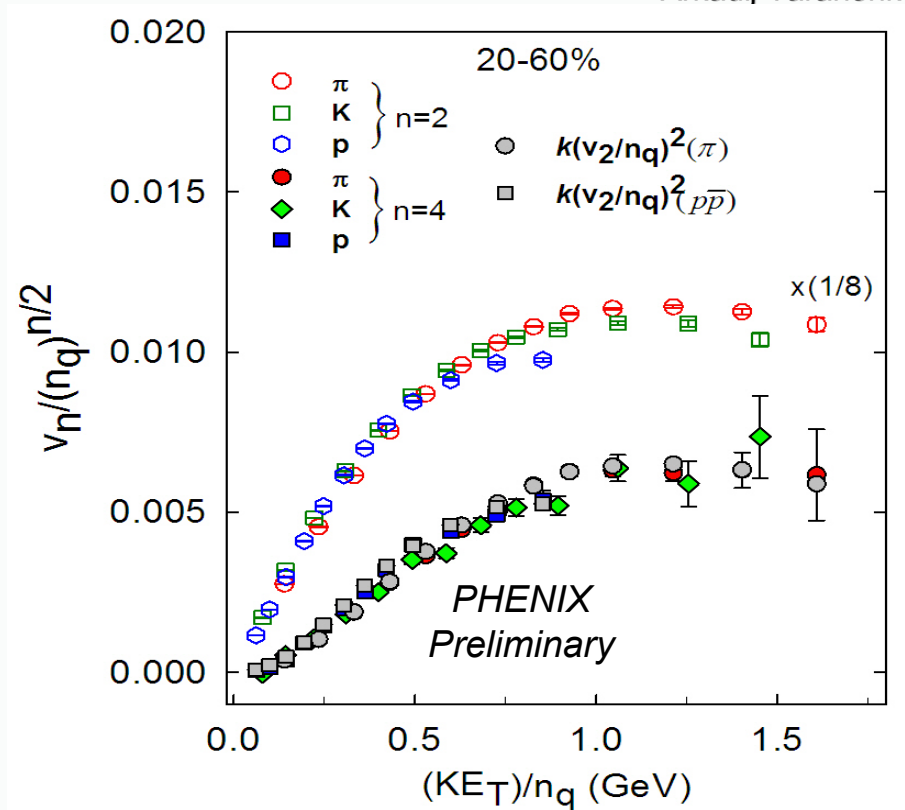


FIG. 3 (color online). (a) v_2/n_q vs p_T/n_q and (b) v_2/n_q vs KE_T/n_q for identified particle species obtained in minimum-bias Au + Au collisions. The STAR data are from Refs. [24,43].

- Excellent scaling over the full range of KE_T/n_q values. ($KE_T = m_T - m$)
- An indication of the inherent quark like degrees of freedom in the flowing matter.

Flow universal?

Arkadij Taranenko, QM2009



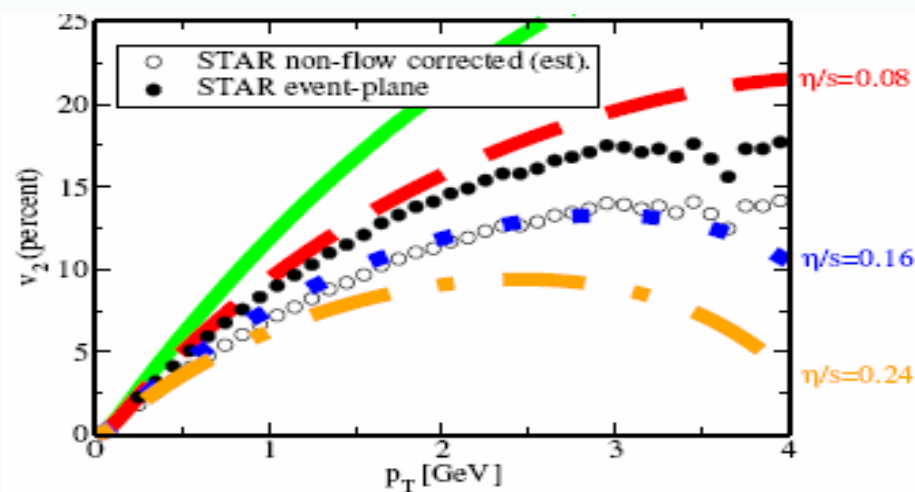
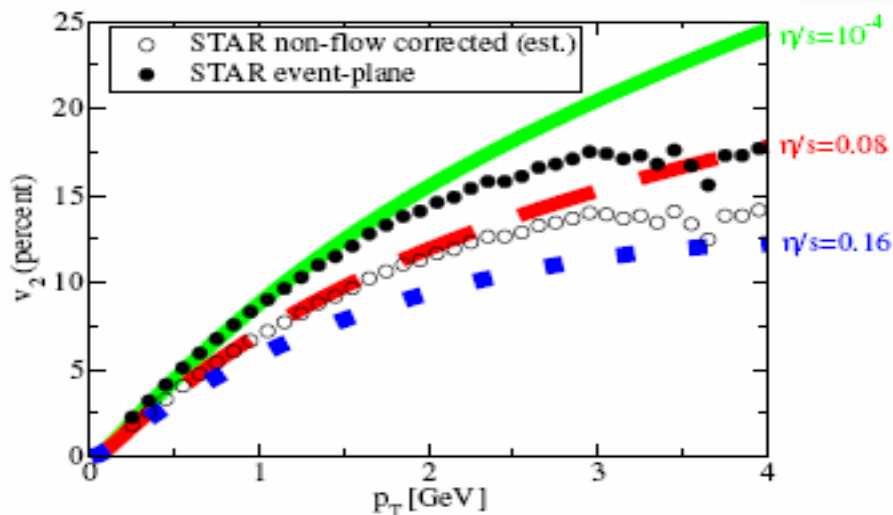
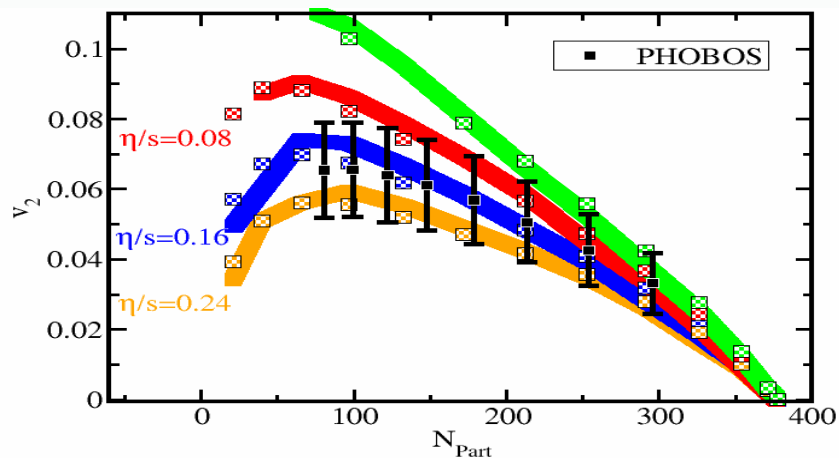
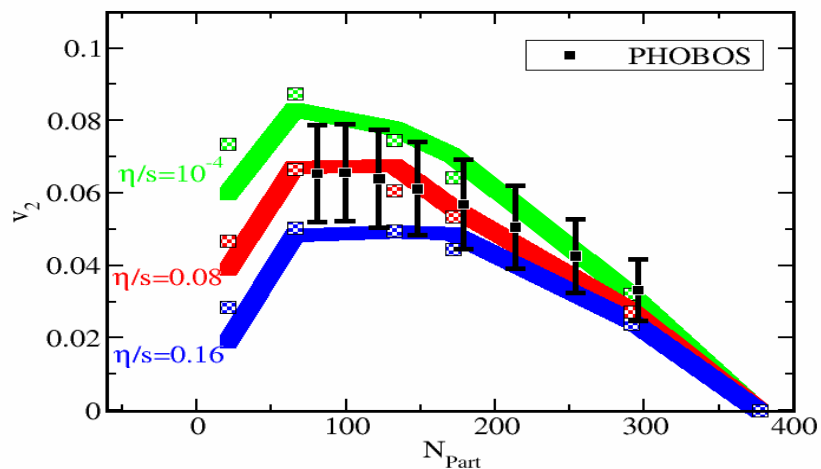
Baryon and meson V_2 & V_4 scale to a universal curve
as a function of $(KE_T)/n_q$

Sensitivity of v_2 to viscosity

Glauber

Luzum & Romatschke, PRC 2008

CGC

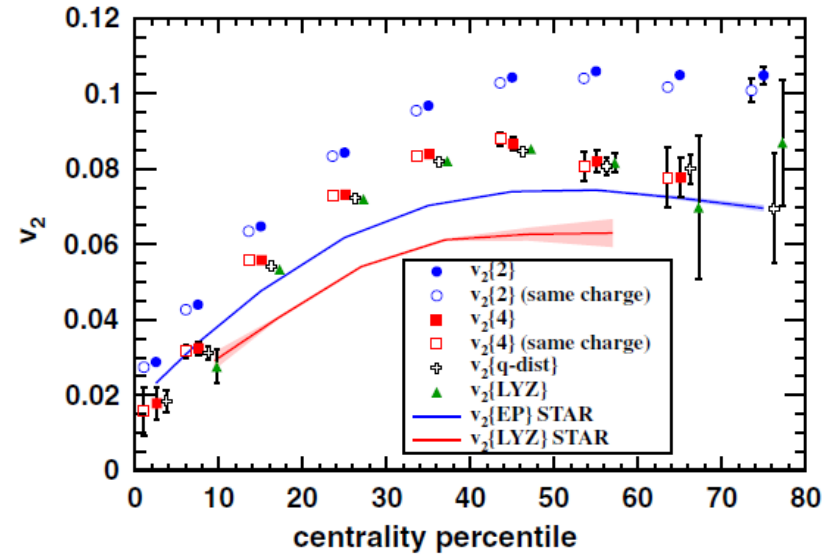
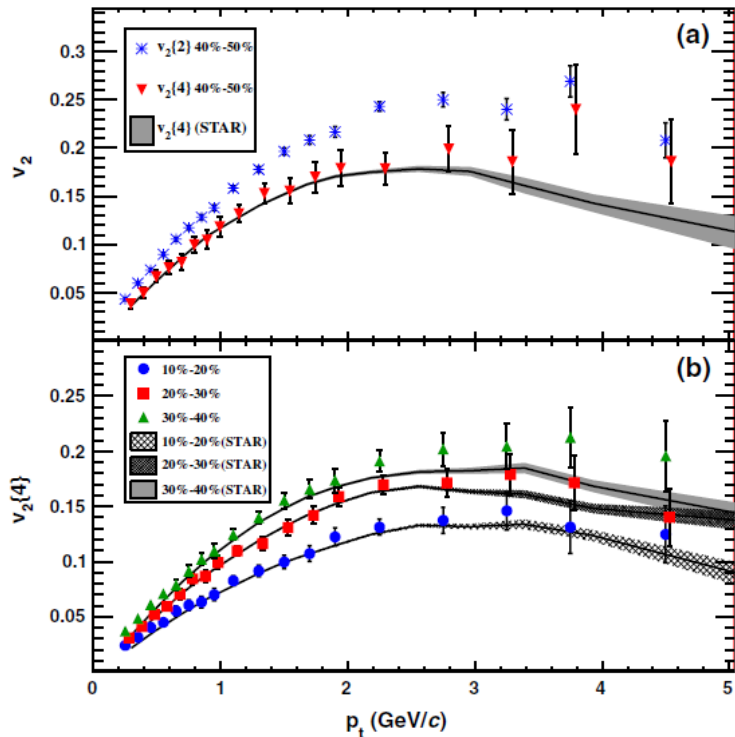


Viscosity reduces v_2

Present conservative upper limit:

$$\eta/s \leq 5 \times (1/4\pi)$$

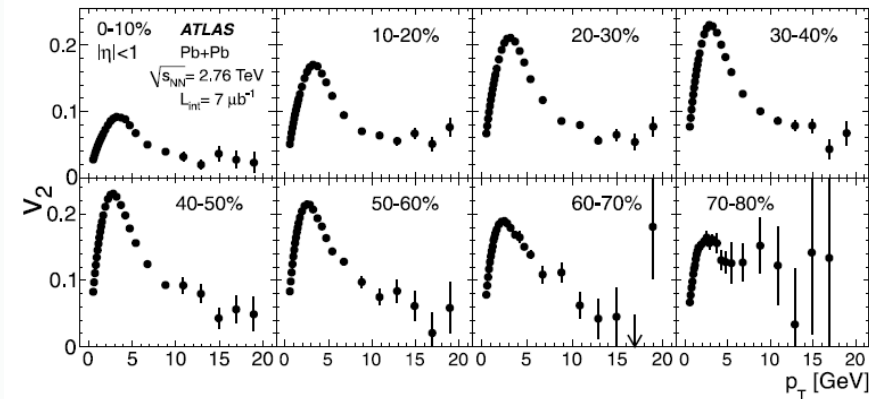
v_2 at LHC: ALICE



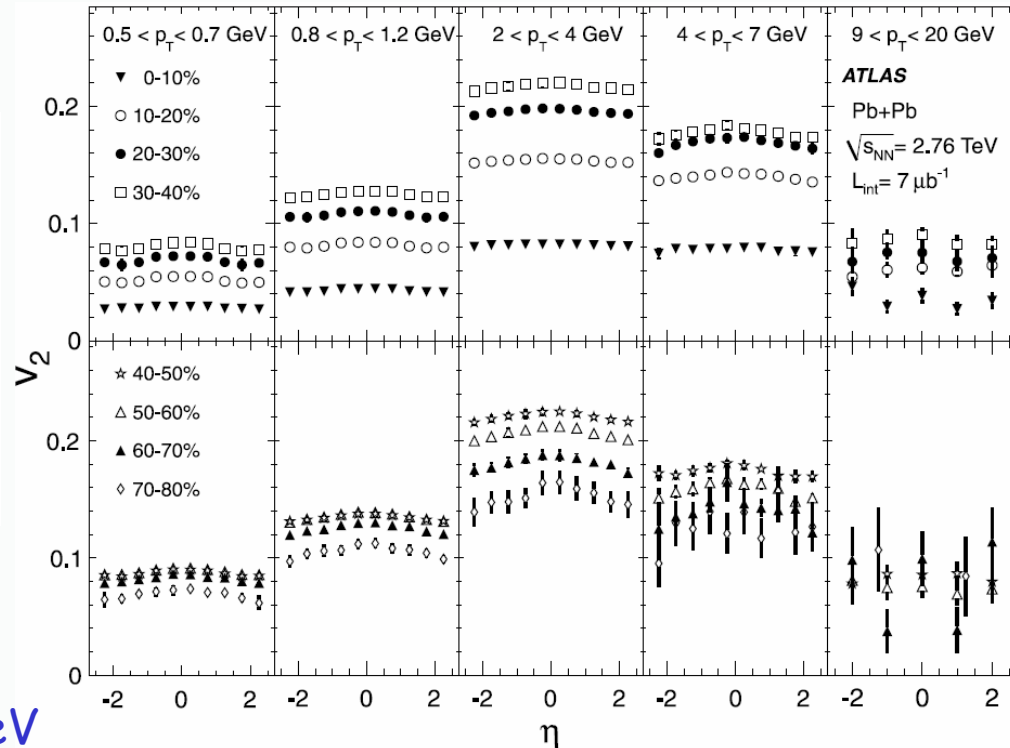
30% increase in the magnitude of integrated v_2 at $\sqrt{s_{NN}}=2.76$ TeV compared to RHIC.
Larger $\langle p_T \rangle$ at LHC.

The value of $v_2(p_T)$ does not change within uncertainties from $\sqrt{s_{NN}}=200$ GeV to 2.76 TeV

v_2 at LHC: ATLAS

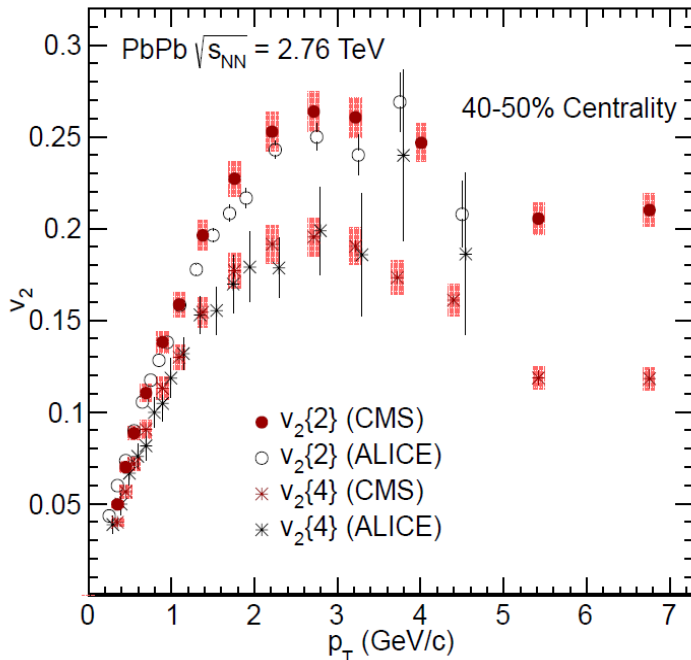


Rapid rise in $v_2(p_T)$ up to $p_T = 3$ GeV,
 A decrease out to 7-8 GeV
 A weak p_T dependence beyond 9-10 GeV

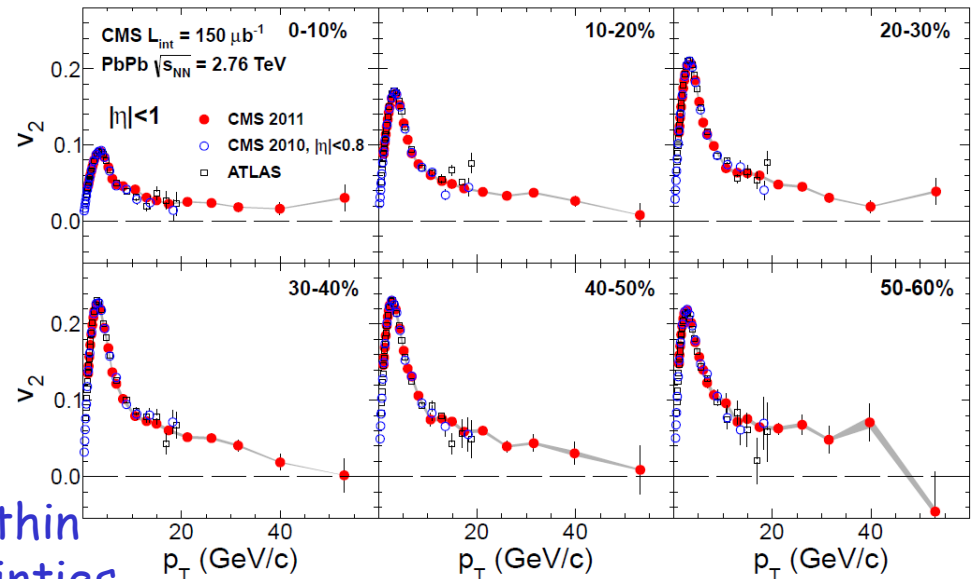
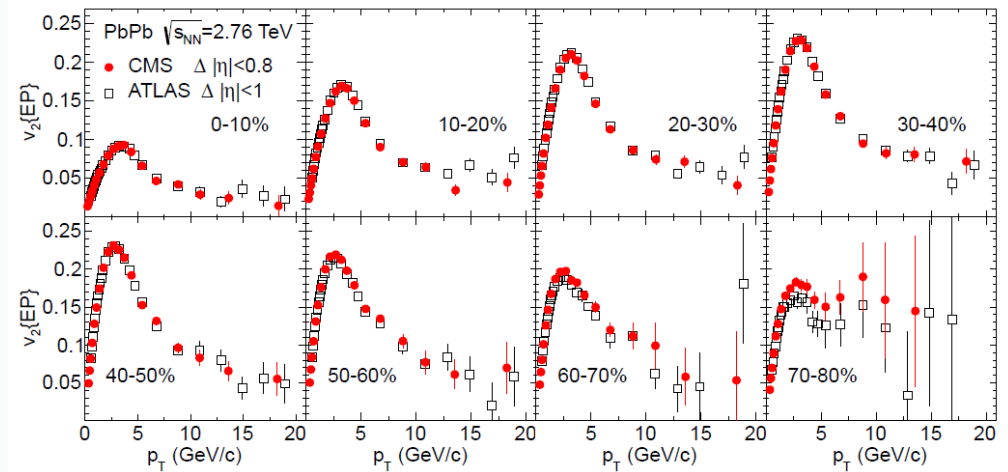


v_2 depends very weakly on η over
 the measured pseudorapidity region

v_2 at LHC



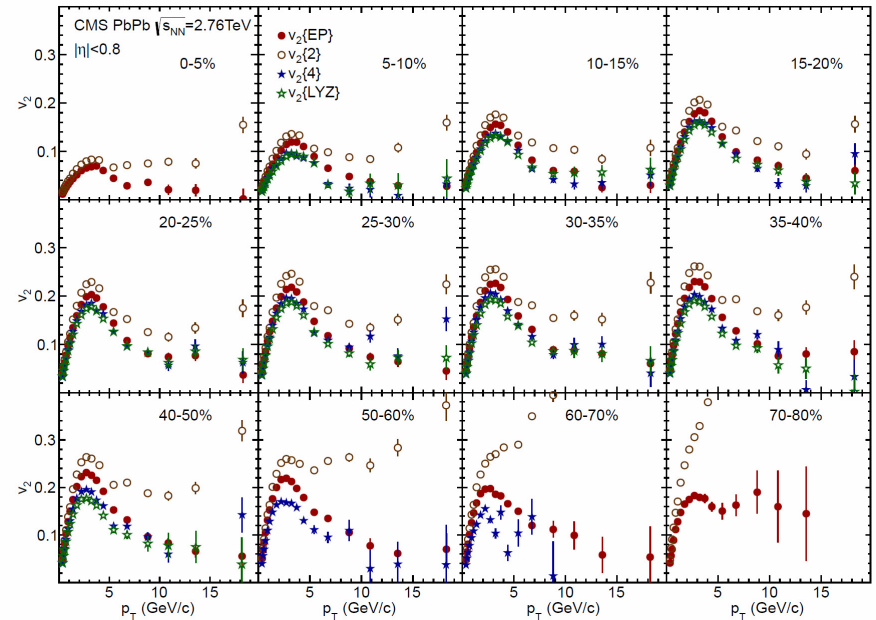
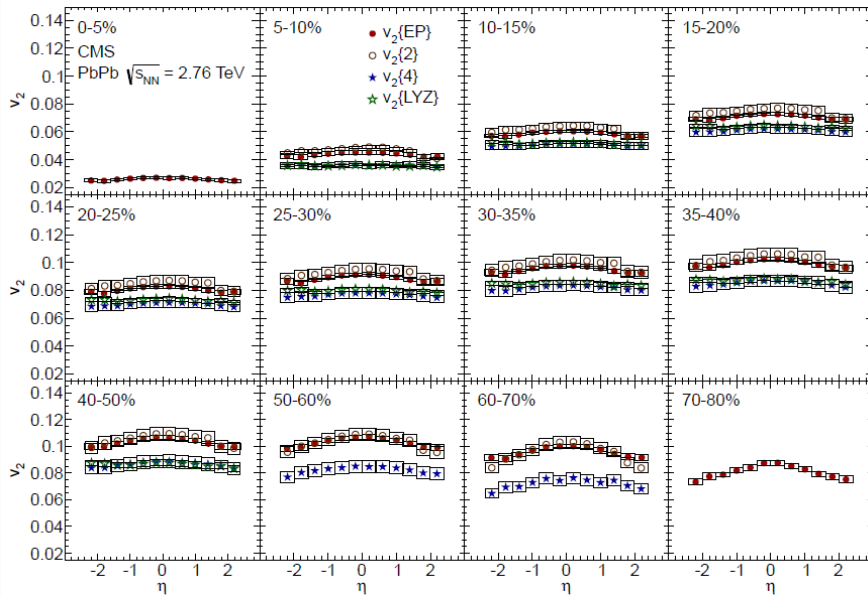
Phys. Rev. Lett. **105** (2010) 252302,
Phys. Lett. B **707** (2012) 330,
arXiv:1204.1409, arXiv:1204.1850

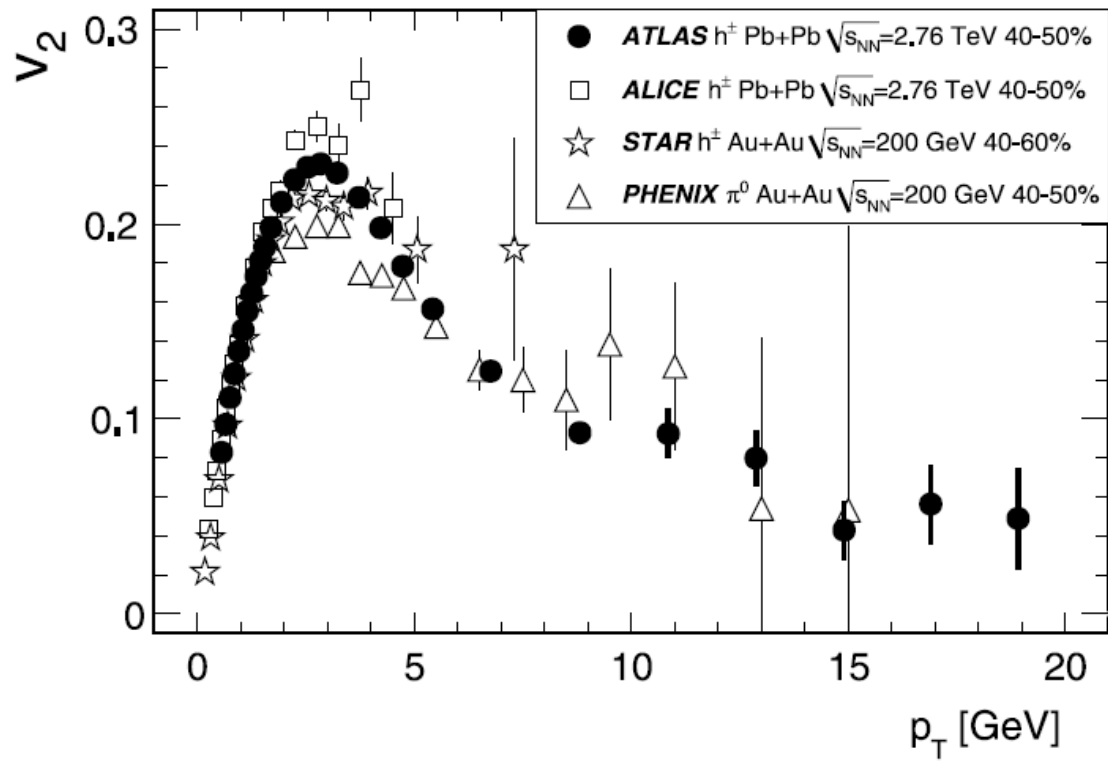


The results are in good agreement within the statistical and systematic uncertainties

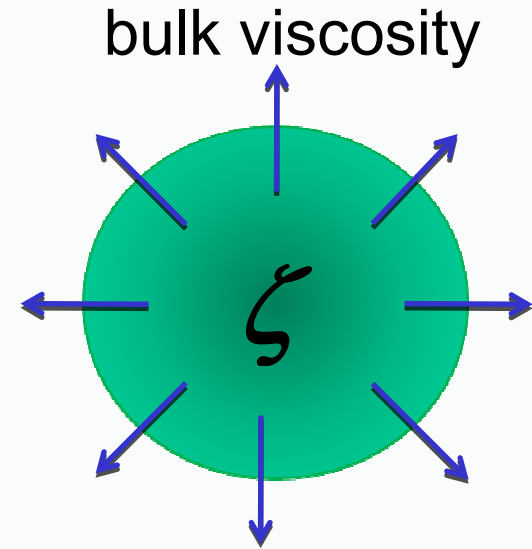
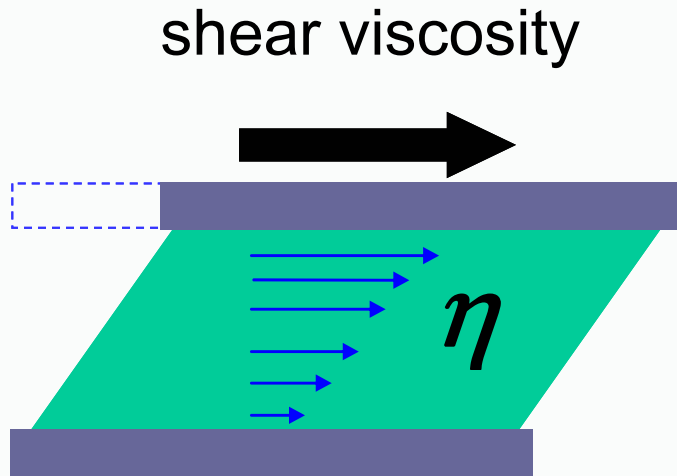
Backup

v_2 at LHC:CMS





Viscous hydro with shear & bulk viscosity



Conservation laws:

$$T^{\mu\nu} = (e + p + \Pi)u^\mu u^\nu - (p + \Pi)g^{\mu\nu} + \pi^{\mu\nu}$$

Evolution equations for shear pressure tensor $\pi^{\mu\nu}$ and bulk pressure:

$$\tau_\pi \Delta^{\alpha\mu} \Delta^{\beta\nu} \dot{\pi}_{\alpha\beta} + \boxed{\pi^{\mu\nu} = 2\eta\sigma^{\mu\nu}} - \frac{1}{2} \pi^{\mu\nu} \frac{\eta T}{\tau_\pi} \partial_\lambda \left(\frac{\tau_\pi}{\eta T} u^\lambda \right)$$

$$\tau_\Pi \dot{\Pi} + \boxed{\Pi = -\zeta(\partial \cdot u)} - \frac{1}{2} \Pi \frac{\zeta T}{\tau_\Pi} \partial_\lambda \left(\frac{\tau_\Pi}{\zeta T} u^\lambda \right)$$

(2nd order shear-bulk -mixing term (Muronga, Rischke) not included.)