

Formulario

Medida. Probabilidad de obtener a al medir el observable A en un sistema que se encuentra en el estado $|\psi\rangle$:

$$P_\psi(a, A) = \sum_{i=1}^{g_a} |\langle \phi_{ai} | \psi \rangle|^2, \quad A|\phi_{ai}\rangle = a|\phi_{ai}\rangle, \quad i = 1, \dots, g_a, \quad g_a = \text{degeneración de } a.$$

Momento angular: $J_\pm = J_1 \pm iJ_2$,

$$\mathbf{J}^2|jm\rangle = \hbar^2 j(j+1)|jm\rangle, \quad J_3 = \hbar m|jm\rangle, \quad J_\pm|jm\rangle = \hbar \sqrt{j(j+1) - m(m \pm 1)} |j, m \pm 1\rangle.$$

Matrices de espín $s = \frac{1}{2}$. Para un vector unitario $\hat{\mathbf{a}}$ cualquiera,

$$\mathbf{S} = \frac{\hbar}{2} \boldsymbol{\sigma}, \quad \sigma_x := \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \quad \sigma_y := \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}, \quad \sigma_z := \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}, \quad e^{i\theta(\hat{\mathbf{a}} \cdot \boldsymbol{\sigma})} = \cos \theta \mathbf{1} + i \sin \theta (\hat{\mathbf{a}} \cdot \boldsymbol{\sigma}).$$

Autoestados de $\mathbf{S}_{\hat{\mathbf{n}}} = \mathbf{S} \cdot \hat{\mathbf{n}}$, con $\hat{\mathbf{n}} = (\sin \theta \cos \phi, \sin \theta \sin \phi, \cos \theta)$ dirección arbitraria en coordenadas esféricas:

$$\mathbf{S}_{\hat{\mathbf{n}}}|\phi_\pm\rangle = \pm \frac{\hbar}{2} |\phi_\pm\rangle \quad |\phi_+\rangle = \begin{pmatrix} \cos \frac{\theta}{2} \\ e^{i\phi} \sin \frac{\theta}{2} \end{pmatrix}, \quad |\phi_-\rangle = \begin{pmatrix} \sin \frac{\theta}{2} \\ -e^{i\phi} \cos \frac{\theta}{2} \end{pmatrix}, \quad \langle \phi_\pm | \mathbf{S}_{\hat{\mathbf{n}}} | \phi_\pm \rangle = \pm \frac{\hbar}{2} \hat{\mathbf{n}}.$$

Matrices de espín $s = 1$:

$$S_x := \frac{\hbar}{\sqrt{2}} \begin{pmatrix} 0 & 1 & 0 \\ 1 & 0 & 1 \\ 0 & 1 & 0 \end{pmatrix}, \quad S_y := \frac{\hbar}{\sqrt{2}} \begin{pmatrix} 0 & -i & 0 \\ i & 0 & -i \\ 0 & i & 0 \end{pmatrix}, \quad S_z := \hbar \begin{pmatrix} 1 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & -1 \end{pmatrix}.$$

Armónicos esféricos:

$$Y_0^0(\theta, \phi) = \sqrt{\frac{1}{4\pi}}; \quad Y_1^1(\theta, \phi) = -\sqrt{\frac{3}{8\pi}} e^{i\phi} \sin \theta, \quad Y_1^0(\theta, \phi) = \sqrt{\frac{3}{4\pi}} \cos \theta, \quad Y_1^{-1}(\theta, \phi) = \sqrt{\frac{3}{8\pi}} e^{-i\phi} \sin \theta.$$

Perturbaciones independientes del tiempo $H = H_0 + H_I$. **Niveles no degenerados** $H_0|\phi_n^{(0)}\rangle = E_n^{(0)}|\phi_n^{(0)}\rangle$:

$$E_n \simeq E_n^{(0)} + \langle \phi_n^{(0)} | H_I | \phi_n^{(0)} \rangle + \sum_{k \neq n} \frac{|\langle \phi_k^{(0)} | H_I | \phi_n^{(0)} \rangle|^2}{E_n^{(0)} - E_k^{(0)}}, \quad |\phi_n\rangle \simeq |\phi_n^{(0)}\rangle + \sum_{k \neq n} \frac{\langle \phi_k^{(0)} | H_I | \phi_n^{(0)} \rangle}{E_n^{(0)} - E_k^{(0)}} |\phi_k^{(0)}\rangle.$$

Niveles degenerados $H_0|\phi_{ni}^{(0)}\rangle = E_n^{(0)}|\phi_{ni}^{(0)}\rangle$: Las correcciones $E_n^{(1)}$ son los autovalores de la matriz $\mathbb{H} = (H_{I,ij})$, con $H_{I,ij} = \langle \phi_{ni}^{(0)} | H_I | \phi_{nj}^{(0)} \rangle$.

Perturbaciones dependientes del tiempo $H = H_0 + V(t)$:

$$P_{i \rightarrow f}(t_0, t) \simeq \frac{1}{\hbar^2} \left| \int_{t_0}^t dt' e^{i\omega_{fi}(t'-t_0)} V_{fi}(t') \right|^2, \quad \omega_{fi} := \frac{E_f^{(0)} - E_i^{(0)}}{\hbar} \quad V_{fi}(t') := \langle \phi_f^{(0)} | V(t') | \phi_i^{(0)} \rangle, \quad f \neq i$$

Método variacional: $E_{\text{fund}} \simeq E(\lambda_{\min})$, con $E(\lambda) = \frac{\langle \psi_\lambda | H | \psi_\lambda \rangle}{\langle \psi_\lambda | \psi_\lambda \rangle}$.

Pozo infinito unidimensional de anchura a centrado en el origen:

$$\phi_n(x) = \sqrt{\frac{2}{a}} \cos\left(\frac{n\pi x}{a}\right), \quad n = 1, 3, 5, \dots; \quad \phi_n(x) = \sqrt{\frac{2}{a}} \sin\left(\frac{n\pi x}{a}\right), \quad n = 2, 4, 6, \dots; \quad E_n = \frac{\hbar^2 \pi^2}{2ma^2} n^2.$$

Oscilador armónico unidimensional:

$$E_n = \hbar \omega \left(n + \frac{1}{2}\right), \quad a|n\rangle = \sqrt{n}|n-1\rangle \quad a^+|n\rangle = \sqrt{n+1}|n+1\rangle, \quad [a, a^+] = 1,$$

$$a = \frac{1}{\sqrt{2}} \left(\alpha x + \frac{i p}{\alpha \hbar}\right), \quad a^+ = \frac{1}{\sqrt{2}} \left(\alpha x - \frac{i p}{\alpha \hbar}\right); \quad x = \frac{1}{\sqrt{2}\alpha} (a + a^+), \quad p = \frac{i\hbar\alpha}{\sqrt{2}} (a^+ - a); \quad \alpha = \sqrt{\frac{m\omega}{\hbar}}.$$

Atomo de hidrógeno: $\psi_{n\ell m_\ell} = R_{n\ell}(r) Y_\ell^{m_\ell}(\theta, \phi)$, $E_n = -\frac{Z^2 e^2}{(4\pi\epsilon_0) 2a_0 n^2}$, $4\pi\epsilon_0 = 1$ en unidades gaussianas.

Integrales útiles: $\int_{-\infty}^{\infty} dx e^{-ax^2 + ibx} = \sqrt{\frac{\pi}{a}} e^{-b^2/4a}, \quad a > 0, b \text{ reales.}$

43. CLEBSCH-GORDAN COEFFICIENTS, SPHERICAL HARMONICS, AND d FUNCTIONS

Note: A square-root sign is to be understood over *every* coefficient, e.g., for $-8/15$ read $-\sqrt{8/15}$.

Notation:

J	J	...
M	M	...
m_1	m_2	
m_1	m_2	Coefficients
$\vdots$	$\vdots$	
$\vdots$	$\vdots$	

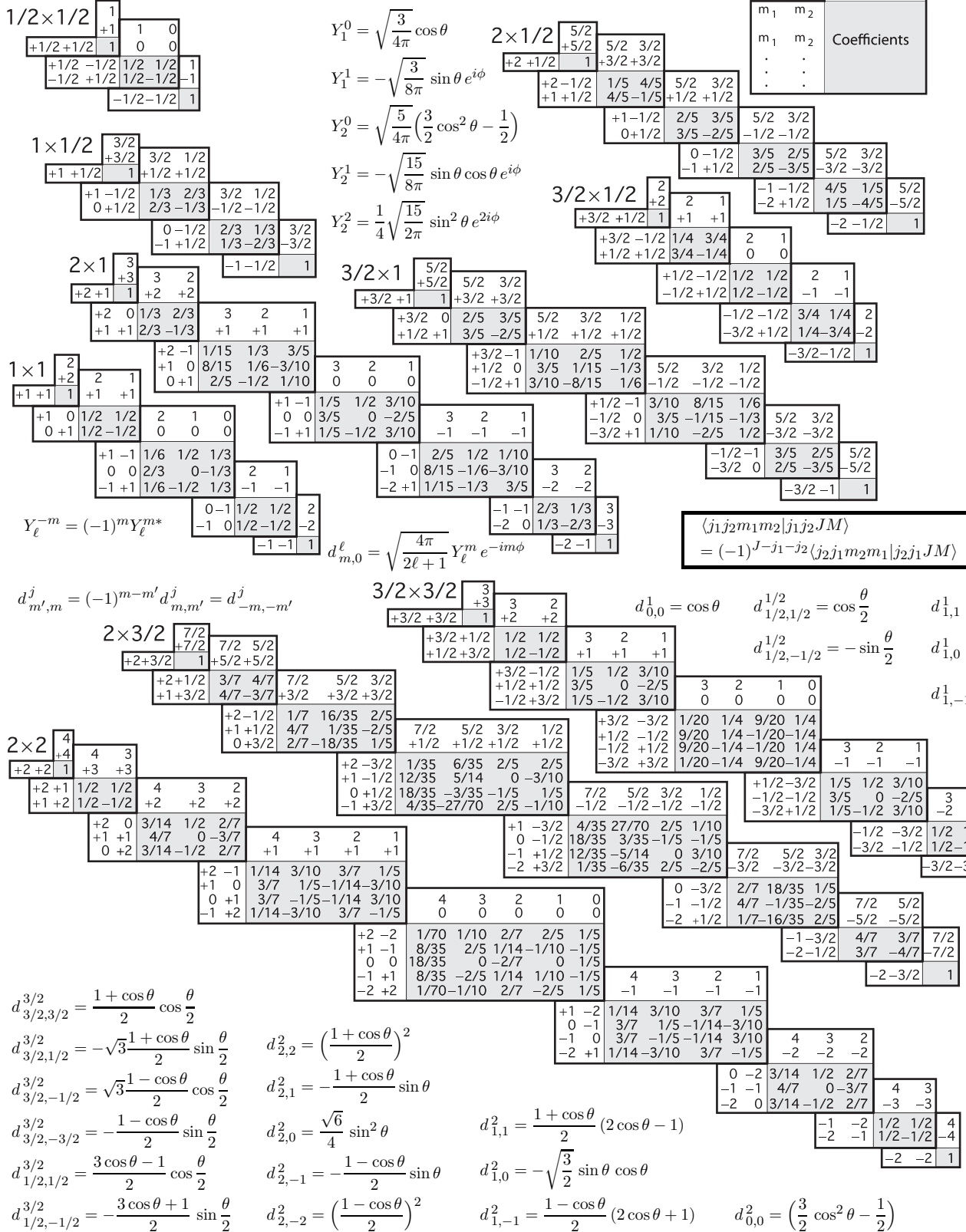


Figure 43.1: The sign convention is that of Wigner (*Group Theory*, Academic Press, New York, 1959), also used by Condon and Shortley (*The Theory of Atomic Spectra*, Cambridge Univ. Press, New York, 1953), Rose (*Elementary Theory of Angular Momentum*, Wiley, New York, 1957), and Cohen (*Tables of the Clebsch-Gordan Coefficients*, North American Rockwell Science Center, Thousand Oaks, Calif., 1974).