# DBI Galileon inflation in the light of Planck 2015 (IberiCos 2015)

## K. Sravan Kumar<sup>(1)</sup>, Juan C. Bueno Sánchez, Celia Escamilla Rivera, João Marto and Paulo Vargas Moniz.

<sup>(1)</sup>Departamento de Física, Universidade da Beira Interior, 6200 Covilhã, Portugal



PROGRAMAS DE DOUTORAMENTO FCT April 1, 2015



(日)

▲□▶ ▲□▶ ▲三▶ ▲三▶ 三三 のへで

## Outlines

- Inflation and Planck 2015
- Status of DBI inflation
- Motivations for DBI Galileon inflation
- Discussion on various limits of DBIG
- Results

## Constraints on inflation

## Spectral index Vs tensor scalar ratio

The current bounds  $n_s = 0.968 \pm 0.006$  and r < 0.09 at 95% CL constrain many inflationary scenarios.



Figure :  $n_s - r$  Planck

- The Planck results indicate that we live in a spatially flat universe,  $\Omega_k = 0.000 \pm 0.0025$ ,
- The perturbation spectrum imprinted in the CMB is Gaussian to a high degree, imposing severe constraints on the bispectrum amplitudes:  $f_{NL}^{loc} = 0.8 \pm 5.0$ ,  $f_{NL}^{eq} = -4 \pm 43$  and  $f_{NL}^{ortho} = -26 \pm 21$  at the 68% confidence level.
- Small running  $dn_s/d \ln k = -0.003 \pm 0.007$
- All these so far not evidently deviated from single field inflation.
- The fact that Starobinsky and Higgs inflationary models are consistent with data indicates the micro-physics of inflation should therefore be described by an EFT.

The new data gives us a wonderful opportunity to construct and validate viable and generic inflatioinary models from UV-completion theories such as SUGRA and string theory.

## Inflationary model building

- The prediction of starobinsky model with r = 0.0033 is well consistent with the data and it has become a sweet spot of present day inflationary cosmology.
- It is important to build models with  $r \sim \mathcal{O}(10^{-3})$  maintaining  $f_{NL} \sim \mathcal{O}(1)$  as the future experiments as the post-Planck sattellites CMBPol,COrE, Prism and LiteBIRD and many other ground based experiments such as Keck/BICEP3 etc are expected to reach sensitivity to detect  $r \sim 2 \times 10^{-3}$  Creminelli et al (2015).
- $r \sim \mathcal{O}(10^{-3})$  maintaining  $f_{NL} \sim \mathcal{O}(1)$  is rather a difficult task for non-canonical models, perhaps models built with features of modified gravity, string theory and SUGRA might be best candidates ,Jerome Martin (2015).
- In the context of inflationary models the commonly heard word is 'generic 'which should mean how well the model consistent with UV-completion theory rather than how small is the parameter space.

#### Status of tensor tilt

- Tensor tilt is one of the crucial parameter to validate inflation. Especially a detection of red tilt can confirm standard inflationary paradigm. Alternatives to inflation such as Ekpyrotic and String gas cosmology predict a blue tensor tilt.
- Although we have not yet detected B-modes, the phenomenological and statistical evidence of BKP, BKP+LIGO find no evidence for blue tensor tilt. Huang and Wang (2015)

parameter	BKP		BKP+LIGO	
	68% C.L.	95% C.L.	68% C.L.	95% C.L.
r	< 0.055	< 0.099	< 0.059	< 0.106
$n_t$	$0.66^{+1.83}_{-1.44}$	$0.66^{+2.92}_{-3.42}$	$-0.76^{+1.37}_{-0.52}$	$-0.76^{+1.63}_{-2.21}$

Figure : The 68% and 95% limits for the parameters r and  $n_t$  from the BKP only and BKP+LIGO datasets.

#### Consistency relation

We need to detect GW's to further screen out and validate inflationary models. A significant detection of tensor tilt  $n_t$  in post Planck experiments could play a vital role in inflationary cosmology.



Figure : The marginalized contour plot and likelihood distributions of the parameters r and  $n_t$  from the BKP only (red) and BKP+LIGO (blue) datasets.

## **DBI** inflation

- DBI inflation has been an attractive model in String cosmology , Silverstein and Tong (2004).
- DBI is a natural non-canonical model from String theory. And it has the potential for smaller tensor scalar ratio and large Non-gaussianity

$$\begin{array}{c} r = 16\epsilon c_s \\ \downarrow \qquad \downarrow \qquad f_{NL} \sim \frac{1}{c_s^2} \end{array}$$



Figure : A Calabi–Yau space with a warped throat region. The dot on the left denotes a D3-brane sitting at the tip of the throat. The dot to the right with the arrow is a mobile D3-brane in the throat. It is attracted by the D3-brane.

#### DBI->DBI Galileon

- DBI inflation is consistent with the current planck data although speed of sound is constrained to be  $c_s > 0.087$ .
- Due to NG constraints, the lower bound on tensor scalar ratio in DBI model is given by r > 0.01, Baumann and McAllister (2006).
- de Rham and Tolley (2010) proposed a very interesting extension of the DBI action which generalizes and unifies single field DBI and the modified gravity model called the Galileon (DBIG).
- Multifield DBI Galileon inflation has been studied in literature by its potential to generate large Orthogonal NG.
   S.Renaux-Petel et al (2011), Koyama et al (2014).
- DBI Galileon model has also been embedded in SUGRA , Sayantan choudhury et al (2013).

## DBIG model

$$S = \int d^4x \left[ rac{m_{
ho}^2}{2} \sqrt{-g} R\left[g
ight] + rac{ ilde{m}^2}{2} \sqrt{-\gamma} R\left[\gamma
ight] + \sqrt{-g} \mathcal{L}_{\textit{brane}} 
ight]$$

 $Einstein-Hilbert + Induced \ gravity + DBI \ action$ 

where  $m_p$  is reduced planck mass  $(m_p = 2.24 \times 10^{18} {
m Gev})$  and  $\tilde{m}$  is a parameter associated with induced gravity. And

$$\mathcal{L}_{\textit{brane}} = -rac{1}{f\left(arphi
ight)}\left(\sqrt{\mathcal{D}}-1
ight) - V\left(arphi
ight)$$

Where

$$\mathcal{D} \equiv \det \left( \delta^{\mu}_{\nu} + f \partial_{\mu} \varphi \partial_{\nu} \varphi \right)$$

The induced metric is

$$\gamma_{\mu\nu} = f^{-1/2} \left( g_{\mu\nu} + f \partial_{\mu} \varphi \partial_{\nu} \varphi \right)$$

From now on f and  $\varphi$  are rescaled warp factor and scalar field,

$$f = \frac{h}{T_3}, \quad \varphi = \sqrt{T_3}\rho \tag{1}$$

▲ロト ▲園ト ▲ヨト ▲ヨト ニヨー わえの

## DBIG model

The modified Friedmann equations for the action are given by

$$3H^{2}\left(m_{P}^{2}+\frac{\tilde{m}^{2}}{c_{D}^{3}}\left(1-\frac{\epsilon_{f}}{2}-\frac{\epsilon_{f}^{2}}{16}\right)\right)=\frac{1}{f}\left(\frac{1}{c_{D}}-1\right)+V$$
$$-m_{P}^{2}\dot{H}+\frac{\tilde{m}^{2}H^{2}}{c_{D}}\left[\epsilon+g_{1}\left(\epsilon_{D},\epsilon_{f},\eta_{f}\right)+\frac{3}{2}\left(\frac{1}{c_{D}^{2}}-1\right)g_{2}\left(\epsilon_{f}\right)\right]=\frac{\dot{\varphi}^{2}}{2c_{D}},$$
$$c_{D}^{2}\equiv1-f\dot{\varphi}^{2},$$

Where the slow roll parameters are

$$\epsilon \equiv -\frac{\dot{H}}{H^2} \quad , \quad \eta \equiv \frac{d \ln \epsilon}{d \ln a} \quad , \quad \epsilon_{\mathcal{D}} \equiv \frac{d \ln c_{\mathcal{D}}}{d \ln a}$$
$$\eta_{\mathcal{D}} \equiv \frac{d \ln \epsilon_{\mathcal{D}}}{d \ln a} \quad \epsilon_f \equiv \frac{d \ln f}{d \ln a} \quad , \quad \eta_f \equiv \frac{d \ln \epsilon_f}{d \ln a}.$$

・ロト・日本・日本・日本・日本・日本

## Studying DBIG in various limits

- DBIG  $\longrightarrow$  DBI as  $\tilde{m} \rightarrow 0$ : In this limit we revisit the constraints on parameter space of DBI inflation.
- DBIG → Galileon as m<sub>p</sub> → 0 (i.e, switching off Einstein gravity): Taking this limit is not generic but it help us to understand the effects of induced gravity.
- Intermediate case: This case contains both the features of induced gravity and DBI.  $c_D$ , f = constant
- DBIG with constantly varying speed of sound  $\eta_{\mathcal{D}} = 0$
- DBIG with constantly varying warp factor  $\eta_f = 0$ : It is interesting to study motion of D3 brane in the varying warped geometries.
- Studying the model with varying warp factor and varying speed of sound is theoretically more generic but it is complicated. We leave this case for future studies.

## Approach

- The parameter space of DBIG inflation is (c<sub>D</sub>, m̃, f) and the free parameters are (ε<sub>D</sub>, η<sub>D</sub>, ε<sub>f</sub>, η<sub>f</sub>).
- We numerically constrain the parameter space using background equations wrt the latest CMB bounds on  $n_s = 0.968 \pm 0.006$  and r < 0.09, and power spectrum at pivot scale  $\mathcal{P}_{\zeta_*} = 2.2 \times 10^{-9}$ .
- Within the constrained parameter space  $(c_{D*}, \tilde{m}, f_*)$  we compute the tensor tilt  $n_t$  to all orders in slowroll parameters. Therefore these values of  $n_t$  can serve as the test of this model for future experiments.
- From the numerical results we predict No blue tensor tilt in this model but the consistency relation is violated from standard one  $r \simeq -8n_t$ .
- We also compute the energy scale of inflation  $\Lambda$  and mass of inflation  $m_{\varphi}$ .

#### Power spectra

- Perturbations in most general single field (Generalized G-inflation) inflationary models has been studied by Kobayashi et al (2011).
- We compute all the inflationary parameters to the full order in slow roll parameters.

The power spectrum of scalar perturbations is

$$\mathcal{P}_{\zeta} = \frac{\gamma_s}{2} \frac{\mathcal{G}_{s*}^{1/2}}{\mathcal{F}_{s*}^{3/2}} \frac{H_*^2}{4\pi^2} \quad , \quad \gamma_s \equiv 2^{2\nu_s - 3} \frac{\Gamma(\nu_s)^2}{\Gamma(3/2)^2} \left(1 - \epsilon_* + \frac{g_{s*}}{2} - \frac{f_{s*}}{2}\right)^2$$

The spectral index  $(n_s)$  is given by

$$n_s - 1 = 3 - 2\nu_s$$

For DBIG  $\mathcal{F}_s \equiv \mathcal{F}_s(\epsilon, c_{\mathcal{D}}, \tilde{m}, \epsilon_{\mathcal{D}})$ ,  $\mathcal{G}_s \equiv \mathcal{G}_s(\epsilon, c_{\mathcal{D}}, \tilde{m}, \epsilon_{\mathcal{D}})$  and  $f_s \equiv \frac{d \ln \mathcal{G}_s}{d \ln a}$ ,  $g_s \equiv \frac{d \ln \mathcal{G}_s}{d \ln a}$ .

#### Power spectra

The power spectrum of tensor perturbations is,

$$\mathcal{P}_{t} = 8\gamma_{t} \frac{\mathcal{G}_{t*}^{1/2}}{\mathcal{F}_{t*}^{3/2}} \frac{\mathcal{H}_{*}^{2}}{4\pi^{2}} \quad , \quad \gamma_{t} \equiv 2^{2\nu_{t}-3} \frac{\Gamma(\nu_{t})^{2}}{\Gamma(3/2)^{2}} \left(1 - \epsilon_{*} + \frac{g_{t*}}{2} - \frac{f_{t*}}{2}\right)^{2}$$

The tensor tilt  $(n_t)$  is given by

$$n_t = 3 - 2\nu_t$$

Where  $\mathcal{F}_t \equiv \mathcal{F}_t(c_{\mathcal{D}}, \tilde{m})$ ,  $\mathcal{G}_t \equiv \mathcal{G}_t(c_{\mathcal{D}}, \tilde{m})$  and  $f_t \equiv \frac{d \ln \mathcal{F}_t}{d \ln a}$   $g_t \equiv \frac{d \ln \mathcal{G}_t}{d \ln a}$ . The tensor scalar ratio is

$$r \equiv \frac{\mathcal{P}_t}{\mathcal{P}_{\zeta}} = 16 \frac{\gamma_t}{\gamma_s} \left(\frac{\mathcal{G}_t}{\mathcal{F}_t}\right)^{1/2} \left(\frac{\mathcal{F}_s}{\mathcal{F}_t}\right)^{3/2}$$

◆□ > ◆□ > ◆豆 > ◆豆 > ̄豆 \_ のへで

▲□▶ ▲□▶ ▲ 三▶ ▲ 三▶ 三三 - のへぐ

## Results

Inflation	r	n <sub>t</sub>	$m_{\varphi}$ (TeV)	$f/m_p^4$
DBI limit	(0.01, 0.1)	(-0.049, -0.038)	$1.46\times10^9$	$(0.07, 0.5) \times 10^{12}$
Galileon limit	(0.13, 0.15)	(-0.05, -0.03)	$m_{arphi}^2 < 0$	$\sim 10^8$
Intermediate	(0.0095, 0.06)	(-0.07, -0.04)	$m_{arphi}^2 < 0$	$(0.2, 1.1)  imes 10^9$
Varying $c_{\mathcal{D}}$	(0.023, 0.033)	-0.05	$m_{arphi}^2 < 0$	$(1.2, 1.5)  imes 10^9$
Varying f	(0.0037, 0.0087)	(-0.068, -0.027)	$(2,8.2)\times10^7$	$(0.54, 1.37) \times 10^7$

Table : Summary of inflationary observables in single field DBI galileon model for No. of efoldings  $N_e = 60$ .

## Conclusion



Figure : Predictions of DBIG  $(n_s - r)$  for constant warp factor (dotted line) and varying warp factor cases (full line).

#### Conclusions

- In all the cases, we restricted the speed of sound  $c_D \lesssim 1$  to avoid large Non-gaussianities. Therefore we attain inflation with  $r \sim \mathcal{O}(10^{-3})$  with  $f_{NL} \sim \mathcal{O}(1)$ .
- We predict tensor tilt of the model  $n_t$  and violation of standard consistency relation  $(r \le -8n_t)$ . This can serve as an experimental test of the model if future experiments detect gravitational waves.
- We obtain the constraints on warp factor to have the model observationally viable. The values of warp factor might be important in understanding warped geometries and also in string phenomenology.