

Heavy Ion Collisions

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Probing Hot QCD Matter with hard probes

Hard Probes:

“highly penetrating observables (particles, radiation)
used to explore properties of matter that cannot be viewed directly!”

$$p_{T,m} > 2 \text{ GeV} \gg \Lambda_{\text{QCD}}$$

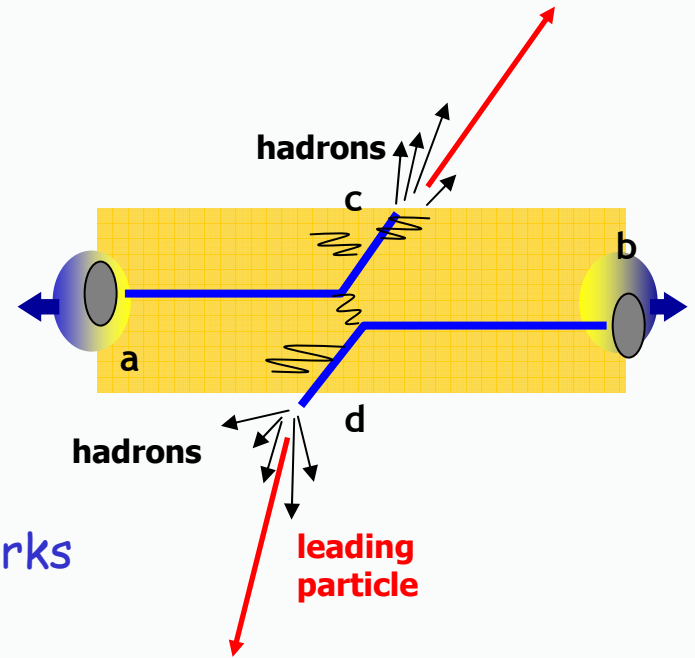
Hard probes

p+p:

- parton scattering \rightarrow fragmentation \rightarrow jet
- can be calculated in perturbative QCD
- collinear factorization

A+A:

- partons traversing medium lose energy
gluon radiation, elastic collisions
- energy loss different for g, light/heavy quarks
(color factor, dead cone effect)



X.-N. Wang, M. Gyulassy, *Phys. Rev. Lett.* **68** (1992) 1480

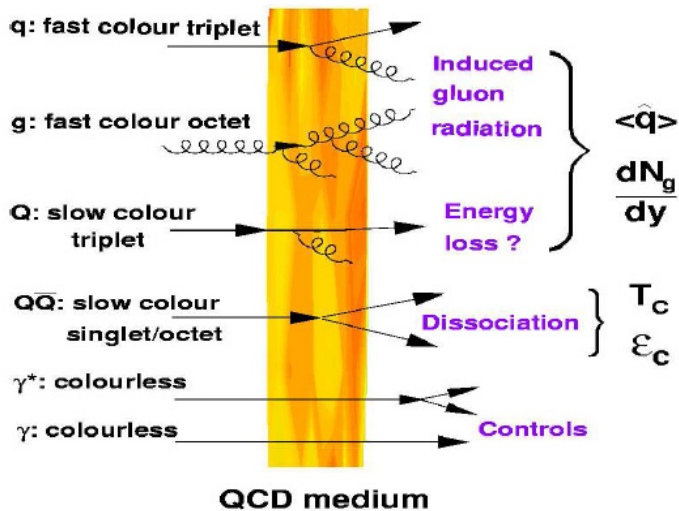
Goal: Use in-medium energy loss to measure medium properties

$$\frac{d\sigma_{pp}^h}{dy d^2p_T} = K \sum_{abcd} \int dx_a dx_b f_a(x_a, Q^2) f_b(x_b, Q^2) \frac{d\sigma}{d\hat{t}}(ab \rightarrow cd) \frac{D_{h/c}^0}{\pi z_c}$$

Parton distribution function	Matrix element	Fragmentation function
measured in DIS initial state (saturation?)	pQCD	e^+e^- final state (energy loss?)

Medium modifications.

$$R_{AA}(\sqrt{s_{NN}}, p_T, y, m; b) = \frac{\text{“hot/dense QCD medium”}}{\text{“QCD vacuum”}} \propto \frac{\Phi_{AA}(\sqrt{s_{NN}}, p_T, y, m; b)}{\Phi_{pp}(\sqrt{s}, p_T, y, m)}$$



Any observed *enhancements* and/or *suppressions* in the $R_{AA}(s_{NN}, p_T, y, m; b)$ ratios can then be directly linked to the properties of strongly interacting matter.

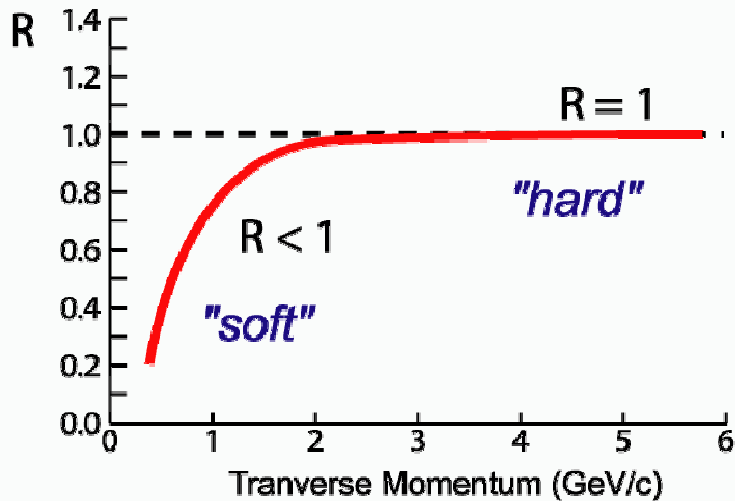
Measurement in pp collisions is essential/ mandatory

Figure 3. Examples of hard probes whose modifications in high-energy AA collisions provide direct information on properties of QCD matter such as the $\langle \hat{q} \rangle$ transport coefficient, the initial gluon rapidity density dN_g/dy , and the critical temperature and energy density.

Medium modifications: R_{AA} , R_{CP}

Nuclear modification factor:

$$R_{AA}(p_t) = \frac{1}{\langle N_{coll} \rangle} \times \frac{dN_{AA} / dp_t}{dN_{pp} / dp_t}$$



No "Effect":

$R < 1$ at small momenta

$R = 1$ at higher momenta where hard processes dominate

Suppression:

$R < 1$

In case pp is not measured, R_{CP} :

$$R_{CP}(p_t) = \frac{1}{\langle N_{coll} \rangle} \times \frac{dN_{AA} / dp_t / \langle T_{AA} \rangle [central]}{dN_{AA} / dp_t / \langle T_{AA} \rangle [peripheral]}$$

Hard cross sections in pp at RHIC

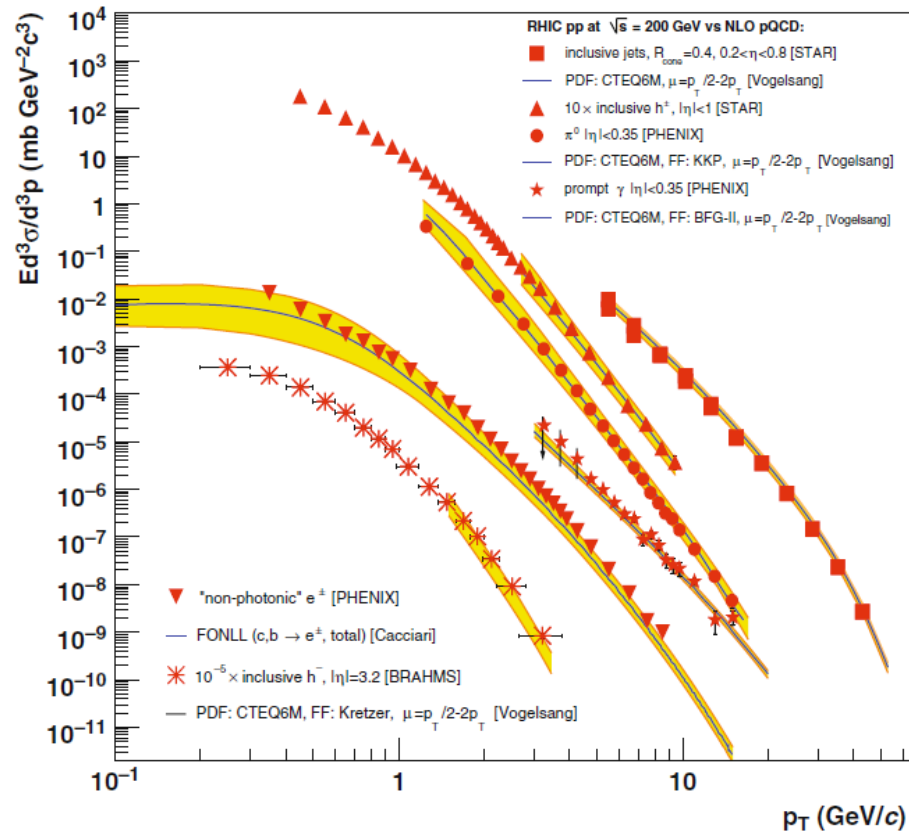


Fig. 13 Compilation of hard cross sections in pp at $\sqrt{s} = 200$ GeV measured by STAR [125, 126], PHENIX [127–129], and BRAHMS [130] (10–30% syst. uncertainties not shown for clarity) compared to NLO [131–138] and NLL [139] pQCD predictions (*yellow bands*)

Discovery of jet quenching at RHIC

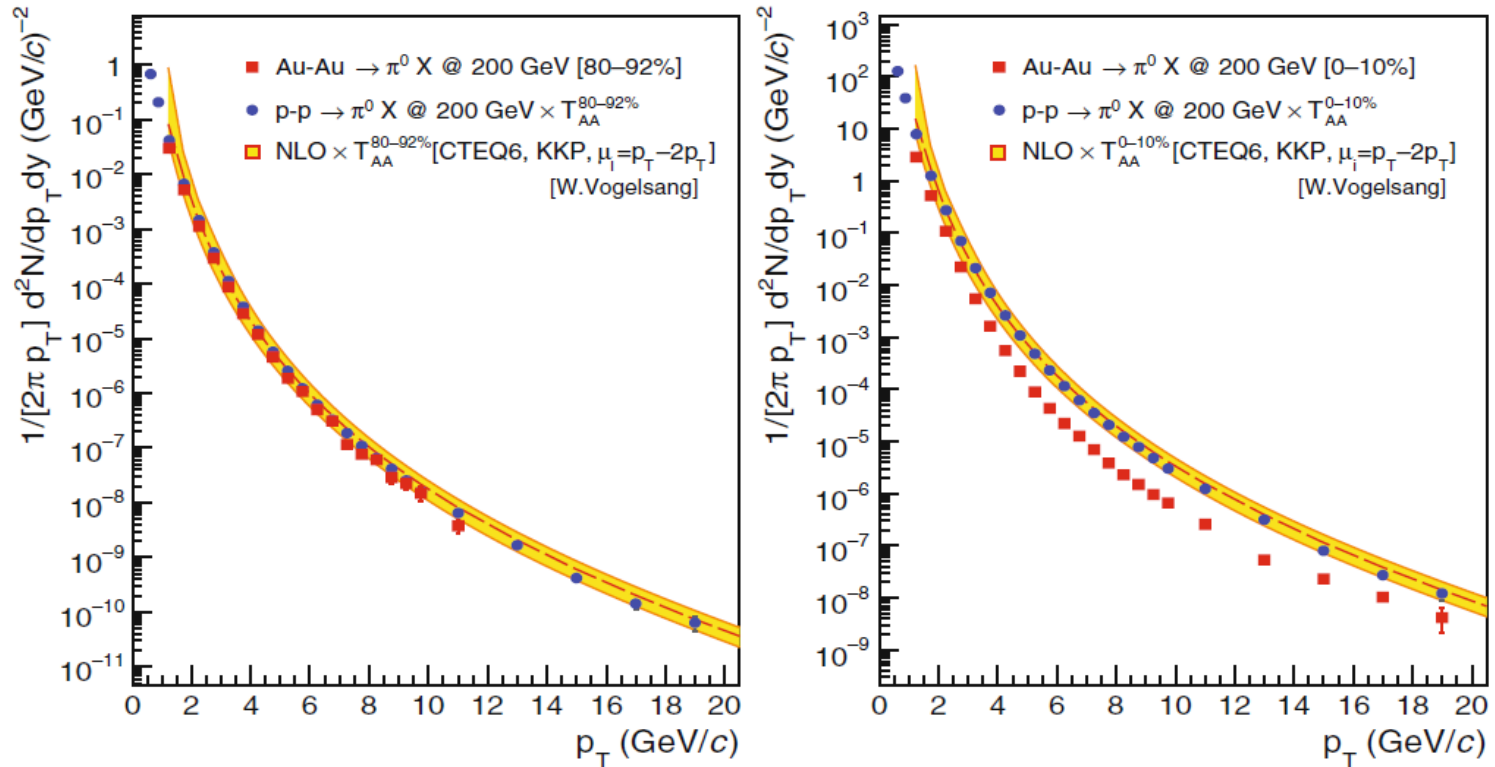
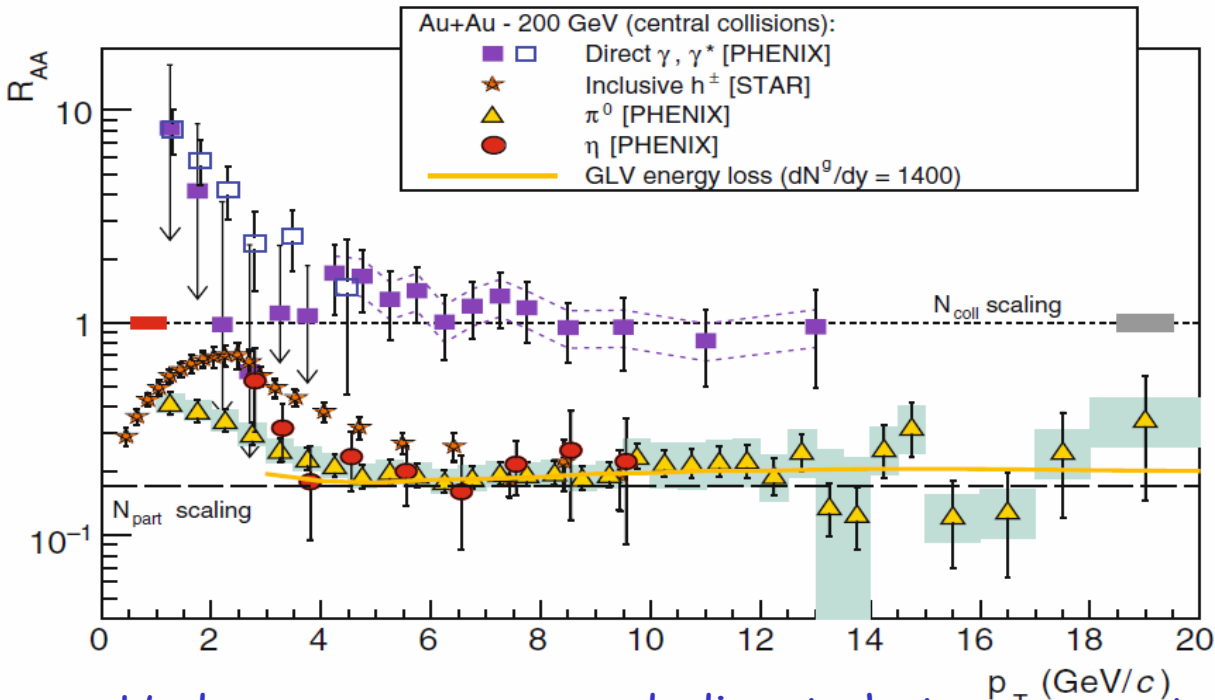


Fig. 15 Invariant π^0 yields measured by PHENIX in peripheral (*left*) and central (*right*) AuAu collisions (*squares*) [100] compared to the (T_{AA} -scaled) pp π^0 cross section (*circles*) [152] and to a NLO pQCD calculation (*curves and yellow band*) [131–133]

Large high- p_T hadron suppression
 ($R_{AA} < 1$) observed in central AuAu compared to pp or dAu reactions.

Discovery of jet quenching at RHIC



PHENIX: Phys.Rev.Lett.88:022301, 2002
 PHENIX: Phys.Rev.Lett.91:072301, 2003
 PHENIX: Phys.Rev.Lett.94:232301, 2005
 STAR: Phys.Rev.Lett.89:202301,2002
 STAR: Phys.Rev.Lett.90:082302,2003
 STAR: Phys.Rev.Lett.91:172302,2003

$$R_{AA}^\gamma \sim 1$$

$$R_{AA}^{\pi^0, \eta} \sim 0.2$$

- Hadrons are suppressed, direct photons are not
- The hadron spectra at RHIC from p+p, Au+Au and d+Au collisions establish existence of *parton energy loss* from strongly interacting, dense QCD matter in central Au-Au collisions

$$\varepsilon_{\text{loss}} \approx 1 - R_{AA}^{1/(n-2)}$$

https://wiki.bnl.gov/TECHQM/index.php/Main_page

Theory-Experiment Collaboration on Hot Quark Matter

$$\langle \hat{q} \rangle = 4 - 13 \text{ GeV}^2 / \text{fm}$$

$$dN^0/dy \sim 1400 \pm 200$$

S. Bass et al. PRC79 (2009) 024901

Discovery of jet quenching at RHIC

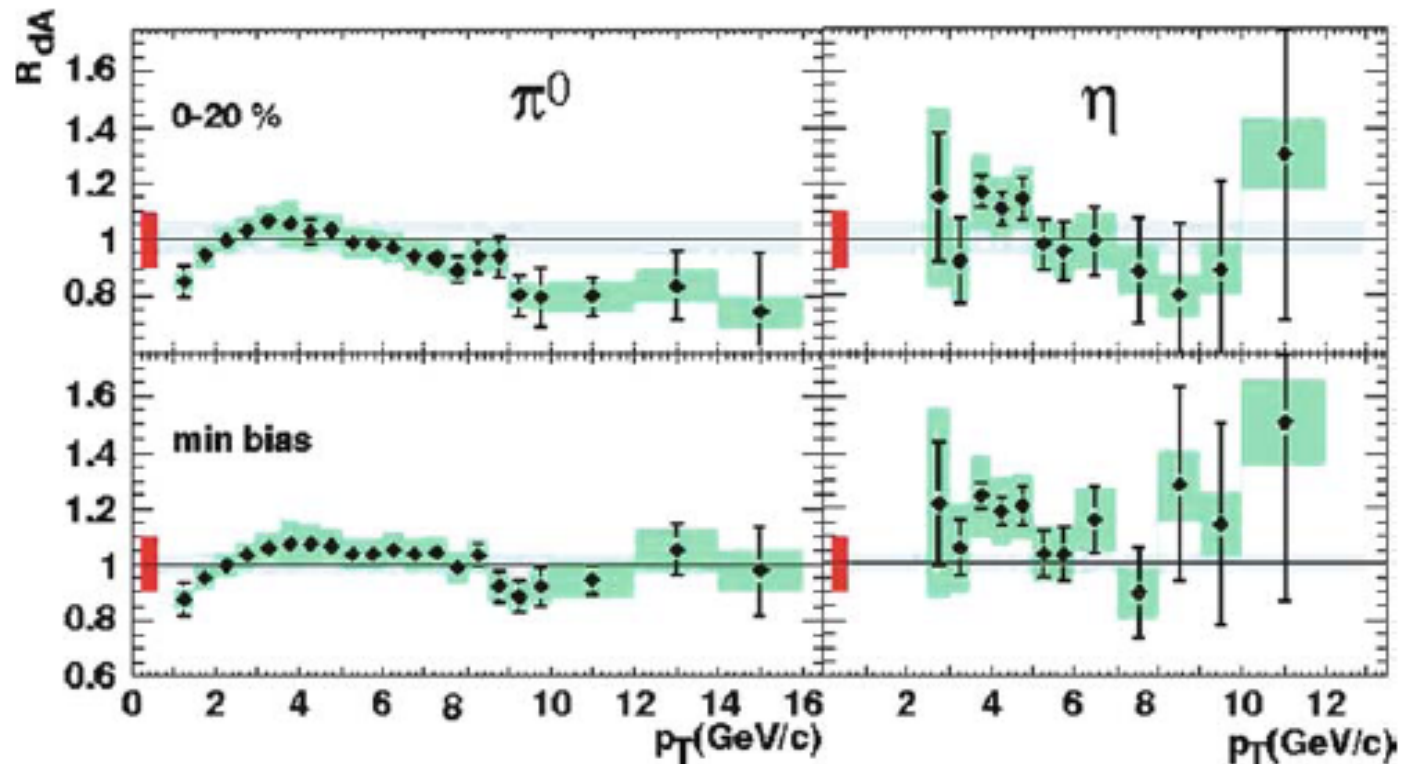


Fig. 14 Nuclear modification factors for high- p_T π^0 (left) and η (right) mesons at midrapidity in dAu collisions at $\sqrt{s_{NN}} = 200$ GeV [143, 144] compared to pQCD calculations [145, 146] with EKS98 [147] nuclear PDFs

No suppression in dAu . Evidence for final state effect.

Energy dependence of the suppression

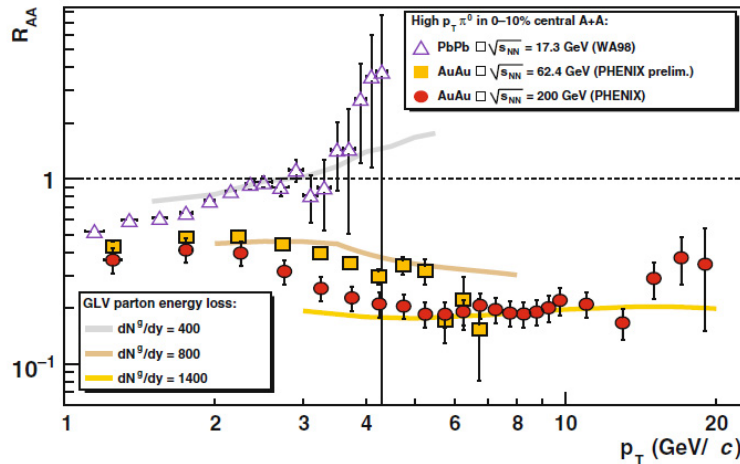


Fig. 18 Nuclear modification factor, $R_{AA}(p_T)$, for neutral pions in central PbPb at $\sqrt{s_{NN}} = 17.3$ GeV [168, 169] and AuAu at $\sqrt{s_{NN}} = 62.4$ GeV [170], 200 GeV [153]; compared to GLV energy loss calculations for initial gluon densities: $dN^g/dy = 400, 800, 1400$ [160, 161, 171], respectively. Experimental normalisation errors, $\mathcal{O}(10\text{--}25\%)$, not shown

Table 2 Initial gluon densities dN^g/dy [160, 161, 171], and transport coefficients $\langle \hat{q} \rangle$ [89] for the dense media produced in central AA collisions at SPS and RHIC energies obtained from parton energy loss calculations reproducing the observed high- p_T π^0 suppression at each $\sqrt{s_{NN}}$. The measured charged particle densities at midrapidity, $dN_{ch}^{exp}/d\eta$ [166, 167], are also quoted

	$\sqrt{s_{NN}}$ (GeV)	$\langle \hat{q} \rangle$ (GeV ² /fm)	dN^g/dy	$dN_{ch}^{exp}/d\eta$
SPS	17.3	3.5	400	312 ± 21
RHIC	62.4	7.	800	475 ± 33
RHIC	130.	~ 11	~ 1000	602 ± 28
RHIC	200.	13	1400	687 ± 37

Heavy quarks as medium probes: Energy Loss

q: colour triplet

u,d,s: $m \sim 0$, $C_R = 4/3$
(difficult to tag at LHC)

g: colour octet

g: $m = 0$, $C_R = 3$
> E loss, dominant at LHC

Q: colour triplet

c: $m \sim 1.5$ GeV, $C_R = 4/3$
small m , tagged by D's

b: $m \sim 5$ GeV, $C_R = 4/3$
large mass, < E loss

QCD medium

Parton Energy Loss by

- medium-induced gluon radiation
- collisions with medium gluons

$$\Delta E(\varepsilon_{QGP}; C_R, m, L)$$

pred: $\Delta E_g > \Delta E_{c \approx q} > \Delta E_b$



$$R_{AA}^{\pi} < R_{AA}^D < R_{AA}^B$$

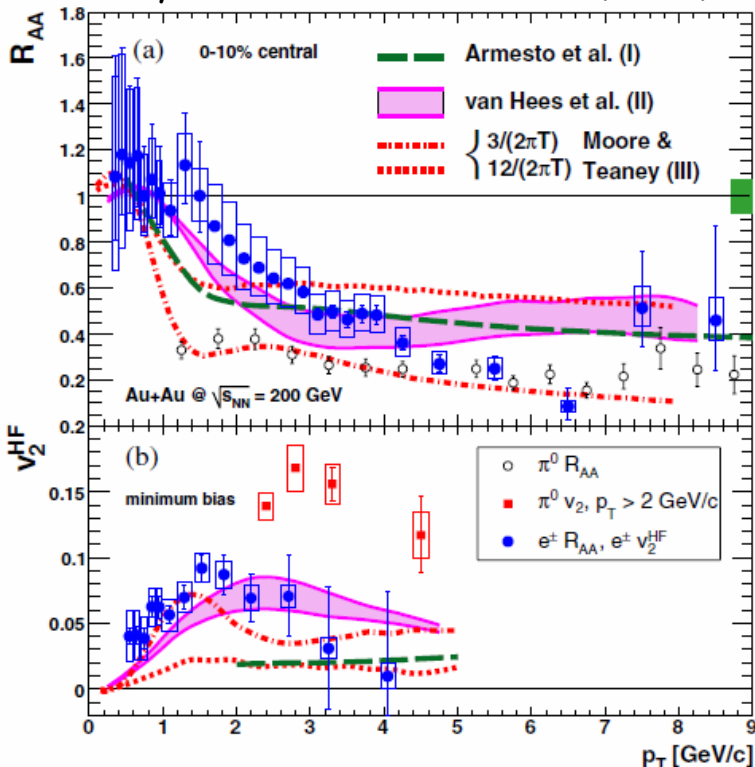
Dokshitzer and Kharzeev, PLB 519 (2001) 199. Armesto, Salgado, Wiedemann, PRD 69 (2004) 114003.

Djordjevic, Gyulassy, Wicks, nucl-th/0512076

Heavy flavour electrons

- The inclusive electron spectra consist of
- (i) "nonphotonic" electrons from heavy-flavor decays,
 - (ii) "photonic" background from Dalitz decays and photon conversions (mainly in the beam pipe), and (iii) nonphotonic background from $K \rightarrow e \pi \nu$ and dielectron decays of vector mesons.

Phys. Rev. Lett. 98, 172301 (2007)



$$\Delta E_{\text{loss}}(g) > \Delta E_{\text{loss}}(q) > \Delta E_{\text{loss}}(Q)$$

(color factor) (dead-cone effect)

But, strong suppression of heavy flavour electrons observed from $2 < p_T < 5$ GeV/c

Models have difficulties to explain both R_{AA} and v_2

van Hees *et al.*: only elastic scattering mediated by resonance excitation of D and B-like states

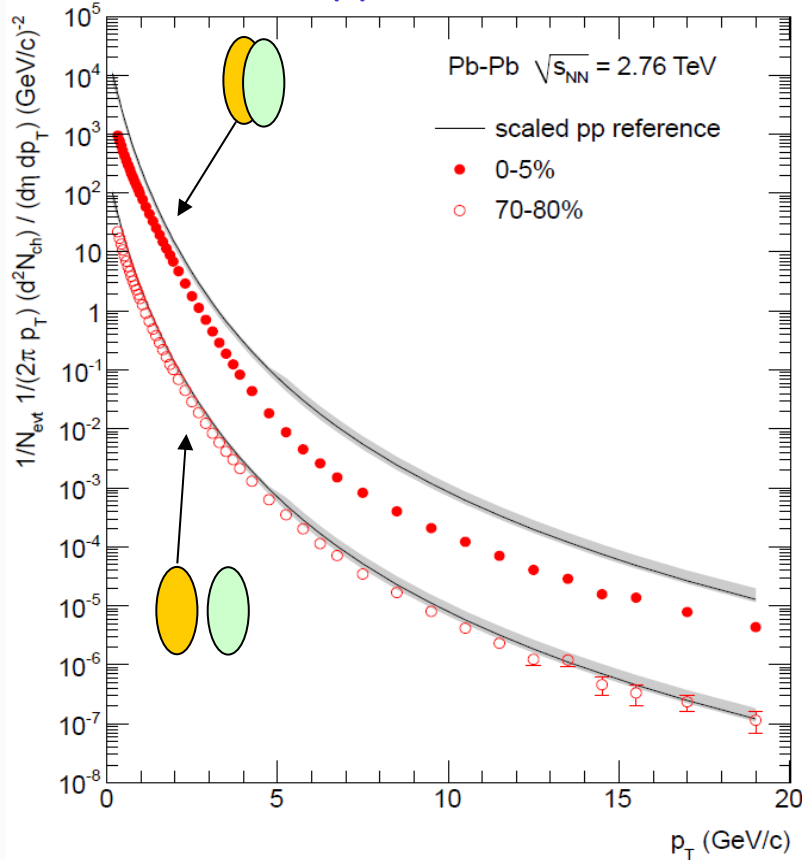
RAA@LHC: ALICE

Suppression of high p_T particles (\sim leading jet fragments)

Rising with p_T !

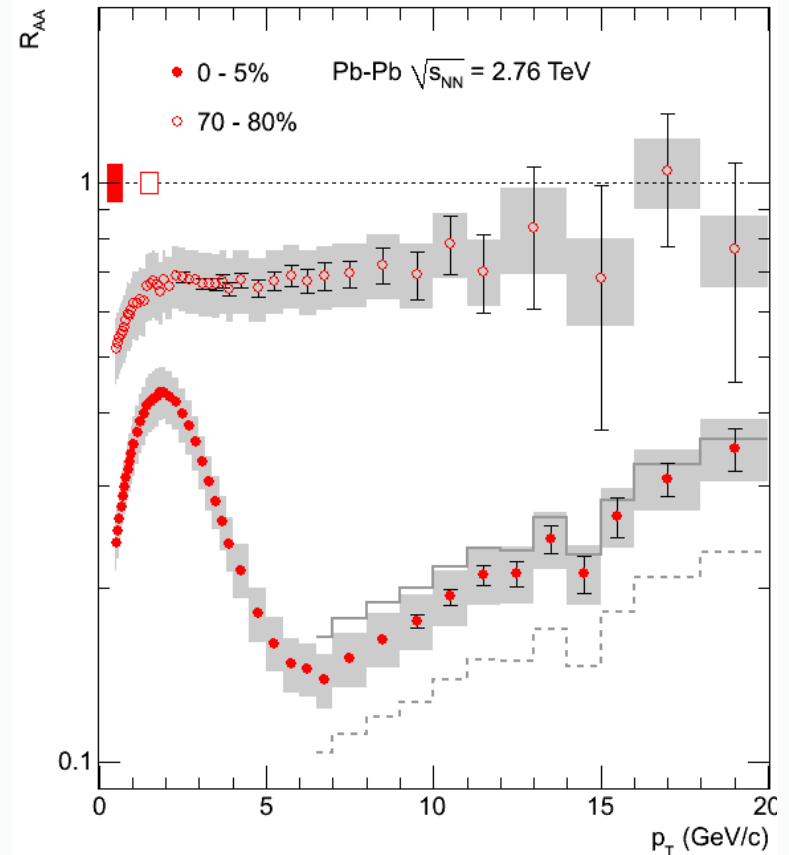
Accuracy limited by pp reference

\Rightarrow need pp at 2.76 TeV !



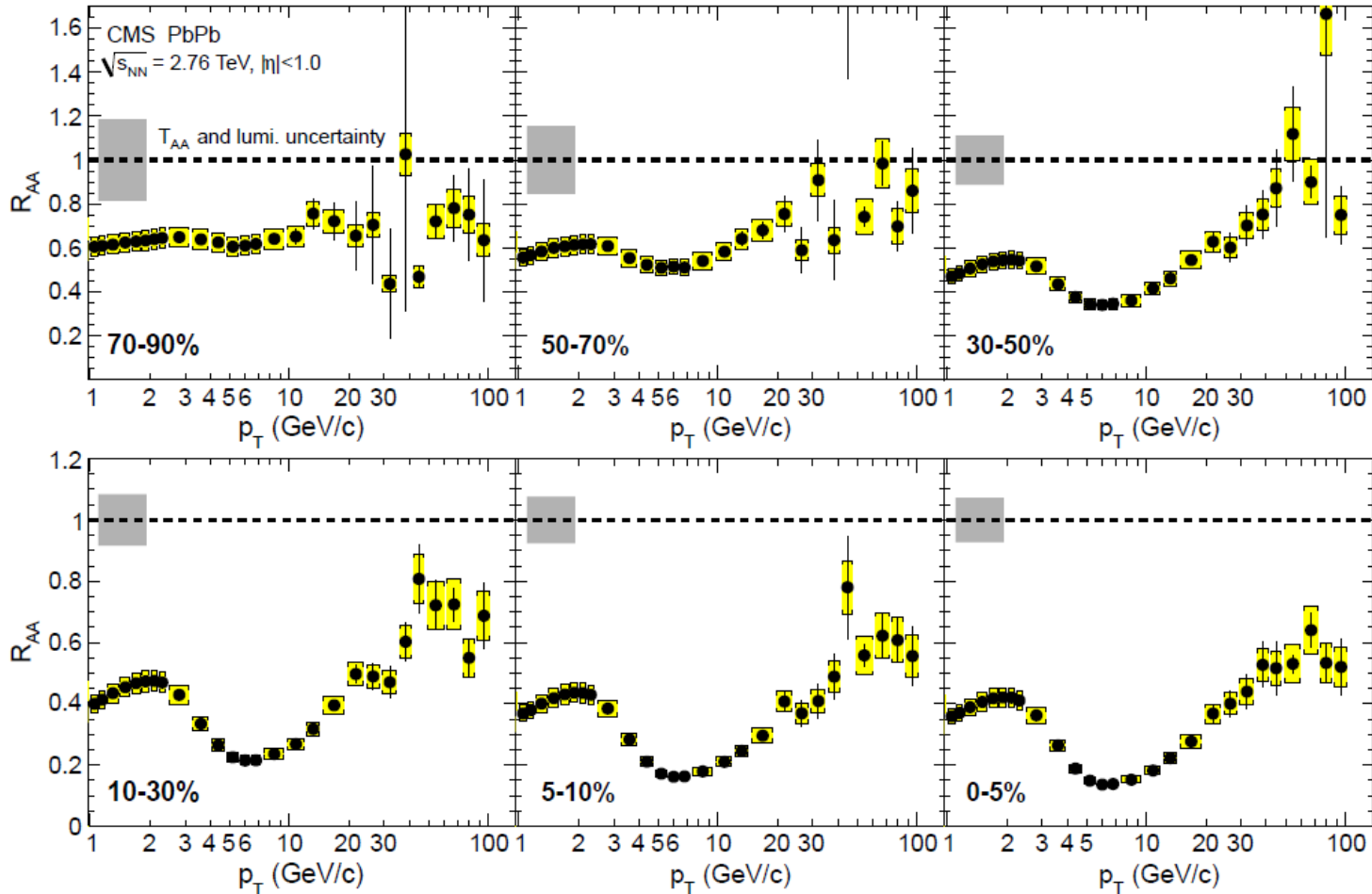
ALICE, Phys. Lett. B 696 (2011) 30

$$R_{AA}(p_T) = \frac{(1/N_{evt}^{AA}) d^2N_{ch}^{AA} / d\eta dp_T}{\langle N_{coll} \rangle (1/N_{evt}^{pp}) d^2N_{ch}^{pp} / d\eta dp_T}$$



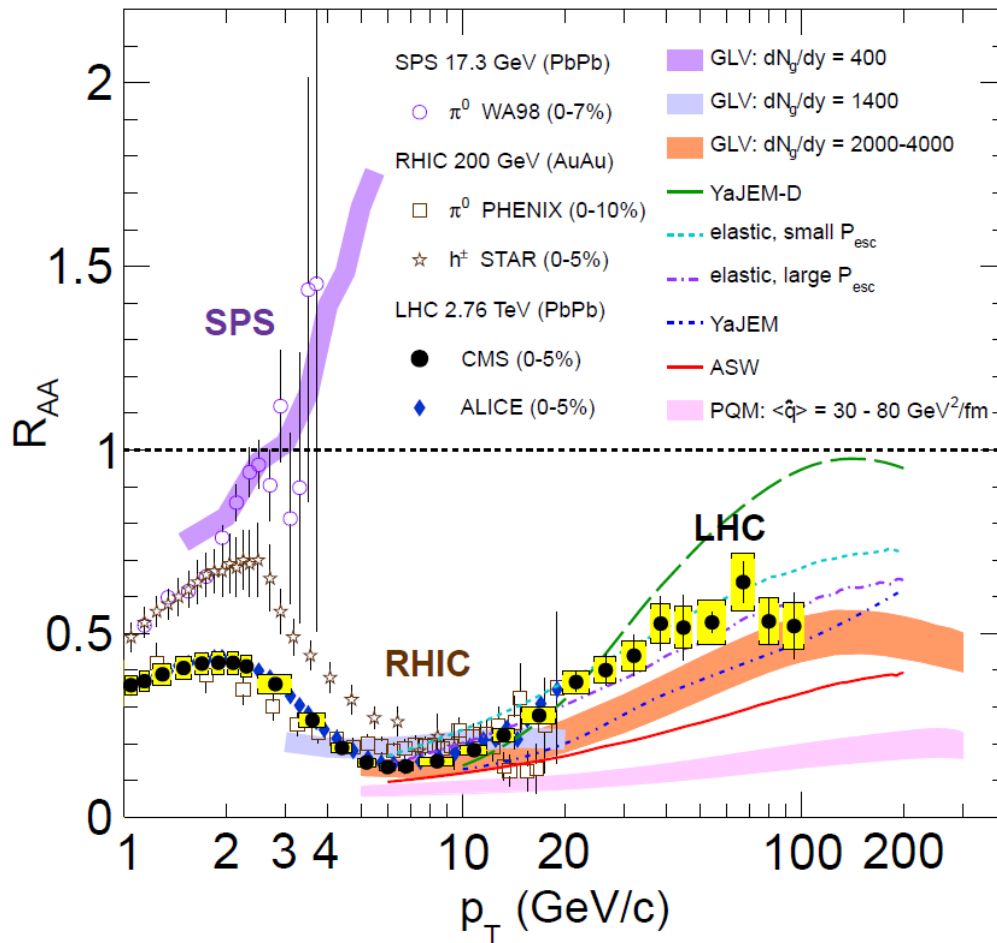
RAA@LHC: CMS

Eur. Phys. J. C (2012) 72:1945



In 0-5% centrality R_{AA} reaches a minimum value of about 0.13 at $p_T = 6-7 \text{ GeV/c}$.
At higher p_T , the value of R_{AA} rises and levels off above 40 GeV/c at ~ 0.5 .

Measurements of R_{AA} in central HIC

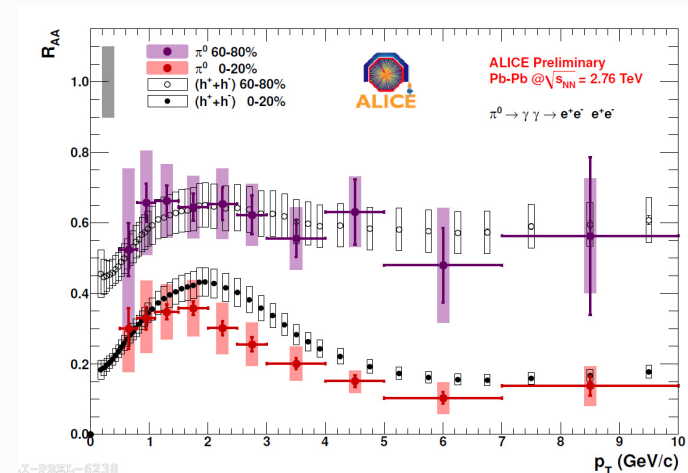
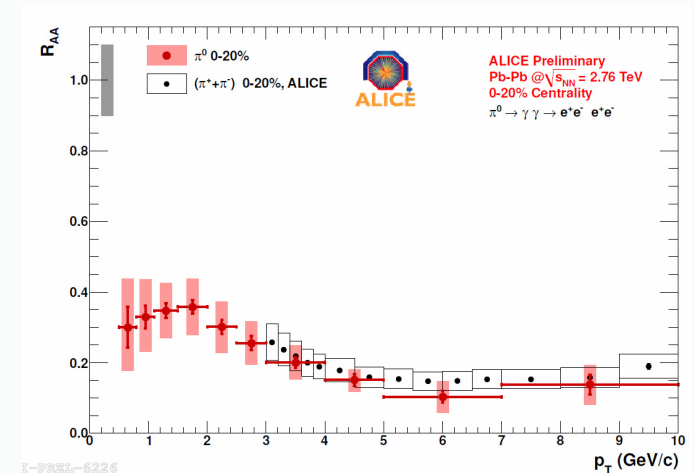
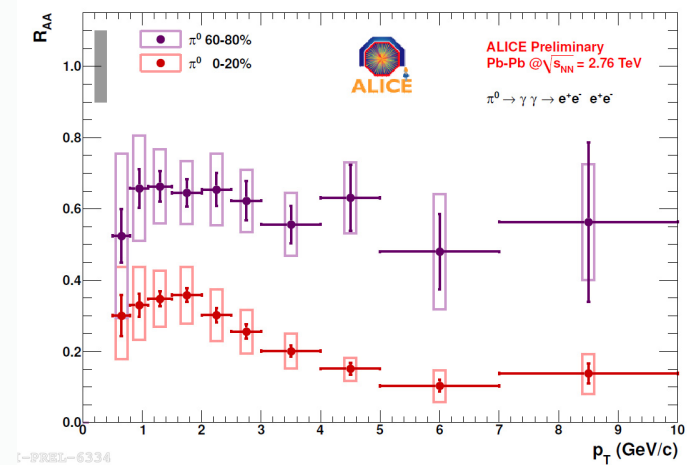
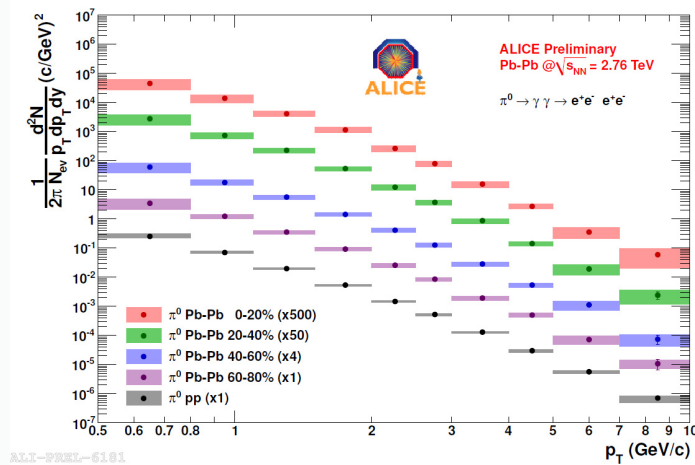


Eur. Phys. J. C (2012) 72:1945

- Larger suppression at LHC than at RHIC.
- Most models predict the generally rising behavior of R_{AA}
- Measurement can be used to constrain the quenching parameters used in these models

Identified particles: π^0 R_{AA}

J.Phys.G G38 (2011) 124117

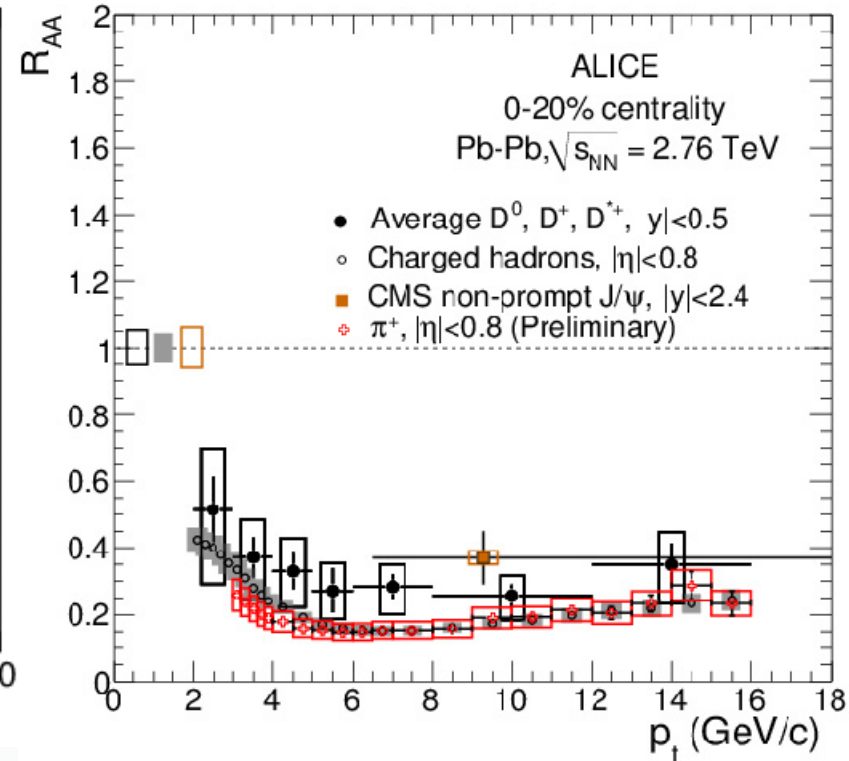
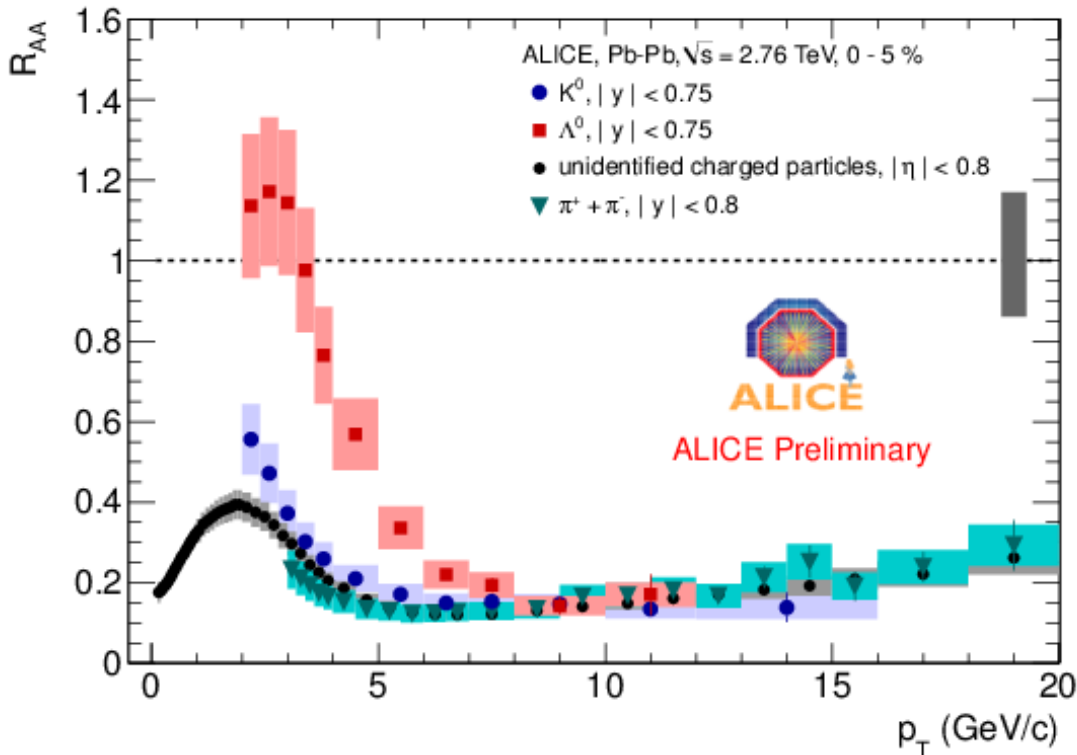


Agreement with charged pion R_{AA}

Different to charged particle R_{AA}

R_{AA} for light/heavy quark hadrons

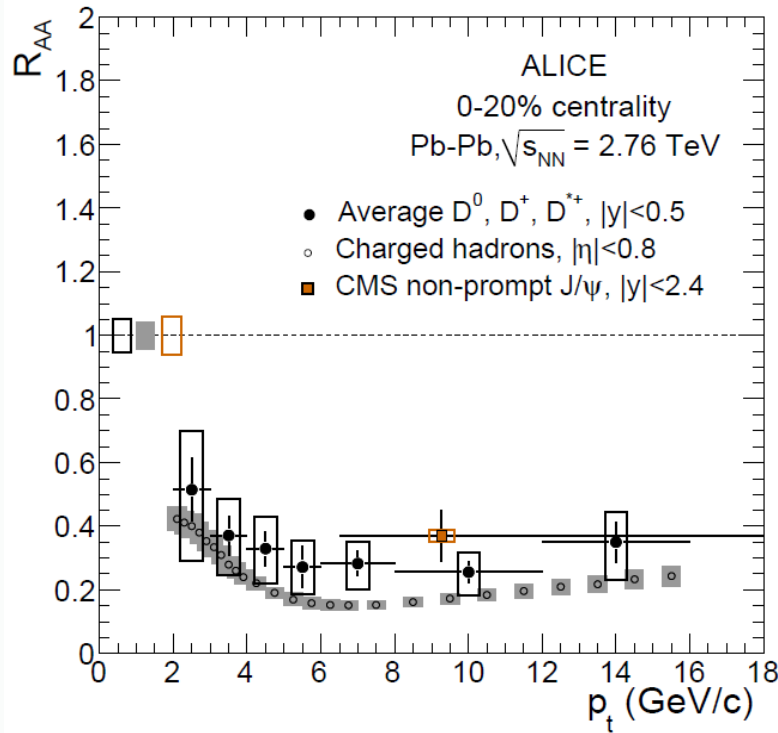
arXiv:1203.2160



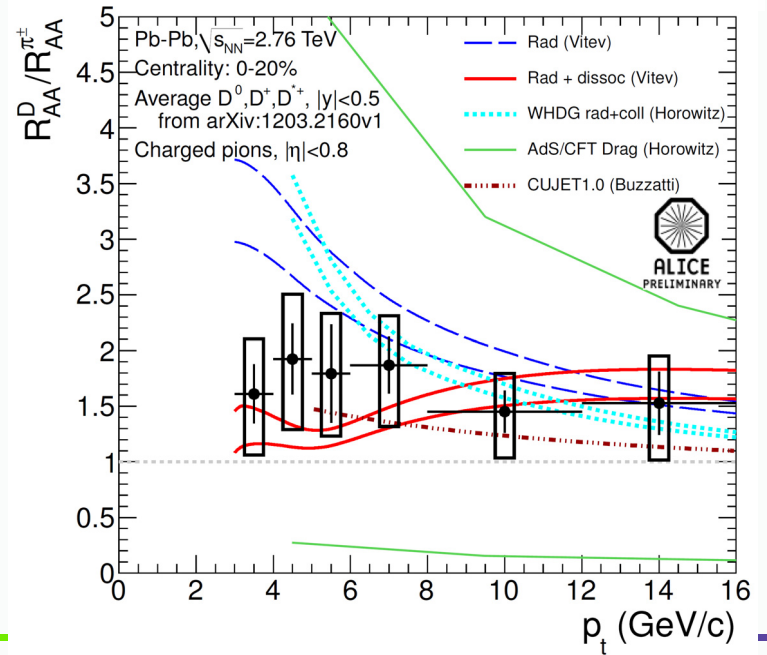
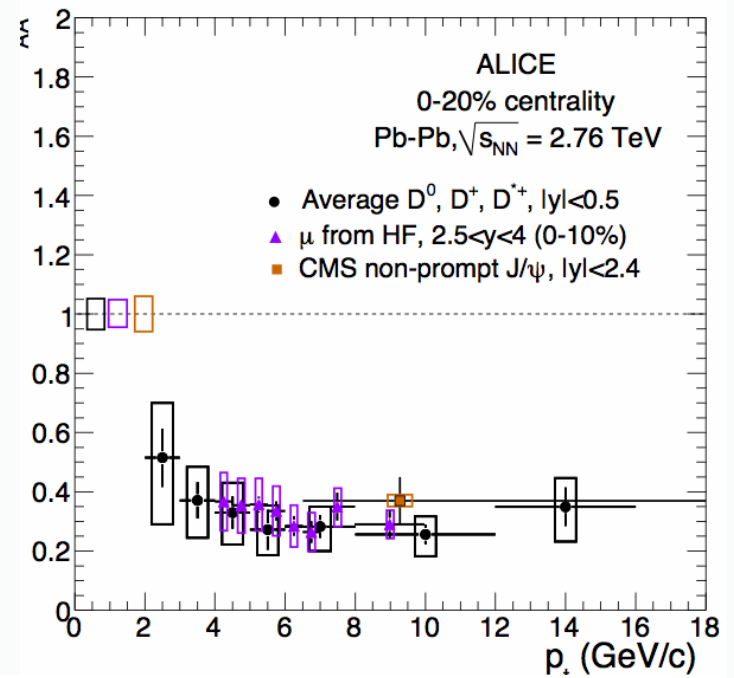
- Light quark hadrons with $p_T > 8$ GeV/c are equally suppressed
- This seem to indicate that medium interactions do not affect fragmentation for $p_T > 8$ GeV/c - fragmentation occurs into vacuum
- Light quark results also provide a baseline for understanding heavy quark energy loss

Heavy/light quark R_{AA}

arXiv:1203.2160

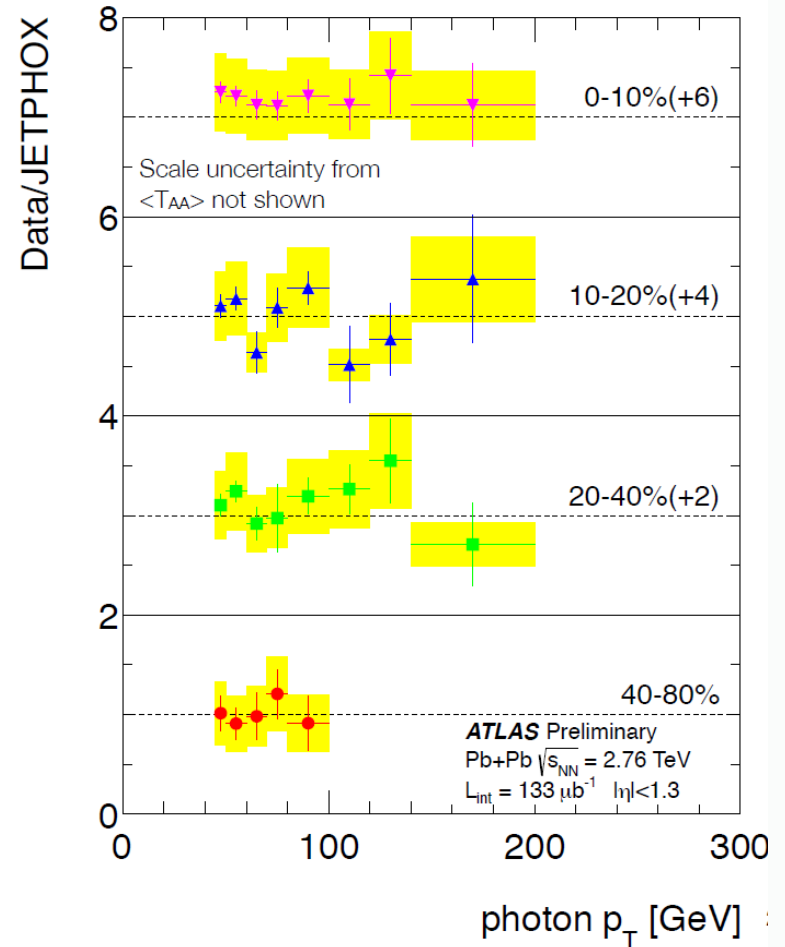
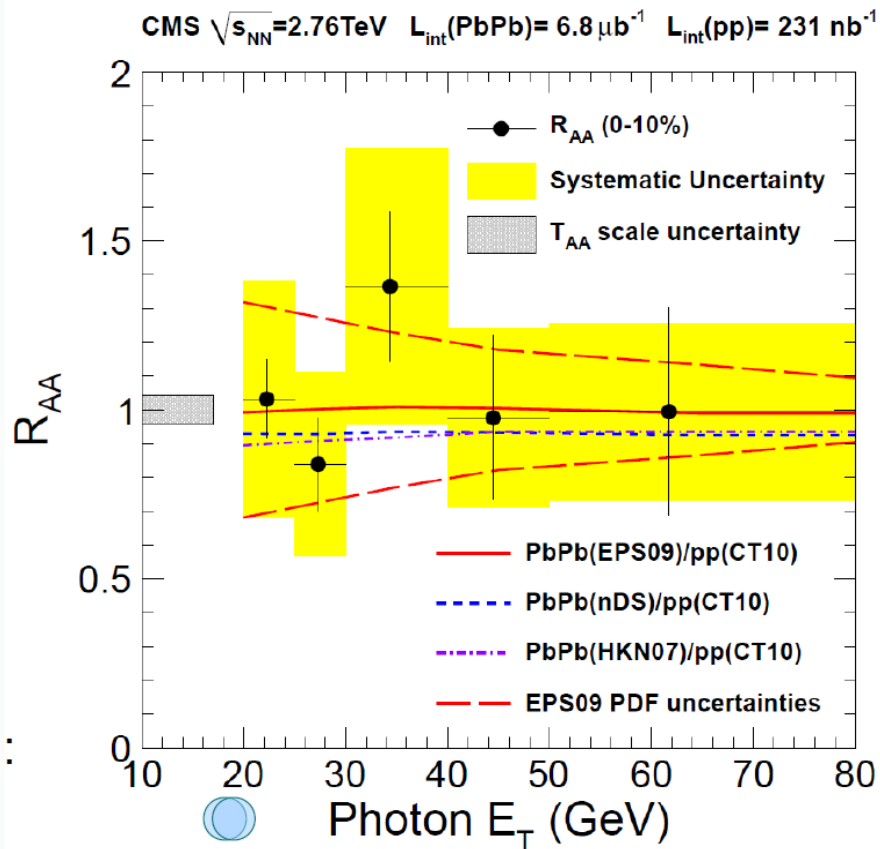


Evidence for $R_{AA}^D / R_{AA}^{\pi} > 1$



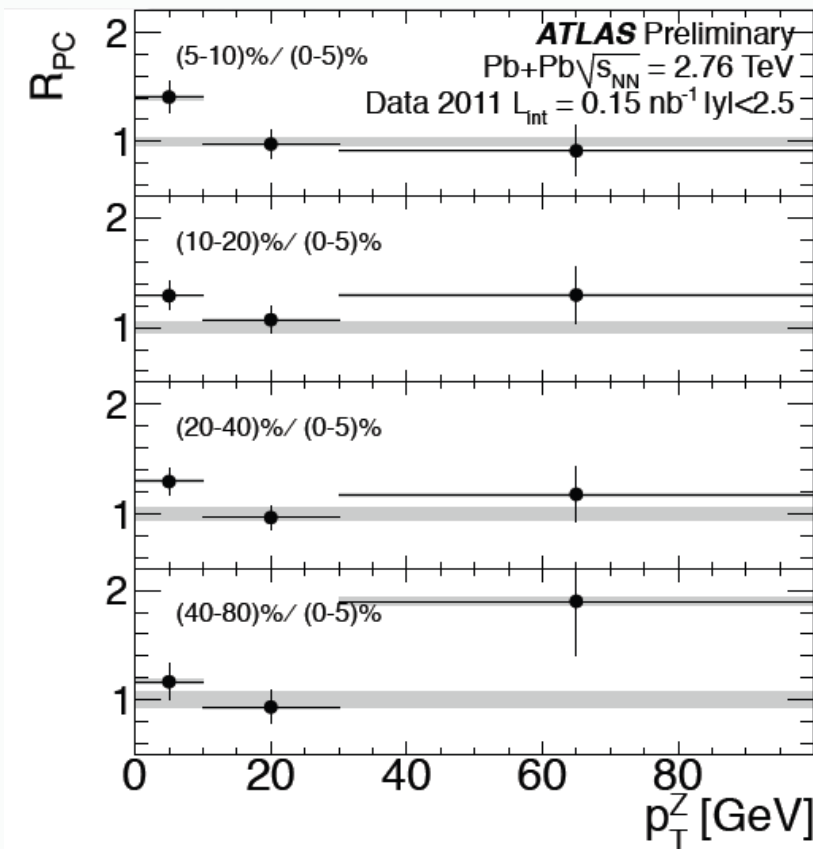
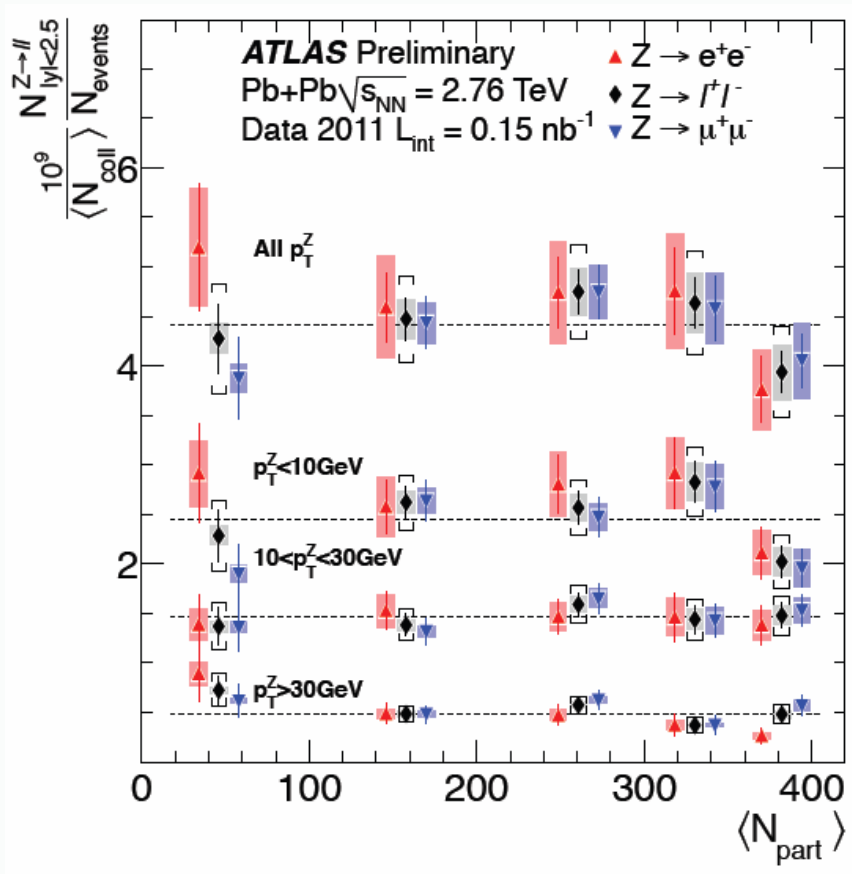
Isolated photon R_{AA} - CMS & ATLAS

PLB 710 (2012) 256



Colorless probes are unaffected

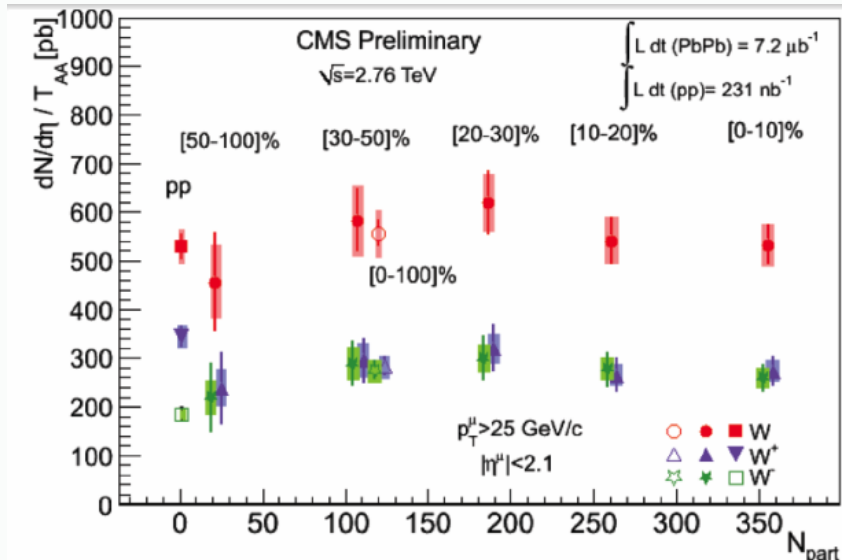
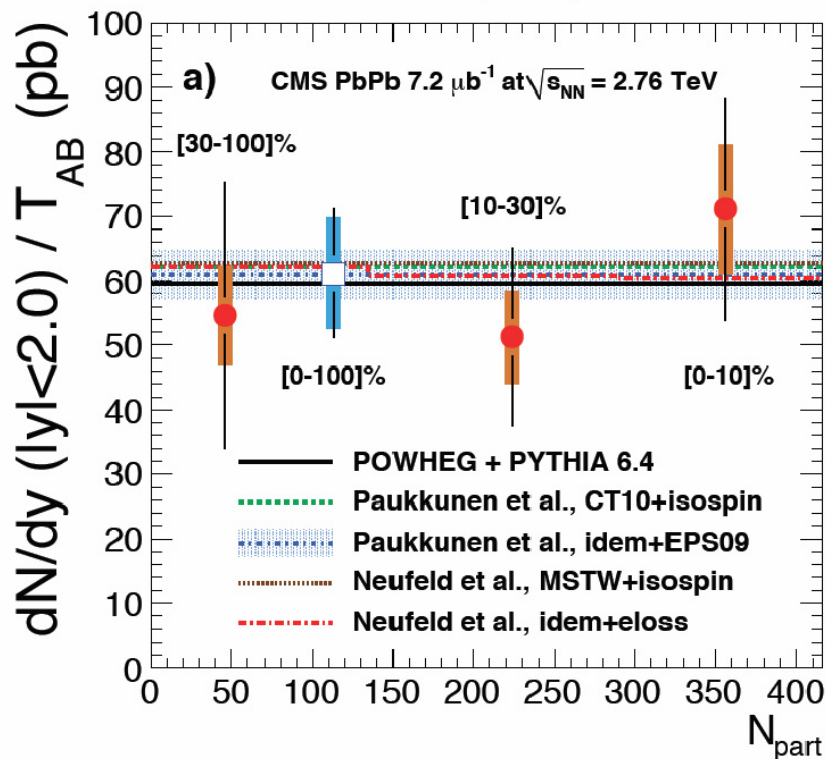
Z boson R_{PC} -ATLAS



Comparing production in different centrality bins consistent with binary collision scaling
 Colorless probes are unaffected

Z and W boson . CMS

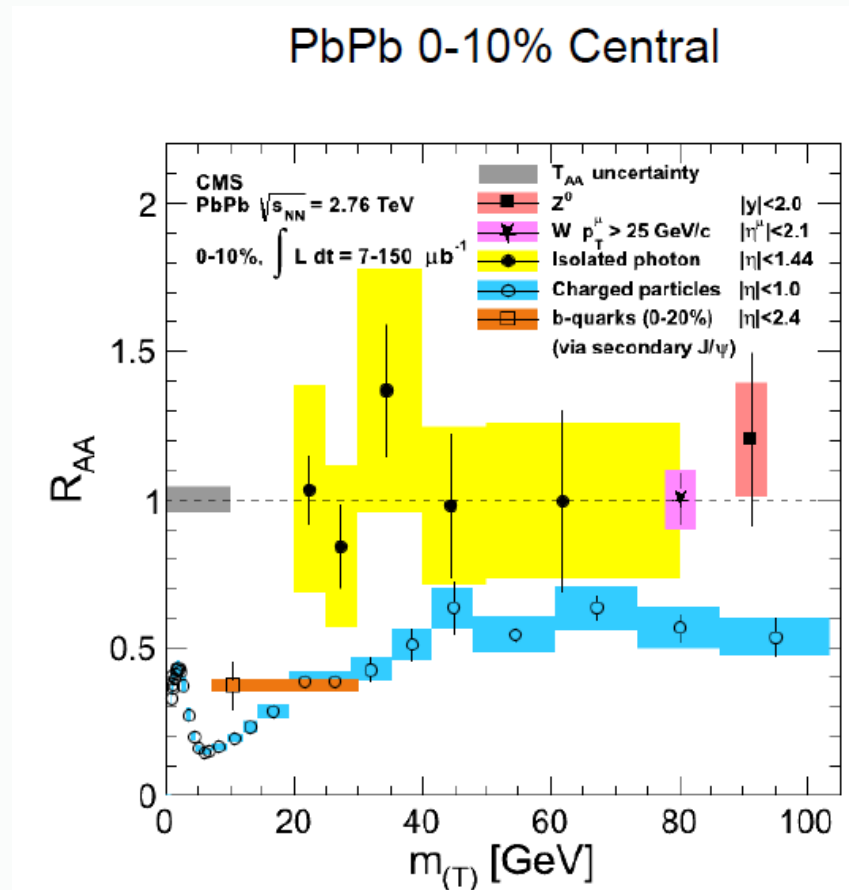
PRL 106 (2011) 212301



Once summed W+ and W- consistent with pp

Colorless probes are unaffected

Compilation of R_{AA} by CMS

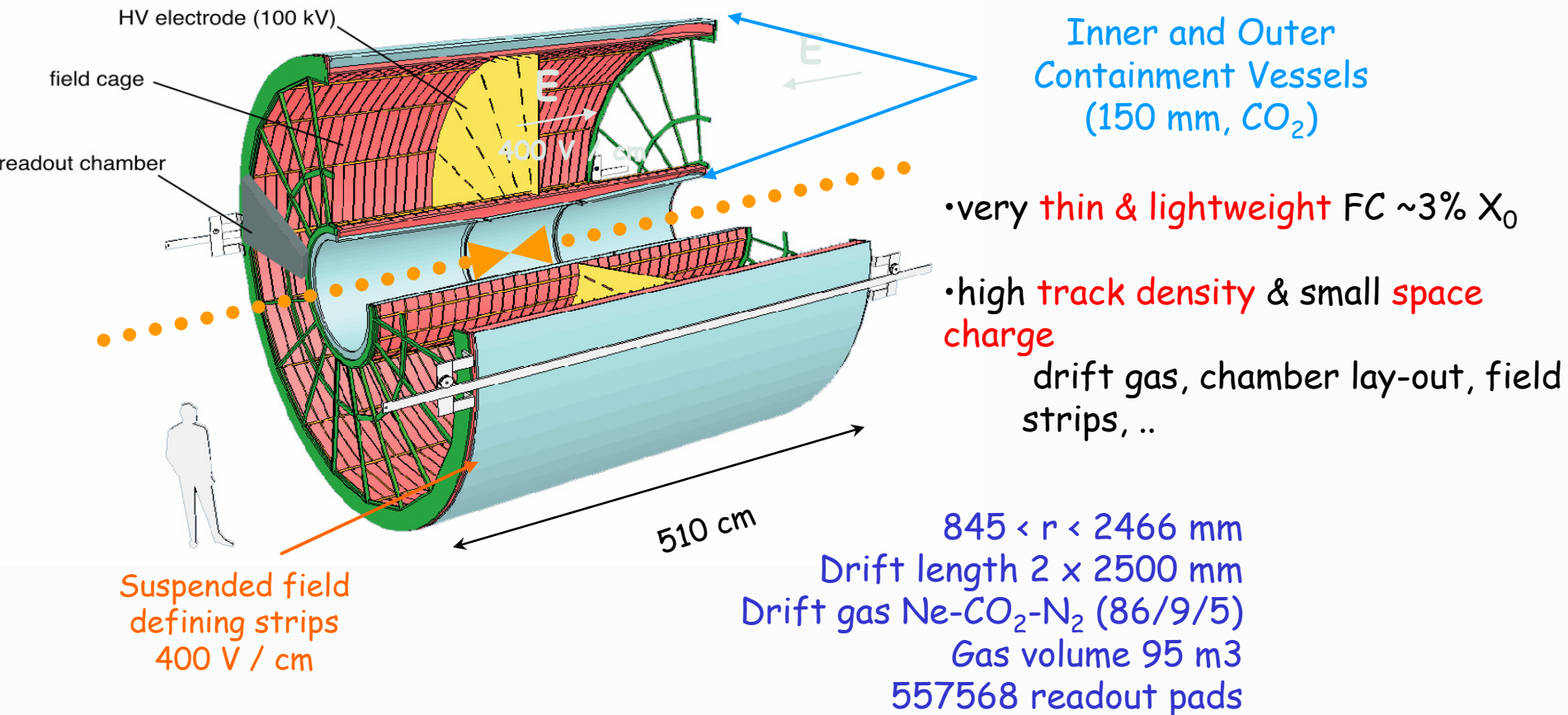


Data Table: <https://twiki.cern.ch/twiki/bin/view/CMSPublic/HIN11002Data>

- **backup**

The Time Projection Chamber

Main tracking detector (charged particles) of the ALICE Central Barrel



Functions, Functions, ...

$$\frac{dN}{p_T dp_T} \propto \left(1 + \frac{p_0}{p_T}\right)^{-n}$$

power law (high- p_T)

$$\frac{dN}{m_T dm_T} \propto m_T K_1\left(\frac{m_T}{T}\right) \xrightarrow{m_T \gg T} \sqrt{m_T} e^{-m_T/T}$$

thermal emission (4π)

$$\frac{dN}{m_T dm_T} \propto m_T e^{-m_T/T}$$

thermal emission ($y=0$)

$$\frac{dN}{m_T dm_T} \propto \int_0^R r dr m_T I_0\left(\frac{p_T \sinh \rho}{T}\right) K_1\left(\frac{m_T \cosh \rho}{T}\right)$$

thermal + flow

$$\frac{dN}{m_T dm_T} \propto e^{-m_T/T}$$

simple

$$\frac{dN}{m_T dm_T} \propto \frac{e^{-m_T/T}}{m_T^\lambda}$$

Empirical parametrization from pp (m_T -scaling)

but also from theoretical model (flux-tube + Schwinger)
(Gatoff, Wong, PRD 46, 997 (1992))

Note: “T” depends on function used

in papers often more than one fit function quoted ...

Assumptions:

- Factorization between the hard part and the non perturbative PDF and fragmentation function $D_q \square H(z_q, Q^2)$
- Universal fragmentation and PDFs (e.g PDF from ep, fragmentation fz. from ee, but used in pp data)

Medium modifications: R_{AA} , R_{CP}

Nuclear modification factor for:

$$R_{AA}(p_t) = \frac{1}{\langle N_{coll} \rangle} \times \frac{dN_{AA} / dp_t}{dN_{pp} / dp_t}$$

- π^0, η , direct γ (not coming from neutral meson decays)
- **Charm. Heavy flavour electrons** (not photonic, or dielectron decays of mesons, or direct γ , or J/ψ and Y)

$$R_{CP}(p_t) = \frac{1}{\langle N_{coll} \rangle} \times \frac{dN_{AA} / dp_t / \langle T_{AA} \rangle [central]}{dN_{AA} / dp_t / \langle T_{AA} \rangle [peripheral]}$$

and status in ALICE...

