

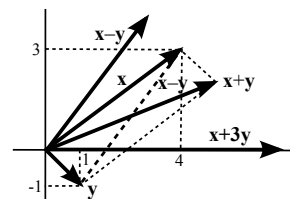
Soluciones de problemas 1 de MM(im) (2020)

1. $\mathbf{x} = (4, 3)$, $\mathbf{y} = (1, -1)$. $\mathbf{x} + \mathbf{y} = (5, 2)$, $\mathbf{x} - \mathbf{y} = (3, 4)$ y $\mathbf{x} + 3\mathbf{y} = (7, 0)$,

$$\|\mathbf{x} + \mathbf{y}\| = \sqrt{29} \leq 6 \leq 5 + \sqrt{2} = \|\mathbf{x}\| + \|\mathbf{y}\|. \quad \|\mathbf{x} - \mathbf{y}\| = 5 = \text{distancia de } \mathbf{x} \text{ a } \mathbf{y}.$$

$$|\mathbf{x} \cdot \mathbf{y}| = 1 \leq 5\sqrt{2} = \|\mathbf{x}\|\|\mathbf{y}\|. \quad \mathbf{x} \cdot \mathbf{y} > 0 \Rightarrow \text{ángulo agudo. } \|\mathbf{x} + 3\mathbf{y}\| = 7.$$

$$\frac{(\mathbf{x} + 3\mathbf{y}) \cdot \mathbf{y}}{\|\mathbf{x} + 3\mathbf{y}\|\|\mathbf{y}\|} = \frac{7}{7\sqrt{2}} = \cos \phi \rightarrow \mathbf{y} \text{ y } \mathbf{x} + 3\mathbf{y} \text{ forman un ángulo } \frac{\pi}{4} \text{ (claro en dibujo).}$$



2. $\mathbf{b} - \mathbf{a} = (0, 1, 1)$, $\mathbf{a} \cdot \mathbf{b} = 1$, $\mathbf{a} \times \mathbf{b} = \begin{vmatrix} \mathbf{i} & \mathbf{j} & \mathbf{k} \\ 1 & 0 & -1 \\ 1 & 1 & 0 \end{vmatrix} = (1, -1, 1)$, $\mathbf{a} \cdot (\mathbf{a} \times \mathbf{b}) = 0$. $\|\mathbf{b} - \mathbf{a}\| = \sqrt{2} = \text{distancia de } \mathbf{a} \text{ a } \mathbf{b}$.
perpendiculares

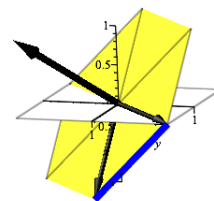
$$\frac{\mathbf{a} \cdot \mathbf{b}}{\|\mathbf{a}\|\|\mathbf{b}\|} = \frac{1}{2} = \cos \phi \rightarrow \mathbf{a} \text{ y } \mathbf{b} \text{ forman un ángulo } \frac{\pi}{3}. \quad \|\mathbf{a} + \mathbf{b}\| = \|(2, 1, -1)\| = \sqrt{6} \leq 2\sqrt{2} = \|\mathbf{a}\| + \|\mathbf{b}\| \quad (6 < 8).$$

Plano que contiene a $\mathbf{a}$ y $\mathbf{b}$: $\mathbf{x} = t(1, 0, -1) + s(1, 1, 0)$, $\begin{cases} x = t+s \\ y = s \\ z = -t \end{cases} \rightarrow \boxed{z = y - x}$.

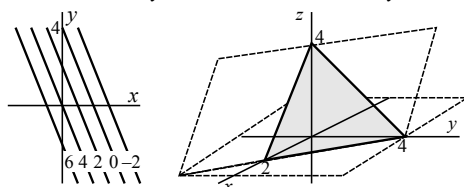
O como $\mathbf{a} \times \mathbf{b}$ es perpendicular a este plano: $(1, -1, 1) \cdot (x, y, z) = 0$ ✓

Segmentos $\mathbf{c}(t) = \mathbf{a} + t(\mathbf{b} - \mathbf{a})$ y $\mathbf{c}_*(t) = \mathbf{b} + t(\mathbf{a} - \mathbf{b})$, ambos con $t \in [0, 1]$. Es decir:

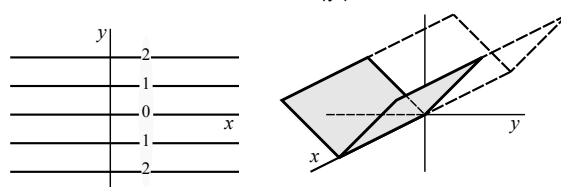
$$\mathbf{c}(t) = (1, 0, -1) + t(0, 1, 1) = \boxed{(1, t, t-1)}, \quad \mathbf{c}_*(t) = (1, 1, 0) + t(0, -1, -1) = \boxed{(1, 1-t, -t)}$$



3. a) $\boxed{z = 4 - 2x - y}$ $x=0 \rightarrow z=4-y$ curvas de nivel: $y=4-C-2x$
 $y=0 \rightarrow z=4-2x$



- b) $\boxed{z = |y|}$ corte con todo plano $x = cte \rightarrow z = |y|$
curvas de nivel: $|y| = C$

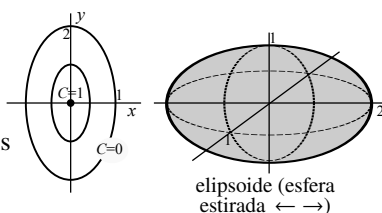


- c) $\boxed{4x^2 + y^2 + 4z^2 = 4}$

$$z = C \rightarrow 4x^2 + y^2 = 4 - 4C^2$$

$$x = 0 \rightarrow \frac{y^2}{2^2} + z^2 = 1 \leftarrow \text{elipses}$$

$$y = 0 \rightarrow x^2 + z^2 = 1$$

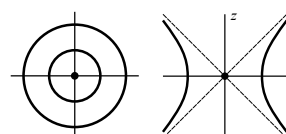


- d) $\boxed{z^2 = x^2 + y^2 - 1}$ de revolución

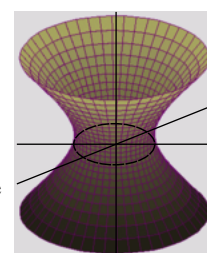
$$x = 0 \rightarrow y^2 - z^2 = 1$$

$$y = 0 \rightarrow x^2 - z^2 = 1$$

hipérbolas



hiperboloide de una hoja (wikipedia)



4. a) $\boxed{f(x, y) = xy}$ $\nabla f = (y, x) = (1, 0) (0, 1) (1, 1) (2, 2) (1, -1) (-1, 1) (-1, -1)$
en $(0, 1) (1, 0) (1, 1) (2, 2) (-1, 1) (1, -1) (-1, -1)$

[silla de montar como la del Ej 8]

Plano tangente: $\boxed{z = x + y - 1}$.

- b) $\boxed{g(x, y) = y^2}$ $\nabla g = (0, 2y) [= (0, 2) \text{ en } (1, 1)]$.

Plano tangente: $z = 1 + 2(y - 1)$, $\boxed{z = 2y - 1}$.

- c) $\boxed{h(x, y) = 2x - y}$ $\nabla h = (2, -1)$ (constante y $\perp$ a las rectas)

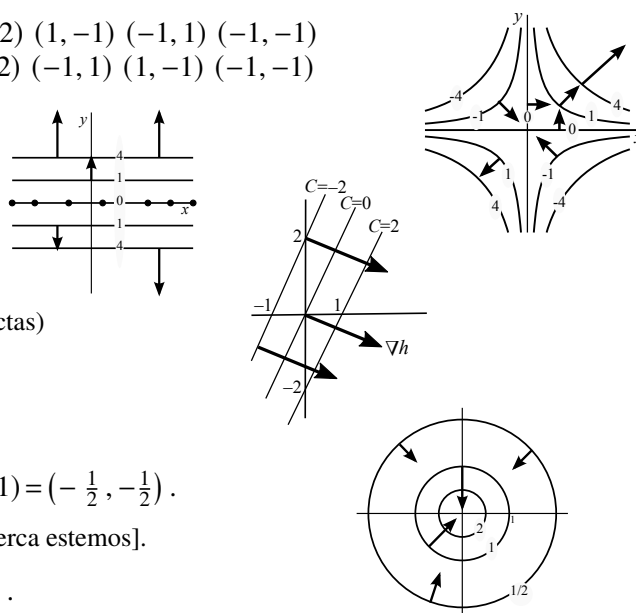
Plano tangente: $z = 1 + 2(x - 1) - (y - 1)$, $\boxed{z = 2x - y}$.

(claro, debía ser el propio plano)

- d) $\boxed{k(x, y) = (x^2 + y^2)^{-1}}$ $\nabla k = -\frac{2}{(x^2 + y^2)^2} (x, y)$. $\nabla k(1, 1) = (-\frac{1}{2}, -\frac{1}{2})$.

[∇k apunta hacia el origen y es mayor cuanto más cerca estemos].

Plano tangente: $z = \frac{1}{2} - \frac{1}{2}(x - 1) - \frac{1}{2}(y - 1)$, $\boxed{z = \frac{3 - x - y}{2}}$.



5. $f(x, y) = x^5 y - 4x - xy^3$ $f_x = 5x^4 y - 4 - y^3$, $f_y = x^5 - 3xy^2$. $f_{xx} = 20x^3 y$, $f_{xy} = f_{yx} = 5x^4 - 3y^4$, $f_{yy} = -6xy$.
 $g(x, y) = x e^{x+y}$ $g_x = (x+1) e^{x+y}$, $g_y = x e^{x+y}$. $g_{xx} = (x+2) e^{x+y}$, $g_{xy} = g_{yx} = (x+1) e^{x+y}$, $g_{yy} = x e^{x+y}$.
 $h(x, y) = \frac{\cos(xy)}{x}$ $h_x = -\frac{1}{x^2} \cos(xy) - \frac{y}{x} \sin(xy)$, $h_y = -\sin(xy)$.
 $h_{xx} = \left[\frac{2}{x^3} - \frac{y^2}{x} \right] \cos(xy) + \frac{2y}{x^2} \sin(xy)$, $h_{xy} = h_{yx} = -y \cos(xy)$, $h_{yy} = -x \cos(xy)$.
 $k(x, y) = \arctan \frac{x}{y}$ $k_x = \frac{y}{x^2+y^2}$, $k_y = -\frac{x}{x^2+y^2}$. $k_{xx} = -\frac{2xy}{(x^2+y^2)^2}$, $k_{xy} = k_{yx} = \frac{x^2-y^2}{(x^2+y^2)^2}$, $k_{yy} = \frac{2xy}{(x^2+y^2)^2}$.

6. $F(x, y, z) = x^3 - 2yz + e^x z^2$. $\nabla F = (3x^2 + z^2 e^x, -2z, 2ze^x - 2y) \xrightarrow{(0,1,2)} 2(2, -2, 1)$.
 Plano: $(2, -2, 1) \cdot (x, y-1, z-2) = 0$. $z = 2y - 2x$.

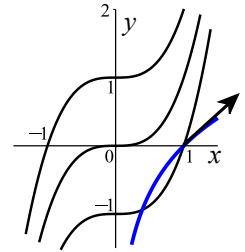
Perpendicular al vector y al gradiente en el punto es $(1, 0, 1) \times (2, -2, 1) = (2, 1, -2) \rightarrow \bar{u} = \left(\frac{2}{3}, \frac{1}{3}, -\frac{2}{3} \right)$ [o $-\bar{u}$].

[O resolviendo $\begin{cases} a+c=0 \\ 2a-2b+c=0 \end{cases} \rightarrow \begin{cases} c=-a \\ a=2b \end{cases}$, vector de la forma $(2b, b, -2b)$ de módulo $9|b|$].

7. $f(x, y) = (x^3 - y)^2$ y $\mathbf{c}(t) = (t, \ln t)$. $f = 0, 1 \rightarrow y = x^3$, $y = x^3 \pm 1$.

$\nabla f = (6x^2(x^3 - y), 2(y - x^3))$, $\nabla f(1, 0) = (6, -2)$,

$\mathbf{c}'(t) = (1, \frac{1}{t}) \xrightarrow{t=1} (1, 1)$, de módulo $\|\cdot\| = \sqrt{2}$. $D_{\mathbf{u}} f(1, 0) = (6, -2) \cdot \left(\frac{1}{\sqrt{2}}, \frac{1}{\sqrt{2}} \right) = 2\sqrt{2}$.



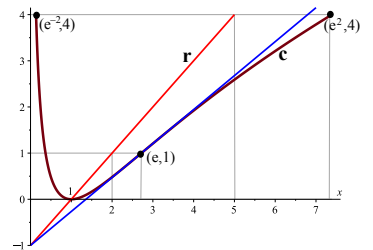
8. $\mathbf{c}(t) = (e^t, t^2)$, $t \in [-2, 2]$. Pasa por los puntos $(e^{-2}, 4)$, $(1, 0)$, $(e, 1)$ y $(e^2, 4)$.
 [La curva es la gráfica de $y = (\log x)^2$, pues $t = \log x$ e $y = t^2$].

i) $\mathbf{c}'(t) = (e^t, 2t)$, $\mathbf{c}'(1) = (e, 2)$. Tangente $\mathbf{x} = (e, 1) + t(e, 2) = (e + te, 1 + 2t)$.

En cartesianas: $t = \frac{x}{e} - 1 = \frac{y-1}{2}$, $y = \frac{2}{e}x - 1$. Normales $(\pm 1, \mp e) \rightarrow \frac{1}{\sqrt{e^2+1}}(\pm 1, \mp e)$.

ii) La recta $\mathbf{r}(s) = (s, s-1)$ corta la curva sólo en $(1, 0)$ ($t=0$ y $s=1$).

$[s = e^t, s-1 = t^2 \Rightarrow 1+t^2 = e^t \Leftrightarrow t=0, \text{ como muestran las gráficas}]$.



9. $g(x, y) = \sqrt{20 - x^2 - y^2}$ $\mathbf{c}(t) = (t, t^2)$

a) $g(x, y) = C \rightarrow x^2 + y^2 = 20 - C^2$. Circunferencias para $0 \leq C < 2\sqrt{5}$.

Para $C = 2\sqrt{5}$ es el punto $(0, 0)$. C es parte de la parábola $y = x^2$.

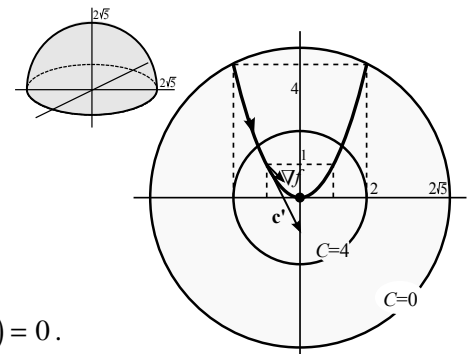
b) $g_x = \frac{-x}{\sqrt{20 - x^2 - y^2}}$, $g_y = \frac{-y}{\sqrt{20 - x^2 - y^2}}$. $\nabla g(-1, 1) = \left(\frac{1}{3\sqrt{2}}, -\frac{1}{3\sqrt{2}} \right)$.

Plano tangente: $z = 3\sqrt{2} + \frac{1}{3\sqrt{2}}(x+1) - \frac{1}{3\sqrt{2}}(y-1)$, $z = \frac{1}{3\sqrt{2}}[20 + x - y]$.

O bien: $G(x, y, z) = x^2 + y^2 + z^2 = 20 \rightarrow (2, 2, 6\sqrt{2}) \cdot (x+1, y-1, z-3\sqrt{2}) = 0$.

c) $\mathbf{c}'(-1) = (1, 2t)|_{t=-1} = (1, -2)$. $h'(t) = g_x(x(t), y(t))x'(t) + g_y(x(t), y(t))y'(t) = \nabla g(\mathbf{c}(t)) \cdot \mathbf{c}'(t)$.

$h'(-1) = \nabla g(-1, 1) \cdot (1, -2) = \frac{1}{\sqrt{2}}$. [Componiendo y derivando: $h(t) = \sqrt{20 - t^2 - t^4}$, $h'(t) = \frac{-t - 2t^3}{\sqrt{20 - t^2 - t^4}} \xrightarrow{t=-1} \frac{1}{\sqrt{2}}$].



10. $h(t) = f(t, -t, t^2)$ $h'(t) = f_x - f_y + 2tf_z$, $h''(t) = (f_x)' - (f_y)' + 2t(f_z)' + 2f_z$
 $= f_{xx} - f_{xy} + 2tf_{xz} - f_{yx} + f_{yy} - 2tf_{yz} + 2tf_{zx} - 2tf_{zy} + 4t^2 f_{zz} + 2f_z$
 $= f_{xx} + f_{yy} + 4t^2 f_{zz} - 2f_{xy} + 4tf_{xz} - 4tf_{yz} + 2f_z$

En concreto, para la f dada es $h(t) = t^4$, $h''(t) = 12t^2$. Con la fórmula: $h''(t) = 8t^2 + 4z|_{z=t^2} = 12t^2$.

11. $\mathbf{g}(x, y, z) = (2x^2 - y + 3z^3, 2y - x^2)$ $\mathbf{Dg} = \begin{pmatrix} 4x & -1 & 9z^2 \\ -2x & 2 & 0 \end{pmatrix} \Rightarrow \mathbf{Dg}(2, -1, 1) = \begin{pmatrix} 8 & -1 & 9 \\ -4 & 2 & 0 \end{pmatrix}$.

$\mathbf{Df} \mathbf{Dg} = \begin{pmatrix} 0 & 3 & 9 \\ 12 & 0 & 18 \end{pmatrix}$.

$\mathbf{f}(u, v) = (e^{u+2v}, 2u+v)$ $\mathbf{Df} = \begin{pmatrix} e^{u+2v} & 2e^{u+2v} \\ 2 & 1 \end{pmatrix}$, $\mathbf{g}(2, -1, 1) = (12, -6) \Rightarrow \mathbf{Df}(12, -6) = \begin{pmatrix} 12 & 2 \\ 2 & 1 \end{pmatrix}$.

$(\mathbf{f} \circ \mathbf{g})(x, y, z) = (e^{3y+3z^3}, 3x^2+6z^3)$, $\mathbf{D}(\mathbf{f} \circ \mathbf{g}) = \begin{pmatrix} 0 & 3e^{-} & 9z^2e^{-} \\ 6x & 0 & 18z^2 \end{pmatrix}$, que en $(2, -1, 1) \uparrow$

12. $f(x, y) = \frac{x+y}{1+xy}$ $f_x = \frac{1-y^2}{(1+xy)^2}$, $f_y = \frac{1-x^2}{(1+xy)^2}$, $\nabla f(0, 2) = (-3, 1)$. Plano tangente: $z = 2 - 3x + y - 2 = y - 3x$.

Recta tangente, perpendicular al gradiente: $(x, y-2) \cdot (-3, 1) = 0 \rightarrow y = 3x+2$.

O directamente: $\frac{x+y}{1+xy} = 2 \rightarrow y = \frac{2-x}{1-2x} \rightarrow y'(0) = \frac{3}{(1-2x)^2} \Big|_{x=0} = 3 \rightarrow y = 2+3x$.

O derivando implícitamente: $\frac{(1+y')(1+xy)-(x+y)(y+xy')}{(1+xy)^2} \Big|_{x=2, y=0} = 1+y'-4 = 0 \rightarrow y'(0) = 3 \dots$

Si $h(u, v) = f(u^3 + v^2 - 1, e^v + 1)$, es $Dh(1, 0) = \nabla h(1, 0) = (\nabla f(0, 2)) (Dg(1, 0)) = (-3, 1) \begin{pmatrix} 3 \cdot 1^2 & 2 \cdot 0 \\ 0 & e^0 \end{pmatrix} = (-9, 1) \Rightarrow$

$$D_{(1/\sqrt{2}, -1/\sqrt{2})} h(1, 0) = (-9, 1) \cdot \left(\frac{1}{\sqrt{2}}, -\frac{1}{\sqrt{2}} \right) = -5\sqrt{2}.$$

13. $f(x, y) = \frac{x}{x^2+y^2}$ $C=0 \rightarrow x=0$. $\frac{x}{x^2+y^2} = C \rightarrow x^2+y^2 = \frac{x}{C}$, $(x - \frac{1}{2C})^2 + y^2 = \frac{1}{4C^2}$ si $C \neq 0$,
circunferencias de centro $(\frac{1}{2C}, 0)$ que pasan por $(0, 0)$ y $(\frac{1}{C}, 0)$.

$$\nabla f(x, y) = \left(\frac{y^2 - x^2}{(x^2+y^2)^2}, \frac{-2xy}{(x^2+y^2)^2} \right), \nabla f(2, 1) = \left(-\frac{3}{25}, -\frac{4}{25} \right) \text{ (}\perp\text{ a la circunferencia por el punto).}$$

i) La derivada direccional es 0 en dirección perpendicular a ∇f : $\mathbf{u} = \left(\pm \frac{4}{5}, \mp \frac{3}{5} \right)$.

ii) Uno claro es $\mathbf{u} = (1, 0)$, pues la parcial $f_x = -\frac{3}{25}$ es precisamente la derivada según **i**.

Con menos vista: $D_{\mathbf{u}} f(2, 1) = \nabla \cdot \mathbf{u} = \left(-\frac{3}{25}, -\frac{4}{25} \right) \cdot (a, b) = -\frac{3}{25} \rightarrow 3a + 4b = 3$, $b = \frac{3}{4}(1-a)$ y además $a^2 + b^2 = 1$
 $\rightarrow a^2 + \frac{9}{16}a^2 - \frac{9}{8}a + \frac{9}{16} = 1$, $25a^2 - 18a - 7 = 0$, $a = 1$ ó $-\frac{7}{25} \rightarrow b = 0$ ó $\frac{24}{25}$. $\mathbf{u} = \left(-\frac{7}{25}, \frac{24}{25} \right)$ otro.

$$\Delta f = -2x(x^2+y^2)^{-2} + 4x(x^2-y^2)(x^2+y^2)^{-3} - 2x(x^2+y^2)^{-2} + 8xy^2(x^2+y^2)^{-3} = \frac{4x}{(x^2+y^2)^3} [-x^2 - y^2 + x^2 - y^2 + 2y^2] = 0.$$

Mejor en polares: $f(r, \theta) = \frac{\cos \theta}{r}$, $f_r = -\frac{\cos \theta}{r^2}$, $f_{rr} = \frac{2 \cos \theta}{r^3}$, $f_{\theta\theta} = -\frac{\cos \theta}{r}$, $\Delta f = f_{rr} + \frac{1}{r} f_r + \frac{1}{r^2} f_{\theta\theta} = 0$.

14. a) $\frac{x^2+y^2}{2+x^2+y^2} = C \rightarrow x^2+y^2 = \frac{2C}{1-C} = 1, 2, 4$ para $C = \frac{1}{3}, \frac{1}{2}, \frac{2}{3}$.

$$z = \frac{y^2}{2+y^2} \xrightarrow{y \rightarrow \pm\infty} 1; \text{ par; } z' = \frac{4y}{(2+y^2)^2}; z(0) = 0, z(1) = \frac{1}{3}, z(2) = \frac{2}{3}.$$

b) $\nabla f = \left(\frac{4x}{(2+x^2+y^2)^2}, \frac{4y}{(2+x^2+y^2)^2} \right)$. $\nabla f(1, -1) = \left(\frac{1}{4}, -\frac{1}{4} \right)$.

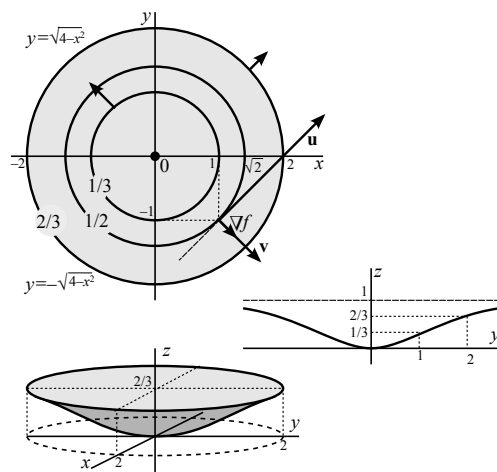
Plano tangente: $z = \frac{1}{2} + \frac{1}{4}(x-1) - \frac{1}{4}(y+1)$, $z = \frac{1}{4}(x-y)$.

$$D_{(2,2)} f(1, -1) = \left(\frac{1}{4}, -\frac{1}{4} \right) \cdot (2, 2) = 0 \text{ [perpendicular } \mathbf{u} \text{ al gradiente]}$$

Máxima en la dirección y sentido de ∇f : $\mathbf{v} = \left(\frac{1}{\sqrt{2}}, -\frac{1}{\sqrt{2}} \right)$.

c) Mejor en polares: $f(r) = \frac{r^2}{2+r^2}$, $f_r = \frac{4r}{(2+r^2)^2}$, $f_{rr} = \frac{8-12r^2}{(2+r^2)^3}$.

$$\Delta f = f_{rr} + \frac{1}{r} f_r = \frac{8(2-r^2)}{(2+r^2)^3} = \frac{8(2-x^2-y^2)}{(2+x^2+y^2)^3}.$$



15. $\mathbf{f}(x, y) = (x^2, 1, y^2)$, $g(x, y, z) = z$ a) $\text{div } \mathbf{f} = 2x$. $\text{rot } \mathbf{f} = \begin{vmatrix} \mathbf{i} & \mathbf{j} & \mathbf{k} \\ \partial/\partial x & \partial/\partial y & \partial/\partial z \\ x^2 & 1 & y^2 \end{vmatrix} = (2y, 0, 0)$. $\nabla g = (0, 0, 1)$. $\Delta g = 0$.

$$\text{rot}(\nabla g) = \mathbf{0}, \text{div}(\text{rot } \mathbf{f}) = 0 \text{ (sin necesidad de calcular nada). } \nabla(\mathbf{f} \cdot \nabla g) = \nabla(y^2) = (0, 2y, 0).$$

$$\text{rot}(\mathbf{f} \times \nabla g) = \text{rot}(1, -x^2, 0) = (0, 0, -2x). \text{rot}(\nabla(\mathbf{f} \cdot \nabla g)) = \mathbf{0}, \text{div}(\text{rot}(\mathbf{f} \times \nabla g)) = 0 \text{ (de nuevo sin calcular).}$$

b) $\text{div}(g \mathbf{f}) = g_x f_1 + g f_{1x} + g_y f_2 + g f_{2y} + g_z f_3 + g f_{3z} = g(f_{1x} + f_{2y} + f_{3z}) + (g_x, g_y, g_z) \cdot (f_1, f_2, f_3) = g \text{div } \mathbf{f} + \nabla g \cdot \mathbf{f}$.

$$\text{rot}(g \mathbf{f}) = \begin{vmatrix} \mathbf{i} & \mathbf{j} & \mathbf{k} \\ \partial/\partial x & \partial/\partial y & \partial/\partial z \\ g f_1 & g f_2 & g f_3 \end{vmatrix} = (g_y f_3 + g f_{3y} - g_z f_2 - g f_{2z}, g_z f_1 + g f_{1z} - g_x f_3 - g f_{3x}, g_x f_2 + g f_{2x} - g_y f_1 - g f_{1y})$$

$$= g(f_{3y} - f_{2z}, f_{1z} - f_{3x}, f_{2x} - f_{1y}) + (g_x, g_y, g_z) \times (f_1, f_2, f_3) = g \text{rot } \mathbf{f} + \nabla g \times \mathbf{f}.$$

Comprobando:

$$\text{div}(g \mathbf{f}) = 2xz + y^2 = z(2x + (0, 0, 1) \cdot (x^2, 1, y^2)). \text{rot}(g \mathbf{f}) = (2yz - 1, x^2, 0) = (2yz, 0, 0) + \begin{vmatrix} \mathbf{i} & \mathbf{j} & \mathbf{k} \\ 0 & 0 & 1 \\ x^2 & 1 & y^2 \end{vmatrix}.$$

16. $\mathbf{F}(x, y, z) = xy \mathbf{i} + y^2 \mathbf{j} + xz \mathbf{k}$. $\text{div } \mathbf{F} = 3y + x$. $\text{rot } \mathbf{F} = \begin{vmatrix} \mathbf{i} & \mathbf{j} & \mathbf{k} \\ \partial/\partial x & \partial/\partial y & \partial/\partial z \\ xy & y^2 & xz \end{vmatrix} = (0, -z, -x)$. $\nabla(\text{div } \mathbf{F}) = (1, 3, 0)$.

$$\text{div}(\text{rot } \mathbf{F}) = 0 \text{ (sabido). } \text{rot}(\text{rot } \mathbf{F}) = \begin{vmatrix} \mathbf{i} & \mathbf{j} & \mathbf{k} \\ \partial/\partial x & \partial/\partial y & \partial/\partial z \\ 0 & -z & -x \end{vmatrix} = (1, 1, 0). \nabla(\mathbf{F} \cdot \mathbf{F}) = (2x(y^2 + z^2), 2x^2 y + 4x^3, 2x^2 z).$$

Soluciones de problemas 2 de MM(im) (2020)

1. a) $\int_0^1 \int_0^1 (x^2 + y^2) dx dy = \int_0^1 \int_0^1 x^2 dx dy + \int_0^1 \int_0^1 y^2 dy dx = \int_0^1 x^2 dx + \int_0^1 y^2 dy = \frac{2}{3}.$

b) $\int_1^2 \int_1^2 \log(xy) dx dy = \int_1^2 \log x dx + \int_1^2 \log y dy = 2[2 \log 2 - 1],$ pues $\int_1^2 \log s ds = s \log s \Big|_1^2 - \int_1^2 1 ds = 2 \log 2 - 1.$

c) $\int_0^1 \int_0^1 y e^{xy} dx dy = \int_0^1 [e^{xy}]_0^1 dy = \int_0^1 [e^y - 1] dy = e - 1 - 1 = e - 2.$

d) $\int_0^1 \int_{x^2}^x xy dy dx = \int_0^1 x \left[\frac{y^2}{2} \right]_{x^2}^x dx = \frac{1}{2} \int_0^1 [x^3 - x^5] dx = \frac{1}{2} \left[\frac{1}{4} - \frac{1}{6} \right] = \frac{1}{24}.$

O bien $\int_0^1 \int_y^{\sqrt{y}} xy dx dy = \int_0^1 y \left[\frac{x^2}{2} \right]_y^{\sqrt{y}} dy = \frac{1}{2} \int_0^1 [y^2 - y^3] dy = \frac{1}{2} \left[\frac{1}{3} - \frac{1}{4} \right] = \frac{1}{24}.$

e) $\int_{-1}^1 x \int_{-\sqrt{1-x^2}}^{\sqrt{1-x^2}} dy dx = 2 \int_{-1}^1 x \sqrt{1-x^2} dx = 0$ ó $\int_0^{2\pi} \cos \theta \int_0^1 r^2 dr d\theta = \frac{1}{2} \int_0^{2\pi} \cos \theta d\theta = 0.$

[Podemos decirlo sin cálculos pues es integral de función impar en recinto simétrico].

f) $\iint_D \sin x dx dy = \int_0^{\pi/2} \int_y^{\pi-y} \sin x dx dy = \int_0^{\pi/2} [-\cos x]_y^{\pi-y} dy = \int_0^{\pi/2} [\cos y - \cos(\pi-y)] dy = \int_0^{\pi/2} [2 \cos y] dy = 2 \sin y \Big|_0^{\pi/2} = 2$

Peor: $\iint_D \sin x dx dy = \int_0^{\pi/2} \int_0^x \sin x dy dx + \int_{\pi/2}^{\pi} \int_0^{\pi-x} \sin x dy dx = \int_0^{\pi/2} x \sin x dx + \int_{\pi/2}^{\pi} (\pi-x) \sin x dx$
 $= -x \cos x \Big|_0^{\pi/2} + \int_0^{\pi/2} \cos x dx - (\pi-x) \cos x \Big|_{\pi/2}^{\pi} - \int_{\pi/2}^{\pi} \cos x dx = 0 + 1 + 0 + 1 = 2$

g) Mucho más corto utilizando polares: $\int_0^{2\pi} \cos^2 \theta d\theta \int_0^1 r^3 dr = \frac{1}{8} \int_0^{2\pi} (1 + \cos 2\theta) d\theta = \frac{\pi}{4}.$

$\int_{-1}^1 x^2 \int_{-\sqrt{1-x^2}}^{\sqrt{1-x^2}} dy dx = 4 \int_0^1 x^2 \sqrt{1-x^2} dx = [x = \sin t] = 4 \int_0^{\pi/2} \sin^2 t \cos^2 t dt = \int_0^{\pi/2} \sin^2 2t dt = \frac{\pi}{4}.$

h) $\int_1^2 \int_0^{2-x} \frac{y}{x^2} dy dx = \frac{1}{2} \int_1^2 \frac{4-4x+x^2}{x^2} dx = \left[-\frac{2}{x} - 2 \ln |x| \right]_1^2 + \frac{1}{2} = \frac{3}{2} - 2 \ln 2.$

[f continua en D pues sólo no lo es en x=0].

Algo más largo: $\int_0^1 \int_1^{2-y} \frac{y}{x^2} dx dy = \int_0^1 y \left[1 - \frac{1}{2-y} \right] dy = \int_0^1 \left[y + 1 - \frac{2}{2-y} \right] dy.$

2. a) $f(x, y) = y^2 - x = 0 \rightarrow x = y^2, f = -1 \rightarrow x = y^2 + 1$ (parábolas). $\nabla f(x, y) = (-1, 2y).$

$f(1, -1) = 0, \nabla f(1, -1) = (-1, -2).$ Plano: $z = -(x-1) - 2(y+1), z = -x - 2y - 1.$

D_u mínima en sentido opuesto a $\nabla f: (1, 2) \rightarrow \vec{u} = \left(\frac{1}{\sqrt{5}}, \frac{2}{\sqrt{5}} \right).$ $D_u f(1, -1) = -\sqrt{5}.$

$f_{xx} + f_{yy} = 2. f(r, \theta) = r^2 \sin \theta - r \cos \theta. f_{rr} + \frac{f_r}{r} + \frac{f_{\theta\theta}}{r^2} = 2s^2 + 2s^2 - \frac{c}{r} + 2c^2 - 2s^2 + \frac{c}{r} = 2.$

b) $\int_0^2 \int_{-y}^{y/2} (y^2 - x) dx dy = \int_0^2 \left[y^2 \left(\frac{y}{2} + y \right) - \frac{y^2/4 - y^2}{2} \right] dy = \int_0^2 \left[\frac{3y^3}{2} + \frac{3y^2}{8} \right] dy = 7.$

Peor: $\int_{-2}^0 \int_{-x}^2 (y^2 - x) dy dx + \int_0^2 \int_{2x}^2 (y^2 - x) dy dx = \int_{-2}^0 \left(\frac{8+x^3}{3} - 2x - x^2 \right) dx + \int_0^2 \left(\frac{8-8x^3}{3} - 2x + 2x^2 \right) dx = 7.$

3. a) $\left[f(x, y) = \frac{y}{x+1} \right] = 0, 1, -1 \rightarrow$ rectas $y=0, y=x+1, y=-x-1.$

$\nabla f = \left(\frac{-y}{(x+1)^2}, \frac{1}{x+1} \right) \xrightarrow{(0,1)} (-1, 1). \Delta f = \frac{2y}{(x+1)^3}. D_v f(0, 1) = (-1, 1) \cdot \left(\frac{3}{5}, \frac{4}{5} \right) = \frac{1}{5}$

b) $\iint_D f = \int_0^1 \int_{\sqrt{x}}^1 \frac{y}{x+1} dy dx = \frac{1}{2} \int_0^1 \frac{1-x}{x+1} dx = \int_0^1 \left[\frac{1}{x+1} - \frac{1}{2} \right] dx = \ln 2 - \frac{1}{2}.$

Más largo: $\int_0^1 \int_0^{y^2} \frac{y}{x+1} dx dy = \int_0^1 y \ln(1+y^2) dy = \frac{y^2}{2} \ln(1+y^2) \Big|_0^1 - \int_0^1 \frac{y^3+y-y}{1+y^2} dy = \frac{1}{2} \ln 2 - \frac{1}{2} + \frac{1}{2} \ln 2.$

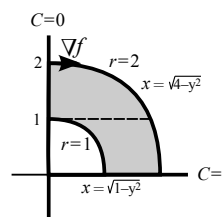
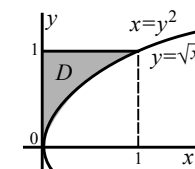
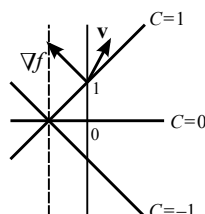
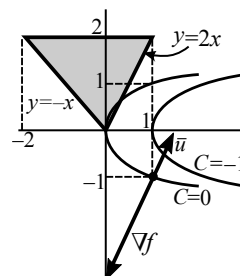
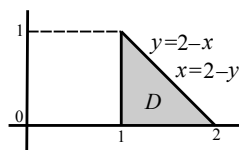
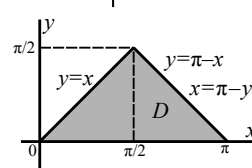
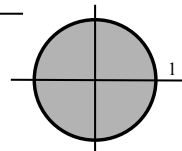
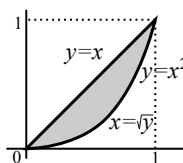
4. $f(x, y) = \frac{x}{\sqrt{x^2+y^2}}$ a) $f=0 \rightarrow x=0, f=1 \rightarrow x^2=x^2+y^2, y=0$ (si $x>0; f=\cos \theta$).

$\nabla f(x, y) = \frac{y}{(x^2+y^2)^{3/2}} (y, -x) \xrightarrow{(0,2)} \left(\frac{1}{2}, 0 \right)$ (perpendicular a la curva de nivel como debía).

Plano tangente: $z = 0 + \frac{1}{2}(x-0) + 0(y-1), z = \frac{x}{2}.$

b) En polares: $\iint_D f dx dy = \int_0^{\pi/2} \int_1^2 r \frac{\cos \theta}{r} dr d\theta = \left[\frac{r^2}{2} \right]_1^2 [\sin \theta]_0^{\pi/2} = \frac{3}{2}.$

[Peor: $\int f dx = \sqrt{x^2+y^2} \rightarrow \int_0^1 \int_{\sqrt{1-y^2}}^{\sqrt{4-y^2}} f dx dy + \int_1^2 \int_0^{\sqrt{4-y^2}} f dx dy = \int_0^1 [2-1] dy + \int_1^2 [2-y] dy = 3 - \left[\frac{y^2}{2} \right]_1^2 = \frac{3}{2}.$]



5. $g(x, y) = y\sqrt{x^2 + y^2}$. a) $\nabla g = (xy(x^2 + y^2)^{-1/2}, (x^2 + y^2)^{1/2} + y^2(x^2 + y^2)^{-1/2}) = \frac{1}{\sqrt{x^2 + y^2}}(xy, x^2 + 2y^2) \xrightarrow{(0, -2)} (0, 4)$.

Plano tangente: $z = -4 + 4(y + 2) = 4y + 4$. $g(r, \theta) = r^2 \sin \theta$. $g_{rr} + \frac{g_r}{r} + \frac{g_{\theta\theta}}{r^2} = 2s + 2s - s = 3 \sin \theta = 3y(x^2 + y^2)^{-1/2}$.

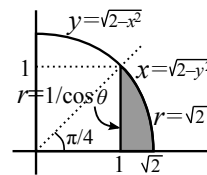
Más largo: $g_{xx} + g_{yy} = y(x^2 + y^2)^{-1/2} - x^2 y(x^2 + y^2)^{-3/2} + 3y(x^2 + y^2)^{-1/2} - y^3(x^2 + y^2)^{-3/2} \uparrow$

b) $\int_{\pi/2}^{\pi} \int_0^2 r^3 \sin \theta = \left[\frac{1}{4} r^4 \right]_0^2 \left[-\cos \theta \right]_{\pi/2}^{\pi} = 4$.

6. i) $\int_1^{\sqrt{2}} \int_0^{\sqrt{2-x^2}} x \, dy \, dx = \int_1^{\sqrt{2}} x \sqrt{2-x^2} \, dx = -\frac{1}{3}(2-x^2)^{3/2} \Big|_1^{\sqrt{2}} = \frac{1}{3}$,

o bien: $\int_0^1 \int_1^{\sqrt{2-y^2}} x \, dx \, dy = \int_0^1 \left[\frac{x^2}{2} \right]_1^{\sqrt{2-y^2}} dy = \int_0^1 \frac{1-y^2}{2} dy = \frac{1}{2} - \frac{1}{6} = \frac{1}{3}$.

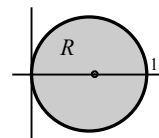
ii) $\int_0^{\pi/4} \int_{1/\cos \theta}^{\sqrt{2}} r^2 \cos \theta \, dr \, d\theta = \frac{1}{3} \int_0^{\pi/4} \left[2\sqrt{2} \cos \theta - \frac{1}{\cos^3 \theta} \right] d\theta = \frac{1}{3} [2\sqrt{2} \sin \theta - \tan \theta]_0^{\pi/4} = \frac{1}{3}$.



7. $M = 2 \int_0^{\pi/2} \int_0^{\cos \theta} r \cos \theta \, dr \, d\theta = \int_0^{\pi/2} (1 - \sin^2 \theta) \cos \theta \, d\theta = \frac{2}{3}$.

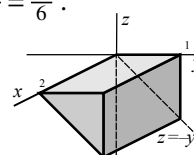
$\bar{x} = \frac{1}{M} 2 \int_0^{\pi/2} \int_0^{\cos \theta} r^2 \cos^2 \theta \, dr \, d\theta = \int_0^{\pi/2} (1 - \sin^2 \theta)^2 \cos \theta \, d\theta = \frac{8}{15}$.

$\bar{y} = 0$ por simetría de R y ρ .



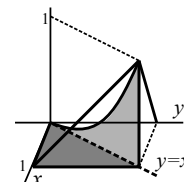
8. z positiva en el rectángulo: $V = \int_0^1 \int_1^2 (x^2 + y) \, dy \, dx = \int_0^1 x^2 \, dx + \int_1^2 y \, dy = \frac{1}{3} + \frac{4-1}{2} = \frac{11}{6}$.

9. a) $\int_0^2 \int_0^1 \int_{-y}^0 e^y \, dz \, dy \, dx = \int_0^2 \int_0^1 y e^y \, dy \, dx = 2 \int_0^1 y e^y \, dy = 2[(y-1)e^y]_0^1 = 2$.



b) $\iiint_V xy^2 z^3 \, dx \, dy \, dz = \int_0^1 \int_0^x \int_0^{xy} xy^2 z^3 \, dz \, dy \, dx = \frac{1}{4} \int_0^1 \int_0^x x^5 y^6 \, dy \, dx = \frac{1}{28} \int_0^1 x^{12} \, dx = \frac{1}{364}$.

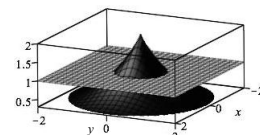
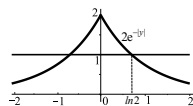
[$z=xy$ es un 'paraboloide hiperbólico' (silla de montar)].



10. $g(x, y) = 2e^{-\sqrt{x^2 + y^2}}$ de revolución. $g(0, y) = 2e^{-|y|}$ con pico.

$2e^{-\sqrt{x^2 + y^2}} = 1 \Leftrightarrow r = \sqrt{x^2 + y^2} = \ln 2$. Polares-cilíndricas:

$V = \int_0^{2\pi} \int_0^{\ln 2} r [2e^{-r} - 1] \, dr \, d\theta = 2\pi [-2(r+1)e^{-r} - \frac{r^2}{2}]_0^{\ln 2} = 2\pi [1 - \ln 2 - \frac{1}{2}(\ln 2)^2]$.



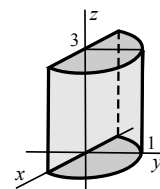
11. $F(x, y, z) = y + 2xz$. a) $\nabla F = (2z, 1, 2x) \xrightarrow{(1, 2, -1)} (-2, 1, 2)$. $D_v = 0$ si $\mathbf{v} \perp \nabla F$. Perpendicular también a $(1, -1, 0)$

$\rightarrow \mathbf{v} = \begin{vmatrix} \mathbf{i} & \mathbf{j} & \mathbf{k} \\ -2 & 1 & 2 \\ 1 & -1 & 0 \end{vmatrix} = (2, 2, 1)$. $\|\mathbf{v}\| = \sqrt{4+4+1} = 3$, $\mathbf{u} = (\frac{2}{3}, \frac{2}{3}, \frac{1}{3})$ (o $-\mathbf{u}$).

b) La integral triple se puede hacer en cilíndricas (mejor) o en cartesianas:

$\int_0^{\pi} \int_0^1 \int_0^3 r^2 (\sin \theta + 2z \cos \theta) \, dz \, dr \, d\theta = \int_0^{\pi} \int_0^1 3r^2 (\sin \theta + 3 \cos \theta) \, dr \, d\theta = \int_0^{\pi} (\sin \theta + 3 \cos \theta) \, d\theta = 2$.

$\int_{-1}^1 \int_0^{\sqrt{1-x^2}} \int_0^3 F \, dz \, dy \, dx = \int_{-1}^1 \int_0^{\sqrt{1-x^2}} (3y + 9x) \, dy \, dx = \int_{-1}^1 \left(\frac{3}{2} - \frac{3}{2}x^2 + 9x\sqrt{1-x^2} \right) dy \, dx = \int_0^1 (3-3x^2) \, dx = 2$.

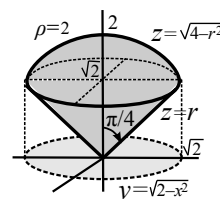


12. El cono $z = \sqrt{x^2 + y^2}$ y la esfera $x^2 + y^2 + z^2 = 4$ piden hacer la integral en esféricas

$\iiint_V z = 2\pi \int_0^{\pi/4} \int_0^2 \rho^3 \sin \phi \cos \phi \, d\rho \, d\phi = \pi \left[\frac{1}{4} \rho^4 \right]_0^2 \left[\sin^2 \phi \right]_0^{\pi/4} = 2\pi$.

En cilíndricas: $2\pi \int_0^{\sqrt{2}} \int_r^{\sqrt{4-r^2}} r z \, dz \, dr = 2\pi \int_0^{\sqrt{2}} r(2-r^2) \, dr = 2\pi \left[r^2 - \frac{r^4}{4} \right]_0^{\sqrt{2}} = 2\pi$.

Peor en cartesianas: $\int_{-\sqrt{2}}^{\sqrt{2}} \int_{-\sqrt{2-x^2}}^{\sqrt{2-x^2}} \int_{\sqrt{x^2+y^2}}^{\sqrt{4-x^2-y^2}} z \, dz \, dy \, dx = \int_{-\sqrt{2}}^{\sqrt{2}} \int_{-\sqrt{2-x^2}}^{\sqrt{2-x^2}} (2-x^2-y^2) \, dy \, dx = \dots$



13. a) $r = f(\theta)$, $\theta \in [\alpha, \beta] \rightarrow \mathbf{c}(\theta) = (f(\theta) \cos \theta, f(\theta) \sin \theta)$, $\mathbf{c}'(\theta) = (f' \cos \theta - f \sin \theta, f' \sin \theta + f \cos \theta)$

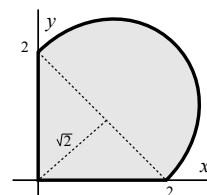
$\rightarrow \|\mathbf{c}'(\theta)\| = \sqrt{f^2(\cos^2 \theta + \sin^2 \theta) + (f')^2(\cos^2 \theta + \sin^2 \theta)} \rightarrow L = \int_{\alpha}^{\beta} \sqrt{[f(\theta)]^2 + [f'(\theta)]^2} \, d\theta$.

b) $r^2 = 2r \cos \theta + 2r \sin \theta$, $x^2 + y^2 = 2x + 2y$, $(x-1)^2 + (y-1)^2 = 2$, circunferencia.

i) $\int_0^{\pi/2} \int_0^{2(\cos \theta + \sin \theta)} r \, dr \, d\theta = \int_0^{\pi/2} (2+4 \sin \theta \cos \theta) \, d\theta = \pi + 2$.

ii) $r^2 + (r')^2 = 4[c^2 + 2sc + s^2 + s^2 - 2sc + c^2] = 8 \rightarrow L = 2+2 + \int_0^{\pi/2} \sqrt{8} \, d\theta = 4 + \pi\sqrt{2}$.

[Las integrales son innecesarias: i) $A = \frac{\pi(\sqrt{2})^2}{2} + \frac{2 \cdot 2}{2} = \pi + 2$; ii) $L = 2+2 + \frac{2\pi\sqrt{2}}{2} = 4 + \pi\sqrt{2}$].



14. a) $f(x, y, z) = yz$, $\mathbf{c}(t) = (t, 3t, 2t)$, $\mathbf{c}'(t) = (1, 3, 2)$; $\int_{\mathbf{c}} f \, ds = \int_1^3 6t^2 \sqrt{14} \, dt = 2\sqrt{14} t^3 \Big|_1^3 = 52\sqrt{14}$.

b) $f(x, y, z) = x + z$, $\mathbf{c}(t) = (t, t^2, \frac{2}{3}t^3)$, $\mathbf{c}'(t) = (1, 2t, 2t^2)$; $\int_{\mathbf{c}} f \, ds = \int_0^1 (t + \frac{2}{3}t^3)(1+2t^2) \, dt = \frac{1}{2}(t + \frac{2}{3}t^3)^2 \Big|_0^1 = \frac{25}{18}$.

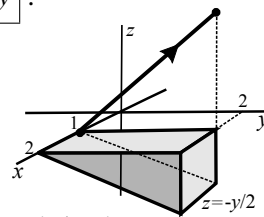
15. $F(x, y, z) = \frac{y+2z}{x}$. a) $\nabla F = \left(-\frac{y+2z}{x^2}, \frac{1}{x}, \frac{2}{x}\right)$. $\nabla F(1, 0, 1) = (-2, 1, 2)$. $\Delta F = 2\frac{y+2z}{x^3}$.

b) La $D_{\mathbf{u}}$ es máxima en la dirección y sentido del gradiente. $\|\nabla F\| = 3$. $\mathbf{u} = \left(-\frac{2}{3}, \frac{1}{3}, \frac{2}{3}\right)$. [Y la $D_{\mathbf{u}}F = 3$].

c) Plano tangente: $0 = (-2, 1, 2) \cdot (x-1, y, z-1) = -2x+2y+2z-2$, es decir, $z = x - \frac{1}{2}y$.
[Claro, la superficie $F=2$ es el plano $y+2z=2x$ y su plano tangente es él mismo].

d) $\iiint_V F = \int_1^2 \int_0^2 \int_{-y/2}^0 \frac{y+2z}{x} dz dy dx = \int_1^2 \frac{1}{x} \int_0^2 \frac{1}{4} y^2 dy dx = \frac{2}{3} \int_1^2 \frac{1}{x} dx = \left[\frac{2}{3} \ln 2\right]$.

e) $\bar{c}(t) = (1, 2t, 2t)$, $t \in [0, 1]$. $F(\bar{c}) = 6t$. $\|\bar{c}'\| = \sqrt{8}$. $\int_C F ds = \int_0^1 12\sqrt{2} t dt = [6\sqrt{2}]$.



16. $\mathbf{f}(x, y, z) = (xy, x, -yz)$. $\text{div } \mathbf{f} = y + 0 - y = 0$. $\text{rot } \mathbf{f} = \begin{vmatrix} \mathbf{i} & \mathbf{j} & \mathbf{k} \\ \partial/\partial x & \partial/\partial y & \partial/\partial z \\ xy & x & -yz \end{vmatrix} = (-z, 0, 1-x)$ [no deriva de un potencial].

$\int_C \mathbf{f} \cdot d\mathbf{s} = \int_0^\pi (\cos t \sin t, \cos t, -\sin t) \cdot (-\sin t, \cos t, 0) dt = \int_0^\pi (\cos^2 t - \sin^2 t \cos t) dt = \frac{\pi}{2} + \left[\frac{\sin 2t}{4} - \frac{\sin^3 t}{3}\right]_0^\pi = \frac{\pi}{2}$.

Como $\|\bar{c}'\| = \sqrt{\sin^2 t + \cos^2 t + 0} = 1$, la longitud de la curva es $L = \int_0^\pi 1 dt = \pi$.

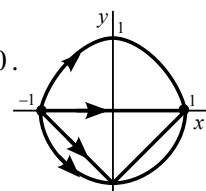
17. $\mathbf{f}(x, y) = (xy, 0)$ entre $(-1, 0)$ y $(1, 0)$. Usamos en todas el parámetro x , $x \in [-1, 1]$:

a) $\mathbf{c}(x) = (x, 0)$, $\mathbf{c}' = (1, 0)$. $\int_{-1}^1 0 dx = 0$. b) $\mathbf{c}(x) = (x, 1-x^2)$, $\mathbf{c}' = (1, -2x)$. $\int_{-1}^1 (x-x^3) dx = 0$.

c) $\mathbf{c}(x) = (x, |x|-1)$, $\mathbf{c}' = \begin{cases} (1, -1), & x < 0 \\ (1, 1), & x > 0 \end{cases}$. $\int_{-1}^0 (-x^2-x) dx + \int_0^1 (x^2-x) dx = 0$.

d) $\mathbf{c}(x) = (x, -\sqrt{1-x^2})$, $\mathbf{c}' = (1, x(1-x^2)^{-1/2})$. $\int_{-1}^1 -x\sqrt{1-x^2} dx = 0$.

Pero no es gradiente de ningún campo escalar pues $f_y = x \neq 0 = g_x$ (sobre otros caminos no se anulará).



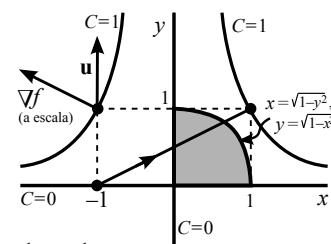
18. a) $f=0 \rightarrow x=0$ o $y=0$; $f=1 \rightarrow y = \frac{1}{x^2}$. ∇f perpendicular a la curva de nivel.

$\nabla f(x, y) = (2xy, x^2)$. $\nabla f(-1, 1) = (-2, 1)$. $\mathbf{u} = \mathbf{j}$ lo cumple.

c) Cartesianas: $\iint_D f = \int_0^1 \int_0^{\sqrt{1-y^2}} x^2 y dy dx = \frac{1}{2} \int_0^1 x^2 (1-x^2) dx = \frac{1}{2} \left[\frac{1}{3} - \frac{1}{5}\right] = \frac{1}{15}$.

Más largo: $\int_0^1 \int_0^{\sqrt{1-y^2}} x^2 y dx dy = \frac{1}{3} \int_0^1 y (1-y^2)^{3/2} dy = -\frac{1}{15} (1-y^2)^{5/2} \Big|_0^1 = \frac{1}{15}$.

En polares: $\iint_D f = \int_0^{\pi/2} \int_0^1 r r^3 \cos^2 \theta \sin \theta dr d\theta = \left[\frac{1}{5} r^5\right]_0^1 \left[-\frac{1}{3} \cos^3 \theta\right]_0^{\pi/2} = \frac{1}{5} \cdot \frac{1}{3} = \frac{1}{15}$.



d) Como $\mathbf{g} = \nabla f$, para hallar la integral de línea basta calcular $f(1, 1) - f(-1, 0) = 1 - 0 = 1$.

Directamente: $\mathbf{c}(t) = (t, \frac{1+t}{2})$, $t \in [-1, 1] \rightarrow \int_C \mathbf{g} \cdot d\mathbf{s} = \int_{-1}^1 (t+t^2, t^2) \cdot (1, \frac{1}{2}) dt = \int_{-1}^1 (t + \frac{3}{2}t^2) dt = 0 + 1 = 1$.

19. $\mathbf{f}(x, y) = (y, x)$. a) $\text{div } \mathbf{f} = f_x + g_y = 0$. Como $g_x = 1 = f_y$ y $\mathbf{f}$ es $C^1(\mathbb{R}^2)$, deriva de un potencial:

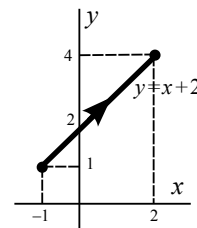
$U_x = y \rightarrow U = xy + p(y) \rightarrow U(x, y) = xy$.
 $U_y = x \rightarrow U = xy + q(x)$

b) i) La recta tiene pendiente 1 y corta $x=0$ en $y=2$. Es, por tanto, $y=x+2$.

Una parametrización sería, pues: $\mathbf{c}(t) = (t, t+2)$, $t \in [-1, 2]$, $\mathbf{c}'(t) = (1, 1) \rightarrow$

$\int_C \mathbf{f} \cdot d\mathbf{s} = \int_{-1}^2 (t+2, t) \cdot (1, 1) dt = \int_{-1}^2 (2t+2) dt = t^2 \Big|_{-1}^2 + 6 = [9]$.

ii) Con el potencial: $\int_C \mathbf{f} \cdot d\mathbf{s} = U(2, 4) - U(-1, 1) = 8 + 1 = [9]$.



20. $\bar{f}(x, y) = (y+3x^2, x)$. a) $\text{div } \bar{f} = 6x + 0 = [6x]$. $g_x = 1 = f_y$ y $\bar{f} \in C^1 \rightarrow$ deriva de un potencial.

b) Para calcular el valor de la integral de línea podemos calcular el potencial:

$U_x = y+3x^2 \rightarrow U = xy + x^3 + p(y) \rightarrow U = xy + x^3$, $\int_C \bar{f} \cdot d\bar{s} = U(-1, 0) - U(1, 0) = [-2]$.
 $U_y = x \rightarrow U = xy + q(x)$

O podemos parametrizar la curva y calcular la integral: $\bar{c}(t) = (\cos t, \sin t)$, $t \in [0, \pi] \rightarrow$

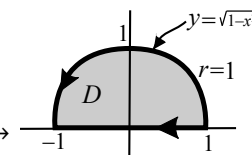
$\int_C \bar{f} \cdot d\bar{s} = \int_0^\pi (s+3c^2, c) \cdot (-s, c) dt = \int_0^\pi [c^2 - s^2 - 3c^2 s] dt = \left[\frac{1}{2} \sin 2t + \cos^3 t\right]_0^\pi = [-2]$.

O, como no depende del camino, hallar la integral sobre el sencillo segmento $\bar{c}_*(x) = (x, 0)$, $x \in [1, -1]$:

$\int_{\bar{c}_*} \bar{f} \cdot d\bar{s} = \int_1^{-1} (3x^2, x) \cdot (1, 0) dx = \int_1^{-1} 3x^2 dx = [x^3]_1^{-1} = [-2]$.

c) $\int_0^1 \int_{-\sqrt{1-y^2}}^{\sqrt{1-y^2}} 6x dx dy = \int_0^1 0 dx = [0]$ o $\int_{-1}^1 \int_0^{\sqrt{1-x^2}} 6x dy dx = \int_{-1}^1 6x \sqrt{1-x^2} dx = -2(1-x^2)^{3/2} \Big|_{-1}^1 = [0]$.

$\int_0^\pi \int_0^1 6r^2 \cos \theta dr d\theta = 2 \int_0^\pi \cos \theta d\theta = 2 \sin \theta \Big|_0^\pi = [0]$. [Debía anularse por ser $6x$ impar y D simétrico].



21. D limitado por $y=-2$, $y=0$, $x+2y=0$ y $x+2y=2$.

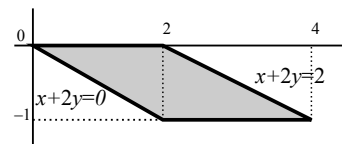
$$a) \iint_D (x+2y) dx dy = \int_{-1}^0 \int_{-2y}^{2-2y} (x+2y) dx dy = \int_{-1}^0 \left(\left[\frac{x^2}{2} \right]_{-2y}^{2-2y} + 4y \right) dy = \int_{-1}^0 2 dy = 2.$$

O bien: $u=x+2y$, $x=u-2v$, $\frac{\partial(x,y)}{\partial(u,v)} = \begin{vmatrix} 1 & -2 \\ 0 & 1 \end{vmatrix} = 1 \rightarrow \int_0^2 \int_{-1}^0 u dv du = \int_0^2 u du = 2.$

b) Como $\mathbf{f}(x, y) = (1, \cos y)$ cumple $(1)_y = 0 = (\cos y)_x \Rightarrow$ deriva de potencial ($U = x + \sin y$) $\Rightarrow \oint_{\partial D} \mathbf{f} \cdot d\mathbf{s} = 0.$

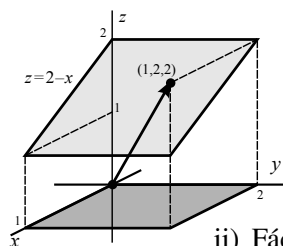
Directamente (largo): $\mathbf{c}_1 = (t, -\frac{t}{2})$, $t \in [0, 2]$, $\mathbf{c}_2 = (t, -1)$, $t \in [2, 4]$, $\mathbf{c}_3 = (t, 1 - \frac{t}{2})$, $t \in [4, 2]$, $\mathbf{c}_4 = (t, 0)$, $t \in [2, 0]$,

$$\int_0^2 [1 - \frac{1}{2} \cos \frac{t}{2}] dt + \int_2^4 dt + \int_4^2 [1 - \frac{1}{2} \cos(1 - \frac{t}{2})] dt + \int_2^0 dt = -\frac{1}{4} \sin 1 + \frac{1}{4} \sin 1 = 0.$$



22. a) $\nabla g = (2ye^{2x-z}, e^{2x-z}, -ye^{2x-z}) \xrightarrow{(1,-1,2)} (-2, 1, 1)$. $D_{(a,b,c)} g(1, -1, 2) = \nabla g(1, -1, 2) \cdot (a, b, c) = b + c - 2a$.

La D_v es nula, por ejemplo, según el vector $(1, 1, 1)$. Haciéndolo unitario: $\mathbf{u} = (\frac{1}{\sqrt{3}}, \frac{1}{\sqrt{3}}, \frac{1}{\sqrt{3}})$ (perpendicular al gradiente).



b) En $[0,1] \times [0,2]$ está $z=2-x$ por encima de $z=0$ (no se precisa el dibujo).

$$\iiint_V g = \int_0^2 \int_0^1 \int_0^{2-x} y e^{2x-z} dz dx dy = \int_0^2 y dy \int_0^1 [e^{2x} - e^{3x-2}] dx = e^2 - 1 - \frac{2}{3} e + \frac{2}{3} e^{-2}.$$

c) i) $\mathbf{c}(t) = (t, 2t, 2t)$ (salta a la vista). $\mathbf{c}'(t) = (1, 2, 2)$. $\|\mathbf{c}'(t)\| = \sqrt{1+4+4} = 3.$

$$\int_C \mathbf{g} \cdot d\mathbf{s} = \int_0^1 \mathbf{g}(\mathbf{c}(t)) \|\mathbf{c}'(t)\| dt = \int_0^1 6t dt = 3.$$

ii) Fácil: $\int_C \nabla g \cdot d\mathbf{s} = g(1, 2, 2) - g(0, 0, 0) = 2$. [Peor: $\int_0^1 (4t, 1, -2t) \cdot (1, 2, 2) dt = \int_0^1 2 dt = 2$].

23. $\mathbf{g}(x, y, z) = (2ye^{2x}, e^{2x}, 2z)$. $\text{div } \mathbf{g} = 4ye^{2x} + 2$. $\nabla(\text{div } \mathbf{g}) = (8ye^{2x}, 4e^{2x}, 0)$. $\text{rot } \mathbf{g} = \begin{vmatrix} \mathbf{i} & \mathbf{j} & \mathbf{k} \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ 2ye^{2x} & e^{2x} & 2z \end{vmatrix} = (0, 0, 2e^{2x} - 2e^{2x}) = \mathbf{0}.$

Podemos usar la definición sobre el camino dado [parametrizando a ojo o usando $\bar{c}(t) = \bar{a} + t(\bar{b} - \bar{a})$]:

$$\bar{c}(t) = (t, 2t, 1), t \in [0, 1] \rightarrow \int_{\bar{c}} \bar{\mathbf{g}} \cdot d\bar{\mathbf{s}} = \int_0^1 (4te^{2t}, e^{2t}, 2) \cdot (1, 2, 0) dt = \int_0^1 (4t+2) e^{2t} dt = [2te^{2t}]_0^1 = 2e^2.$$

O podemos calcular la función potencial [rot $\mathbf{g} = \mathbf{0}$ y $\mathbf{g} \in C^1$] y evaluarla en los puntos final e inicial:

$$U_x = 2ye^{2x} \rightarrow U = ye^{2x} + p(y, z)$$

$$U_y = e^{2x} \rightarrow U = ye^{2x} + q(x, z) \rightarrow U = ye^{2x} + z^2 \rightarrow \int_{\bar{c}} \bar{\mathbf{g}} \cdot d\bar{\mathbf{s}} = U(1, 2, 1) - U(0, 0, 1) = 2e^2 + 1 - 1 = 2e^2.$$

$$U_z = 2z \rightarrow U = z^2 + r(x, y)$$

24. $\bar{\mathbf{g}}(x, y, z) = (z, y^2, x)$. a) $\text{div } \bar{\mathbf{g}} = 2y$. $\nabla(\text{div } \bar{\mathbf{g}}) = (0, 2, 0)$. $\text{rot } \bar{\mathbf{g}} = \begin{vmatrix} \mathbf{i} & \mathbf{j} & \mathbf{k} \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ z & y^2 & x \end{vmatrix} = (0-0, 1-1, 0-0) = \bar{\mathbf{0}}.$ (y $\bar{\mathbf{g}} \in C^1 \Rightarrow \bar{\mathbf{g}}$ conservativo).

b) Con la definición: $\int_{\bar{c}} \bar{\mathbf{g}} \cdot d\bar{\mathbf{s}} = \int_0^{\pi/2} (-s, s^2, c) \cdot (-s, c, -c) dt = \int_0^{\pi/2} [s^2 - c^2 + s^2 c] dt = [\frac{1}{3} \sin^3 t - \frac{1}{2} \sin 2t]_0^{\pi/2} = \frac{1}{3}.$

$$U_x = x \rightarrow U = xz + p(y, z)$$

$$\text{Hallando el potencial: } U_y = y^2 \rightarrow U = \frac{1}{3} y^3 + q(x, z) \rightarrow U = xz + \frac{1}{3} y^3 \rightarrow \int_{\bar{c}} \bar{\mathbf{g}} \cdot d\bar{\mathbf{s}} = U(0, 1, -1) - U(1, 0, 0) = \frac{1}{3}.$$

$$U_z = z \rightarrow U = xz + r(x, y)$$

$$\text{Integrando sobre el segmento } \bar{c}_*(t) = (1-t, t, -t), t \in [0, 1]: \int_{\bar{c}_*} \bar{\mathbf{g}} \cdot d\bar{\mathbf{s}} = \int_0^1 (-t, t^2, t) \cdot (-1, 1, -1) dt = \int_0^1 t^2 dt = \frac{1}{3}.$$

25. $\mathbf{f}(x, y, z) = (z^2, 2y, cxz)$. $\text{div } \mathbf{f} = 2+cx$. $\text{rot } \mathbf{f} = \begin{vmatrix} \mathbf{i} & \mathbf{j} & \mathbf{k} \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ z^2 & 2y & cxz \end{vmatrix} = (0, 2z - cz, 0) = \mathbf{0}$ si $c=2$ y $\mathbf{f} \in C^1(\mathbf{R}^3).$

Si $c=2$ existe el U : $U_x = z^2 \rightarrow U = xz^2 + p(y, z)$
 $U_y = 2y \rightarrow U = y^2 + q(x, z)$, $U = xz^2 + y^2 \rightarrow \int_C \mathbf{f} \cdot d\mathbf{s} = U(1, 0, 1) - U(0, 0, 0) = 1.$
 $U_z = 2xz \rightarrow U = xz^2 + r(x, y)$

Sin necesidad de hacer integrales de línea.

[Una parametrización sería: $\mathbf{c}(t) = (t, 0, t)$, $t \in [0, 1] \rightarrow \int_C \mathbf{f} \cdot d\mathbf{s} = \int_0^1 (t^2, 0, 2t^2) \cdot (1, 0, 1) dt = \int_0^1 3t^2 dt = 1$].

26. a) $\mathbf{f}(x, y, z) = 2xz \mathbf{i} + \mathbf{j} + x^2 \mathbf{k}$. $\text{div } \mathbf{f} = 2z$. $\text{rot } \mathbf{f} = \begin{vmatrix} \mathbf{i} & \mathbf{j} & \mathbf{k} \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ 2xz & 1 & x^2 \end{vmatrix} = \mathbf{0}$. $\nabla(\text{div } \mathbf{f}) = (0, 0, 2)$.
 $\Delta(\mathbf{f} \cdot \mathbf{f}) = \Delta(4x^2 z^2 + 1 + x^4) = 20x^2 + 8z^2.$

b) $\iiint_V 2z dx dy dz = \int_0^3 \int_0^2 \int_0^{2-x} 2z dz dx dy = 3 \int_0^2 (2-x)^2 dx = -(2-x)^3 \Big|_0^2 = 8.$

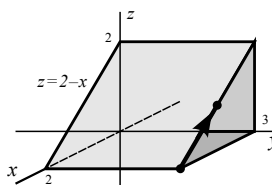
c) Como $\text{rot } \mathbf{f} = \mathbf{0}$ y $\mathbf{f}$ es C^1 en $\mathbf{R}^3$, existe la función potencial.

$$U_x = 2xz \rightarrow U = x^2 z + p(y, z)$$

$$U_y = 1 \rightarrow U = y + q(x, z)$$

$$U_z = x^2 \rightarrow U = x^2 z + r(x, y)$$

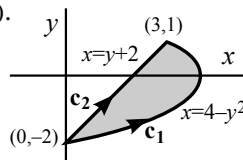
[Con $\mathbf{c}(t) = (2-t, 3, t)$, $t \in [0, 1] \rightarrow \int_C \mathbf{f} \cdot d\mathbf{s} = \int_0^1 (4t-2t^2, 1, (2-t)^2) \cdot (-1, 0, 1) dt = \int_0^1 (4-8t+3t^2) dt = 1$].



27. a) $\mathbf{f}(x, y) = (y^2, 2x)$, $\mathbf{c}_1 = (4-t^2, t)$, $t \in [-2, 1]$, $\mathbf{c}_2 = (t, t-2)$, $t \in [0, 3]$ (sentido opuesto).

$$\int_{\partial D} \mathbf{f} \cdot d\mathbf{s} = \int_{-2}^1 (-2t^3 + 8 - 2t^2) dt - \int_0^3 (t^2 - 2t + 4) dt = 26 - \frac{1}{2} - 12 = \frac{27}{2}.$$

$$g_x - f_y = 2 - 2y, \int_{-2}^1 \int_{y+2}^{4-y^2} (2-2y) dx dy = \int_{-2}^1 (4-6y+2y^3) dy = \frac{27}{2}.$$

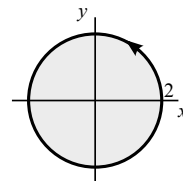


- b) $\mathbf{f}(x, y) = (0, xy^2)$. $\mathbf{c}(t) = (2 \cos t, 2 \sin t)$, $t \in [0, 2\pi]$

$$\rightarrow \int_C \mathbf{f} \cdot d\mathbf{s} = (0, 8cs^2) \cdot (-2s, 2c) dt = \int_0^{2\pi} 16 s^2 c^2 dt = 4\pi - \left[\frac{\sin 4t}{2} \right]_0^{2\pi} = 4\pi.$$

$$g_x - f_y = y^2, \iint_D y^2 dx dy = \int_0^{2\pi} \int_0^2 r^3 \sin^2 \theta dr d\theta = \left[\frac{1}{4} r^4 \right]_0^2 \cdot \frac{1}{2} \int_0^{2\pi} (1 - \cos 2\theta) d\theta = 4\pi.$$

[En cartesianas cálculos muy largos: $\int_{-2}^2 \int_{-\sqrt{4-x^2}}^{\sqrt{4-x^2}} y^2 dy dx = \frac{2}{3} \int_{-2}^2 (4-x^2)^{3/2} dx = \dots$].



- c) $g_x - f_y = -y$, $\iint_D -y dx dy = -\int_{\pi/4}^{5\pi/4} \int_0^{\sqrt{2}} r^2 \sin \theta dr d\theta = \left[\frac{1}{3} r^3 \right]_0^{\sqrt{2}} \left[\cos \theta \right]_{\pi/4}^{5\pi/4} = \frac{2\sqrt{2}}{3} \left[-\frac{1}{\sqrt{2}} - \frac{1}{\sqrt{2}} \right] = -\frac{4}{3}.$

En cartesianas (de las dos formas) es más complicado. Por ejemplo:

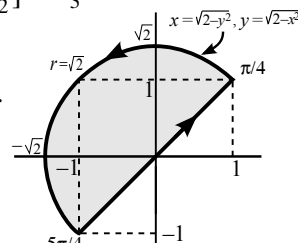
$$\iint_D = -\int_{-\sqrt{2}}^{-1} \int_{-\sqrt{2-x^2}}^{\sqrt{2-x^2}} y dy dx - \int_{-1}^1 \int_{-\sqrt{2-x^2}}^{\sqrt{2-x^2}} y dy dx = -\int_{-\sqrt{2}}^{-1} 0 dx + \int_{-1}^1 [x^2 - 1] dx = -\frac{4}{3}.$$

Parametrizaciones sencillas: $\mathbf{c}_1(t) = (t, t)$, $t \in [-1, 1]$ (en sentido correcto).

$$\mathbf{c}_2(t) = (\sqrt{2} \cos t, \sqrt{2} \sin t), t \in \left[\frac{\pi}{4}, \frac{5\pi}{4} \right] \text{ (también en buen sentido).}$$

$$\oint_{\partial D} \mathbf{f} \cdot d\mathbf{s} = \int_{-1}^1 (t^2, t^2) \cdot (1, 1) dt + \int_{\pi/4}^{5\pi/4} (2 \sin^2 t, 2 \sin t \cos t) \cdot (-\sqrt{2} \sin t, \sqrt{2} \cos t) dt$$

$$= \int_{-1}^1 2t^2 dt + 2\sqrt{2} \int_{\pi/4}^{5\pi/4} [s c^2 - s^3] dt = \frac{4}{3} + 2\sqrt{2} \left[\cos t - \frac{2}{3} \cos^3 t \right]_{\pi/4}^{5\pi/4} = \frac{4}{3} + 2 \left[1 + 1 - \frac{1}{3} - \frac{1}{3} \right] = -\frac{4}{3}.$$



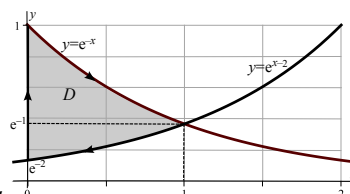
28. a) $\int_0^1 \int_{e^{x-2}}^{e^{-x}} x e^x dy dx = \int_0^1 [x - x e^{2x-2}] dx = \left[\frac{x^2}{2} - \frac{x}{2} e^{2x-2} \right]_0^1 + \frac{1}{2} \int_0^1 e^{2x-2} dx = \frac{1-e^{-2}}{4}.$

- b) $g_x - f_y = -x e^x$. Según Green, la $\oint_{\partial D} \mathbf{f} \cdot d\mathbf{s}$ vale lo de arriba. Directamente:

$$\mathbf{c}_1(t) = (0, t), y \in [e^{-2}, 1]; \mathbf{c}_2(t) = (t, e^{-t}), t \in [0, 1]; \mathbf{c}_3(t) = (t, e^{t-2}), t \in [1, 0].$$

$$\int_{\partial D} \mathbf{f} \cdot d\mathbf{s} = \int_{e^{-2}}^1 (0, 1) \cdot (0, 1) dt + \int_0^1 (t, 1) \cdot (1, -e^{-t}) dt + \int_1^0 (te^{2t-2}, 1) \cdot (1, e^{t-2}) dt$$

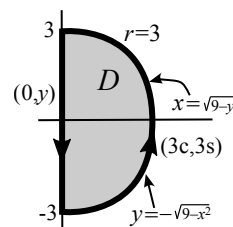
$$= 1 - e^{-2} + \int_0^1 (t - e^{-t}) dt + \int_1^0 (te^{2t-2} + e^{t-2}) dt = \frac{3}{2} - e^{-2} + [e^{-t}]_0^1 + \left[\frac{t}{2} e^{2t-2} - \frac{1}{4} e^{2t-2} + e^{t-2} \right]_1^0 = \frac{1-e^{-2}}{4}.$$



29. a) i) $\int_{-\pi/2}^{\pi/2} \int_0^3 r^2 \cos \theta dr d\theta = 2 [\sin \theta]_0^{\pi/2} \left[\frac{r^3}{3} \right]_0^3 = 18.$

ii) $\int_{-3}^3 \int_0^{\sqrt{9-y^2}} x dx dy = \frac{2}{2} \int_0^3 (9-y^2) dy = 27 - \left[\frac{y^3}{3} \right]_0^3 = 18.$

O bien: $\int_0^3 \int_{-\sqrt{9-x^2}}^{\sqrt{9-x^2}} x dy dx = \int_0^3 2x \sqrt{9-x^2} dx = -\frac{2}{3} (9-x^2)^{3/2} \Big|_0^3 = \frac{2}{3} 27 = 18.$



- b) i) Como $g_x - f_y = x$ el campo **no es conservativo** (debía ser $\equiv 0$). ii) $\text{div } \mathbf{f} = f_x + g_y = 1 - y.$

iii) Según Green el valor de la integral de línea en sentido antihorario coincide con el de la integral de a): 18.

Calculando la integral de línea directamente: $\mathbf{c}(t) = (3 \cos t, 3 \sin t)$, $t \in \left[-\frac{\pi}{2}, \frac{\pi}{2} \right]$. $\mathbf{c}_*(t) = (0, y)$, $y \in [3, -3]$.

$$\oint_{\partial D} \mathbf{f} \cdot d\mathbf{s} = \int_{-\pi/2}^{\pi/2} (-9cs, 3s) \cdot (-3s, 3c) dt + \int_3^{-3} (0, y) \cdot (0, 1) dy = \int_{-\pi/2}^{\pi/2} (27s^2 c + 9sc) dt - \int_3^{-3} y dy = 18s^3 \Big|_0^{\pi/2} = 18.$$

30. a) $y(1-2x)=0 \rightarrow$ rectas $y=0, x=\frac{1}{2}$. $y(1-2x)=1 \rightarrow y=\frac{1}{1-2x}$ $y=0, x=\frac{1}{2}$ asíntotas.

$$\nabla f(x, y) = (-2y, 1-2x), \nabla f(0, 1) = (-2, 1) \text{ (perpendicular curva } C=1).$$

$$D_{\mathbf{u}}=0 \text{ en la dirección perpendicular al gradiente: } \mathbf{u} = \left(\frac{1}{\sqrt{5}}, \frac{2}{\sqrt{5}} \right) \text{ o } \mathbf{u} = \left(-\frac{1}{\sqrt{5}}, -\frac{2}{\sqrt{5}} \right).$$

$$\text{div}(\nabla f) = \Delta f(x, y) = f_{xx} + f_{yy} = 0 + 0 = 0.$$

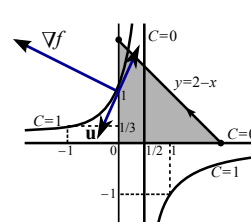
- b) $\iint_D f = \int_0^2 \int_0^{2-x} y(1-2x) dy dx = \frac{1}{2} \int_0^2 (2-x)^2 (1-2x) dx = \int_0^2 (2-6x+\frac{9}{2}x^2-x^3) dx = 0.$

Más corto: $\int_0^2 \int_0^{2-y} y(1-2x) dx dy = \int_0^2 y [(2-y) - (2-y)^2] dy = \int_0^2 [-2y+3y^2-y^3] dy = 0.$

- c) $\mathbf{g}(x, y) = (xy^2, xy)$. Como $g_x - f_y = y - 2xy \neq 0$, $\mathbf{g}$ no deriva de un potencial.

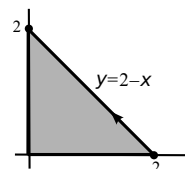
Una posible parametrización es: $\mathbf{c}(x) = (x, 2-x)$, $x \in [2, 0]$, $\mathbf{c}'(x) = (1, -1) \rightarrow$

$$\int_C \mathbf{g} \cdot d\mathbf{s} = \int_2^0 (x(2-x)^2, x(2-x)) \cdot (1, -1) dx = \int_2^0 (2x-3x^2+x^3) dx = x^2-x^3+\frac{1}{4}x^4 \Big|_2^0 = 0.$$



[Se podría hacer con Green en el triángulo, cuyo borde está formado por nuestro segmento y los ejes donde $\mathbf{g}=0$.

$$\oint_{\partial D} \mathbf{g} \cdot d\mathbf{s} = 0 + 0 + \int_C \mathbf{g} \cdot d\mathbf{s} = \iint_D [g_x - f_y] dx dy = \iint_D [y - 2xy] dx dy = 0 \text{ (calculada en b))}.$$



Soluciones de problemas 3 de MM(im) (2020)

1. a) $\boxed{y' = -\frac{2y}{x}}$ $y = C e^{\int a(x) dx} = C e^{-2 \int dx/x} = C e^{-2 \ln x} = C e^{\ln x^{-2}} = \boxed{\frac{C}{x^2}}$.
 - b) $\boxed{y' = \frac{y}{x^2}}$ $y = C e^{\int a(x) dx} = C e^{\int dx/x^2} = \boxed{C e^{-1/x}}$.
 - c) $\boxed{y' = y+5}$, $\boxed{y = C e^x - 5}$ (y_p a ojo). [Sin ojo: $y = C e^x + e^x \int 5 e^{-x} dx = C e^x - 5$]. [También es separable].
 - d) $\boxed{y' = \frac{y}{3x} + 2}$ $e^{\int dx/3x} = e^{\frac{1}{3} \ln x} = x^{1/3}$, $y = C x^{1/3} + x^{1/3} \int 2 x^{-1/3} dx = \boxed{C x^{1/3} + 3x}$.
[También es de la forma $y' = f(\frac{y}{x})$: $z = \frac{y}{x} \rightarrow xz' + z = \frac{z}{3} + 2$, $\int \frac{3dz}{3-z} = -3 \ln(3-z) = 2 \ln x + C$, $3-z = C x^{-2/3} \dots$].
 - e) $\boxed{\frac{dy}{dx} = 2xy^2}$ separable: $\int \frac{dy}{y^2} = \int 2x dx + C$, $-\frac{1}{y} = x^2 + C \rightarrow \boxed{y = \frac{1}{C-x^2}}$.
 - f) $\boxed{\frac{dy}{dx} = x + \frac{x}{y}}$ separable: $\int \frac{y dy}{1+y} = \int x dx + C \rightarrow \boxed{y - \log|1+y| = \frac{x^2}{2} + C}$ (no se puede despejar y).
 - g) $\boxed{\frac{dy}{dx} = x + \frac{y}{x}}$ lineal: $e^{\int dx/x} = x$, $y = Cx + x \int \frac{x dx}{x} \rightarrow \boxed{y = Cx + x^2}$.
 - h) $\boxed{\frac{dy}{dx} = e^{x-y}}$ separable: $\int e^y dy = \int e^x dx + C \rightarrow \boxed{y = \log(C + e^x)}$ [mucho más largo $z = x - y$, $z' = 1 - e^z \dots$]
2. $\boxed{\frac{dy}{dx} = -\frac{12x+5y}{5x+2y}}$ $12x+5y + (5x+2y) \frac{dy}{dx} = 0$, $M_y = 5 = N_x$ exacta. $U = 6x^2 + 5xy + p(y) = 5xy + y^2 + q(x) \downarrow$
O también: $z = \frac{y}{x}$, $xz' + z = -\frac{12+5z}{5+2z}$, $\int \frac{(2z+5) dz}{z^2+5z+6} = \int \frac{-2 dx}{x} + C$, $z^2 + 5z + 6 = \frac{C}{x^2}$, $\boxed{y^2 + 5xy + 6x^2 = C}$.
 3. a) $\boxed{T' = 4 - 4T}$, $T(0) = 0$. La de la homogénea es Ce^{-4t} y la $T_p = 1$ salta a la vista. General: $\boxed{T = Ce^{-4t} + 1}$.
O bien, con fvc: $T = Ce^{-4t} + e^{-4t} \int 4e^{4t} dt = Ce^{-4t} + 1$. $T(0) = C + 1 = 0 \rightarrow C = -1$, $\boxed{T = 1 - e^{-4t}}$.
 - b) $\boxed{T' + 9T = t}$, $T(0) = 0$. $T = Ce^{-9t} + e^{-9t} \int t e^{9t} dt = \boxed{Ce^{-9t} + \frac{t}{9} - \frac{1}{81}}$ (o probando $T_p = At + B \rightarrow A + 9At + 9B = t$, $A = \frac{1}{9}$, $B = -\frac{1}{81}$).
 $\frac{t}{9} e^{9t} - \frac{1}{9} \int e^{9t} dt$ $T(0) = C - \frac{1}{81} = 0 \rightarrow \boxed{T = \frac{1}{81}(e^{-9t} + 9t - 1)}$.
 - c) $\boxed{T' + (1+2t)T = 0}$, $T(0) = 0$. $T = Ce^{-\int (1+2t) dt} = \boxed{Ce^{-(t+t^2)}}$ $\xrightarrow{T(0)=0} \boxed{T \equiv 0}$.
[Las lineales homogéneas (de coeficientes continuos) con datos 0 tienen siempre por solución la trivial $T \equiv 0$].
 - d) $\boxed{T' = T + 2 \sin t}$ $T = Ce^t + 2e^t \int e^{-t} \sin t dt$. Integrando 2 veces por partes: $\boxed{T = Ce^t - \sin t - \cos t}$.
[Sin integrar: $T_p = A \cos t + B \sin t \rightarrow -As + Bc = Ac + Bs + 2s \rightarrow B = A$, $-A = B + 2 \rightarrow A = B = -1$].
 $T(0) = C - 1 = 0 \rightarrow \boxed{T = e^t - \sin t - \cos t}$.
4. a) $\boxed{y'' + 2y' + 5y = 0}$ $\mu^2 + 2\mu + 5 = 0 \rightarrow \mu = -1 \pm 2i$, $\boxed{y = (c_1 \cos 2x + c_2 \sin 2x) e^{-x}}$.
 - b) $\boxed{x^2 y'' - 3xy' + 3y = 0}$ Euler. $\mu(\mu-1) - 3\mu + 3 = \mu^2 - 4\mu + 3 = 0 \rightarrow \mu = 1, 3 \rightarrow \boxed{y = c_1 x + c_2 x^3}$.
 - c) $\boxed{x^2 y'' - x(x+2)y' + (x+2)y = 0}$. No es de coeficientes constantes, ni de Euler, ni falta el término en y .
Buscamos una solución que salte a la vista: $y_1 = x \rightarrow y_2 = x \int \frac{e^{\int \frac{x+2}{x^2} dx}}{x^2} = x e^x$, $\boxed{y = c_1 x + c_2 x e^x}$.
[Si no se puede ver esa solución (que es lo habitual), hay que acudir a series].
5. $\boxed{x^2 y'' + xy' - n^2 y = 0}$. Ecuación de Euler. $\mu(\mu-1) + \mu - 1 = n^2$, $\mu = \pm n \rightarrow$
 $\boxed{y(x) = c_1 x^n + c_2 x^{-n} \text{ si } n > 0}$
 $\boxed{y(x) = c_1 + c_2 \ln x \text{ si } n = 0}$

6. a) $y'' - 3y = e^{2x}$ $\mu = \pm\sqrt{3}$, $y_p = Ae^{2x}$, $4A - 3A = 1 \rightarrow y = c_1 e^{\sqrt{3}x} + c_2 e^{-\sqrt{3}x} + e^{2x}$.
 [Con la **fv** casi siempre es más largo: $|W| = -2\sqrt{3}$, $y_p = e^{-\sqrt{3}x} \int \frac{e^{(2+\sqrt{3})x}}{-2\sqrt{3}} - e^{\sqrt{3}x} \int \frac{e^{(2-\sqrt{3})x}}{-2\sqrt{3}} = \frac{1}{2\sqrt{3}} e^{2x} \left[\frac{1}{2-\sqrt{3}} - \frac{1}{2+\sqrt{3}} \right]$.
- b) $y'' + y = x \cos x$ $\mu = \pm i$, $y_p = (Ax^2 + Bx)c + (Cx^2 + Dx)s$, $y_p'' + y_p = (2A + 2D + 4Cx)c + (2C - 2B - 4Ax)s$
 $\rightarrow y = c_1 \cos x + c_2 \sin x + \frac{1}{4}x^2 \sin x + \frac{1}{4}x \cos x$.
- c) $y'' + y = 6 \cos^2 x$ $\begin{vmatrix} c & s \\ -s & c \end{vmatrix} = 1$, $y_p = s \int 6c^3 - c \int 6sc^2 = \dots \rightarrow y = c_1 \cos x + c_2 \sin x + 4 \sin^2 x + 2 \cos^2 x$.
 O bien, $6 \cos^2 x = 3 + 3 \cos 2x \rightarrow y_p = A + B \cos 2x + C \sin 2x \rightarrow y_p = 3 - \cos 2x$ ↑
- d) $y'' + 4y' + 5y = x$ Coeficientes constantes. $\mu^2 + 4\mu + 5 = 0$, $\mu = -2 \pm i \rightarrow y = e^{-2x}(c_1 \cos x + c_2 \sin x) + y_p$.
 $\mu = 0$ no autovalor $\rightarrow y_p = Ax + B \rightarrow 4A + 5Ax + 5B = x$, $y = e^{-2x}(c_1 \cos x + c_2 \sin x) + \frac{x}{5} - \frac{4}{25}$.
 [La **fv** nos llevaría a integrales bastante largas].
- e) $y'' - 2y' + y = x^2$ $\mu^2 - 2\mu + 1 = 0 \rightarrow \mu = 1$ doble, $y = (c_1 + c_2 x)e^x + y_p$. Como $\mu = 0$ no autovalor:
 $y_p = Ax^2 + Bx + C \rightarrow 2A - 4Ax - 2B + Ax^2 + Bx + C = x^2$, $A = 1$, $B = 4$, $C = 6$, $y = (c_1 + c_2 x)e^x + x^2 + 4x + 6$.
7. $y'' + 2y' - 3y = f(x)$ $\mu^2 + 2\mu - 3 = 0$, $\mu = -1 \pm \sqrt{4} = 1, -3 \rightarrow y = c_1 e^x + c_2 e^{-3x} + y_p$.
 a) $\mu = -1$ no autovalor. Probamos $y_p = Ae^{-x}$: $Ae^{-x} - 2Ae^{-x} - 3Ae^{-x} = e^{-x}$, $A = \frac{1}{4}$, $y = c_1 e^x + c_2 e^{-3x} - \frac{1}{4}e^{-x}$.
 b) $\mu = 1$ sí lo es. Probamos $y_p = Axe^x$: $A(x+2) + 2A(x+1) - 3Ax = 1$, $A = \frac{1}{4}$, $y = c_1 e^x + c_2 e^{-3x} - \frac{1}{4}xe^{-x}$.
 c) $\mu = \pm i$ no lo son. $y_p = A \cos x + B \sin x \rightarrow (2B - 4A)c - (2A + 4B)s = s$. $y = c_1 e^x + c_2 e^{-3x} - \frac{1}{10} \cos x - \frac{1}{5} \sin x$.
8. a) $y'' + 2y' + 2y = 2$ $\mu^2 + 2\mu + 2 = 0 \rightarrow \mu = -1 \pm i$, $y = (c_1 \cos x + c_2 \sin x)e^{-x} + 1$ (y_p a simple vista).
 $y(0) = y'(0) = 0 \rightarrow \begin{cases} c_1 + 1 = 0 \\ c_2 - c_1 = 0 \end{cases} \rightarrow y = 1 - e^{-x}(\cos x + \sin x)$.
 b) $y'' + 2y' + y = x + 2$ $\mu^2 + 2\mu + 1 = 0$, $\mu = -1$ doble $\rightarrow y = (c_1 + c_2 x)e^{-x} + y_p$.
 $y_p = Ax + B \rightarrow 2A + Ax + B = x + 2 \rightarrow A = 1$, $B = 0$. La solución general es: $y(x) = (c_1 + c_2 x)e^{-x} + x$.
 Imponiendo los datos: $\begin{cases} c_1 = 0 \\ c_2 - c_1 + 1 = 0, c_2 = -1 \end{cases}$. La solución es, pues: $y(x) = x(1 - e^{-x})$.
 c) $y'' + y = x^2$ $\mu^2 + 1 = 0$, $\mu = \pm i$. $y_p = Ax^2 + Bx + C$ ($\mu = 0$ no autovalor) $\rightarrow 2A + Ax^2 + Bx + C = x^2$, $A = 1$, $B = 0$, $C = -2$.
 $y = c_1 \cos x + c_2 \sin x + x^2 - 2 \xrightarrow{d.i.} \begin{cases} y(0) = c_1 - 2 = 0 \\ y'(0) = c_2 = 0 \end{cases}$, $y = 2 \cos x + x^2 - 2$.
 d) $y'' + y = 2x e^x$ $\mu = \pm i \rightarrow y = c_1 \cos x + c_2 \sin x + y_p$. $y_p = (Ax + B)e^x$, $y_p' = (Ax + B + A)e^x$, $y_p'' = (Ax + B + 2A)e^x$
 $(2Ax + 2B + 2A)e^x = 2x e^x \rightarrow A = 1$, $B = -1$. $y = c_1 \cos x + c_2 \sin x + (x - 1)e^x$. $\begin{cases} c_1 - 1 = 0 \\ c_2 = 0 \end{cases}$, $y = \cos x + (x - 1)e^x$.
 e) $y'' + y = \frac{2}{\cos^3 x}$ $\mu = \pm i$, $\begin{vmatrix} c & s \\ -s & c \end{vmatrix} = 1$, $y_p = s \int \frac{2}{c^2} - c \int \frac{s}{c^3} = \frac{1}{c} - 2c \rightarrow y = c_1 \cos x + c_2 \sin x + \frac{1}{\cos x}$
 $y(0) = y'(0) = 0 \xrightarrow{\rightarrow} y = \frac{1}{\cos x} - \cos x = \frac{\sec^2 x}{\cos x}$.
9. a) $\frac{dy}{dx} = -\frac{y-2}{x}$ i) Como lineal: $\frac{dy}{dx} = \frac{c}{x} + \frac{1}{x} \int 2 dx = -\frac{y}{x} + \frac{2}{x} \rightarrow y = \frac{c}{x} + 2$, $x(y-2) = C$.
 ii) Como separable: $\int \frac{dy}{y-2} = -\int \frac{dx}{x} + C$, $\ln(y-2) = C - \ln x$, $y-2 = \frac{C}{x}$. ↗
 iii) $(y-2) + x \frac{dy}{dx} = 0$, $M_y \equiv 1 \equiv N_x$. $\begin{matrix} U_x = y-2, U = xy-2x+p(y) \\ U_y = x, U = xy+q(x) \end{matrix} \rightarrow xy-2x = C$.
 b) $x^2 y'' + xy' = 2x$ i) $y' = v \rightarrow v' = -\frac{v}{x} + \frac{2}{x} \rightarrow v = \frac{c}{x} + 2$ (como arriba), $y = K + C \ln x + 2x$.
 ii) $\mu(\mu-1) + \mu = 0$, $\mu = 0$ doble, $y = c_1 + c_2 \ln x + y_p$. Para hallar la y_p :
 $\begin{vmatrix} 1 & \ln x \\ 0 & x^{-1} \end{vmatrix} = x^{-1}$, $y_p = \ln x \int \frac{1 \cdot 2/x}{x^{-1}} - 1 \int \frac{\ln x \cdot 2/x}{x^{-1}} = 2x$. O mejor, $y_p = Ax$ (Ae^s) $\rightarrow A = 2$.

10. a) $\boxed{xy'' + 2y' = x}$ $\mu(\mu-1) + 2\mu = 0 \rightarrow y = c_1 + c_2x^{-1} + y_p$ solución de la no homogénea.
 $\begin{vmatrix} 1 & x^{-1} \\ 0 & -x^{-2} \end{vmatrix} = -x^{-2}$, $y_p = \frac{1}{x} \int \frac{1 \cdot 1}{-x^{-2}} - 1 \int \frac{1/x \cdot 1}{-x^{-2}} = \frac{x^2}{6}$. O mejor, $y_p = Ax^2$ (Ae^{2s}) $\rightarrow 2A + 4A = 1$.
O también: $xv' + 2v = x \rightarrow v = \frac{C}{x^2} + \frac{x}{3} \rightarrow \boxed{y = \frac{1}{6}x^2 - \frac{C}{x} + K}$, como antes (con otro nombre de constantes).
- b) $\boxed{x^2y'' - 2y = 2}$ Euler. $\mu(\mu-1) + 0\mu - 2 = \mu^2 - \mu - 2 = (\mu-2)(\mu+1) = 0$ e $y_p = -1$ a ojo $\rightarrow \boxed{y = c_1x^2 + \frac{c_2}{x} - 1}$.
Sin vista: $|W| = \begin{vmatrix} x^2 & x^{-1} \\ 2x & -x^{-2} \end{vmatrix} = -3$. $y_p = x^{-1} \int \frac{x^2 2x^{-2}}{-3} dx - x^2 \int \frac{x^{-1} 2x^{-2}}{-3} dx = -\frac{2}{3x} \int dx + \frac{x^2}{3} \int \frac{2dx}{x^3} = -\frac{2}{3} - \frac{1}{3}$.
- c) $\boxed{x^2y'' + 4xy' + 2y = e^x}$ $y = \frac{c_1}{x} + \frac{c_2}{x^2} + y_p$, $|W| = -x^{-4}$, $y_p = -x^{-2} \int x e^x + x^{-1} \int e^x = \frac{e^x}{x^2}$. $\boxed{y = \frac{c_1}{x} + \frac{c_2}{x^2} + \frac{e^x}{x^2}}$.
11. $\boxed{(x+1)y'' - y' = (x+1)^2}$ $y' = v$, $v' = \frac{v}{x+1} + x+1$, $v = c_1(x+1) + x^2 + x$, $\boxed{y = c_1(x^2 + 2x) + c_2 + \frac{x^3}{3} + \frac{x^2}{2}}$.
O bien, $x+1 = s$, $sx'' - x' = s^2$ (Euler), $y_p = As^3$, $y = k_1s^2 + k_2 + \frac{s^3}{3} = k_1(x+1)^2 + k_2 + \frac{(x+1)^3}{3} \nearrow$
12. $\boxed{x^2y'' - 3xy' = 4}$. Como Euler: $\mu(\mu-1) - 3\mu = 0$, $\mu = 0, 4 \rightarrow y = c_1 + c_2x^4 + y_p$. Con la fvc:
 $\begin{vmatrix} 1 & x^4 \\ 0 & 4x^3 \end{vmatrix} = 4x^3$. $y_p = x^4 \int \frac{1 \cdot 4/x^2}{4x^3} dx - 1 \int \frac{x^4 \cdot 4/x^2}{4x^3} dx = x^4 \int \frac{dx}{x^5} - \int \frac{dx}{x} = -\frac{1}{4} - \ln|x|$, $\boxed{y = c_1 + c_2x^4 - \ln|x|}$.
[Haciendo $x = e^s \rightarrow y'' - 4y' = 4 \xrightarrow{y_p = As} y = c_1 + c_2e^{4s} - s \nearrow$].
O bien: $y' = v$, $v' = \frac{3}{x}v + \frac{4}{x^2}$, $e^{3 \ln x} = x^3$, $v = Cx^3 + x^3 \int \frac{4}{x^5} dx = Cx^3 - \frac{1}{x}$, $y = Cx^4 + K - \ln|x|$, como arriba.
Imponiendo los datos ($y' = 4c_2x^3 - \frac{1}{x}$): $\begin{cases} c_1 + c_2 = 3 \\ 4c_2 - 1 = 3 \end{cases} \rightarrow \boxed{y = 2 + x^4 - \ln|x|}$.
13. $\boxed{y'' - y = 3e^{-2x}}$. $\mu = \pm 1$ e $y_p = Ae^{-2x}$, $4A - A = 3 \rightarrow$ solución general $y = c_1e^x + c_2e^{-x} + e^{-2x}$.
i) $y(0) = 1$, $y'(0) = 0$. Imponiendo los datos: $\begin{cases} c_1 + c_2 + 1 = 1 \\ c_1 - c_2 - 2 = 0 \end{cases} \rightarrow c_1 = 1$, $c_2 = -1$. $\boxed{y = e^x - e^{-x} + e^{-2x}}$.
ii) $y'(0) = y'(1) = 0$. $\begin{cases} c_1 - c_2 = 2 \\ ec_1 - e^{-1}c_2 = 2e^{-2} \end{cases}$, sistema con solución única. Es $y = e^{-2x} - \frac{2}{e(e+1)} [e^x - (e^2 + e + 1)e^{-x}]$.
[Viendo 4.3: como el problema homogéneo tiene sólo la solución $y \equiv 0$ (por no ser $\lambda = -1$ autovalor), sin necesidad de echar cuentas se sabe que el no homogéneo tiene **solución única**].
14. a) $\boxed{y'' + y' = 2x - 1}$. i) $\mu^2 + \mu = 0 \rightarrow c_1 + c_2e^{-x} + y_p$. $y_p = Ax^2 + Bx$ ($\mu = 0$ autovalor) $\rightarrow$
 $2A + 2Ax + B = 2x - 1$, $A = 1$, $B = -3$, $\boxed{y = c_1 + c_2e^{-x} + x^2 - 3x}$, $y' = -c_2e^{-x} + 2x - 3$
ii) $y' = v \rightarrow v' = -v + 2x - 1 \rightarrow v = Ce^{-x} + e^{-x} \int e^x(2x-1) dx = Ce^{-x} + 2x - 3$, $y = \int v = K - Ce^{-x} + x^2 - 3x$.
b) $\begin{cases} y(0) = c_1 + c_2 = 0, & c_1 = 3 \\ y'(0) = -c_2 - 3 = 0, & c_2 = -3 \end{cases} \rightarrow \boxed{y = x^2 - 3x + 3 - 3e^{-x}}$. c) $\begin{cases} y'(0) = -c_2 - 3 = 0, & c_2 = -3 \\ y'(1) = -c_2e^{-1} - 1 = 0, & c_2 = -e \end{cases}$ Imposible. **No hay solución**.
[O utilizando los resultados de 4.3: $\begin{cases} y = c_1 + c_2e^{-x} \\ y' = -c_2e^{-x} \end{cases} \rightarrow \begin{cases} -c_2 = 0 \\ -c_2e^{-1} = 0 \end{cases}$. **Infinitas soluciones** del homogéneo $\{1\}$.
 $[e^x y']' = e^x(2x-1)$. $\int_0^1 e^x(2x-1) dx = e^x(2x-3)|_0^1 = 3 - e \neq 0 \Rightarrow$ el no homogéneo **no tiene solución**].
15. a) $\boxed{y'' + y' = 2e^x}$ $\mu = 0, -1$. $y_p = Ae^x \rightarrow A + A = 2$. $y = c_1 + c_2e^{-x} + e^x \xrightarrow{d.i.} \begin{cases} y(0) = c_1 + c_2 + 1 = 0 \\ y'(0) = -c_2 + 1 = 2 \end{cases}$, $\boxed{y = e^x - e^{-x}}$.
($\mu = 1$ no autovalor)
O bien: $y' = v \rightarrow v' = -v + 2e^x$, $v = Ce^{-x} + e^{-x} \int 2e^{2x} dx = Ce^{-x} + e^x$. $y = K - Ce^{-x} + e^x$, como antes.
- b) $\boxed{y'' + y' + \lambda y = 0}$ $y'(0) = y'(1) = 0$ $\mu^2 + \mu + \lambda = 0$. $\lambda = 0$, $y = c_1 + c_2e^{-x} \xrightarrow{cc} \begin{cases} -c_2 = 0 \\ -c_2e^{-1} = 0 \end{cases}$, $\forall c_1$. **Autovalor** con $\boxed{y_0 = \{1\}}$.
 $\lambda = -2$, $\mu = 1, -2$, $y = c_1e^x + c_2e^{-2x}$, $y' = c_1e^x - 2c_2e^{-2x} \xrightarrow{cc} \begin{cases} c_1 = 2c_2 \\ c_1e - 2c_2e^{-2} = c_1(e - e^{-2}) = 0, & c_1 = 0 \end{cases}$ **No autovalor**.
[Como la ecuación en forma autoadjunta es $[y'e^x]' + \lambda e^x y = 0$, al ser $q = \alpha\alpha' = \beta\beta' = 0$ sabemos que no había $\lambda < 0$].

16. $y'' + 2y' + y = 4e^x$. a) $\mu^2 + 2\mu + 1 = (\mu + 1)^2 = 0$, $y_p = Ae^x$ (1 no autovalor), $A + 2A + A = 4$, $y = (c_1 + c_2x)e^{-x} + e^x$.

Imponiendo datos $[y' = (c_2 - c_1 - c_2x)e^{-x} + e^x]$: $\begin{cases} c_1 + 1 = 0 \\ c_2 - c_1 + 1 = 2 \end{cases}$, $c_1 = -1$, $c_2 = 0$. $y = e^x - e^{-x}$.

- b) $\lambda = 1 \rightarrow y = (c_1 + c_2x)e^{-x}$ $\begin{cases} y(0) = c_1 = 0 \\ y'(1) = -c_1 e^{-1} = 0 \end{cases}$, $c_1 = 0$, $\forall c_2$. **Autovalor** con autofunción $\{xe^{-x}\}$.

- $\lambda = 2$, $\mu = -1 \pm i \rightarrow y = (c_1 \cos x + c_2 \sin x)e^{-x}$ $\begin{cases} y(0) = c_1 = 0 \\ y'(1) = c_2(\cos 1 - \sin 1)e^{-1} = 0 \end{cases}$, $c_2 = 0$. **No autovalor.**

17. a) $\alpha\alpha' = \beta\beta' = 0$, $q \equiv 0 \Rightarrow \lambda \geq 0$. [Directamente: $\lambda < 0$: $y = c_1 e^{px} + c_2 e^{-px} \rightarrow \begin{cases} c_1 + c_2 = 0 \\ pc_1[e^{p/2} + e^{-p/2}] = 0 \end{cases} \rightarrow c_1 = c_2 = 0$].

$\lambda = 0$: $y = c_1 + c_2x \rightarrow \begin{cases} y(0) = c_1 = 0 \\ y'(\frac{1}{2}) = c_2 = 0 \end{cases} \rightarrow \lambda = 0$ no es autovalor.

$\lambda > 0$: $y = c_1 \cos wx + c_2 \sin wx$, $y' = -wc_1 \sin wx + wc_2 \cos wx$. $y(0) = 0 \rightarrow c_1 = 0 \rightarrow y'(\frac{1}{2}) = wc_2 \cos \frac{w}{2} = 0$

$\rightarrow w_n = (2n-1)\pi$, $\lambda_n = (2n-1)^2\pi^2$, $y_n = \{\sin(2n-1)\pi x\}$, $n = 1, 2, \dots$ (como en formulario).

$\langle y_n, y_n \rangle \stackrel{r=1}{=} \int_0^{1/2} \sin^2 w_n x dx = \frac{1}{2} \int_0^{1/2} [1 - \cos 2w_n x] dx = \frac{1}{4} - \frac{\sin(2n-1)\pi}{2(2n-1)\pi} = \left[\frac{1}{4}\right]$ (como en formulario).

- b) $c_n = 4 \int_0^{1/2} x \sin(2n-1)\pi x dx = -\frac{4x}{(2n-1)\pi} \cos(2n-1)\pi x \Big|_0^{1/2} + \frac{4}{(2n-1)\pi} \int_0^{1/2} \cos(2n-1)\pi x dx = \frac{4}{(2n-1)^2\pi^2} \sin \frac{(2n-1)\pi}{2}$.

Por tanto: $x = \frac{4}{\pi^2} \sum_{n=1}^{\infty} \frac{(-1)^{n+1}}{(2n-1)^2} \sin(2n-1)\pi x = \frac{4}{\pi^2} \left[\sin \pi x - \frac{1}{9} \sin 3\pi x + \frac{1}{25} \sin 5\pi x + \dots \right]$.

18. a) $f(x) = 1$ $b_n = 2 \int_0^1 \sin n\pi x dx = \frac{2[1 - (-1)^n]}{n\pi}$, $\frac{4}{\pi} \sum_{m=1}^{\infty} \frac{\sin(2m-1)\pi x}{2m-1}$.

La serie tiende hacia la extensión 2-periódica de $f(x) = \begin{cases} -1, & -1 < x < 0 \\ 1, & 0 < x < 1 \end{cases}$, y la suma es 0 si $x \in \mathbb{Z}$. Cerca de ellos aparecerán picos.

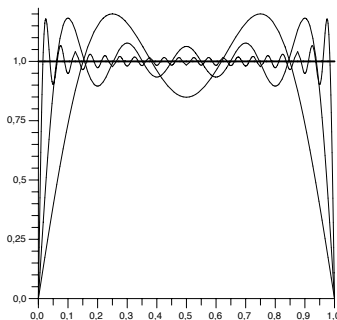
La serie en cosenos es la propia constante $1 = 1 + 0 + 0 + \dots$ (es una de las autofunciones).

- b) $f(x) = x^2$ $b_n = 2 \int_0^1 x^2 \sin n\pi x dx = -\frac{2x^2 \cos n\pi x}{n\pi} \Big|_0^1 + \frac{4}{n\pi} \int_0^1 x \cos n\pi x dx = \frac{2(-1)^{n+1}}{n\pi} + \frac{4x \sin n\pi x}{n^2\pi^2} \Big|_0^1 - \frac{4}{n^2\pi^2} \int_0^1 \sin n\pi x dx$

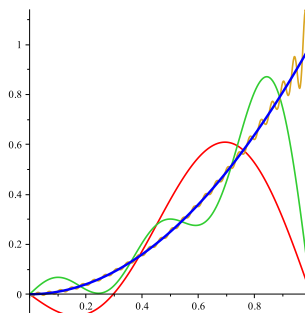
$$a_0 = 2 \int_0^1 x^2 dx = \frac{2}{3}, a_n = 2 \int_0^1 x^2 \cos n\pi x dx = \frac{2x^2 \sin n\pi x}{n\pi} \Big|_0^1 + \frac{4}{n\pi} \int_0^1 x \sin n\pi x dx = -\frac{4x \cos n\pi x}{n^2\pi^2} \Big|_0^1 - \frac{4}{n^2\pi^2} \int_0^1 \cos n\pi x dx$$

Por tanto, $x^2 = \sum_{n=1}^{\infty} \left[\frac{2(-1)^{n+1}}{n\pi} - \frac{4[(-1)^n - 1]}{\pi^3 n^3} \right] \sin n\pi x = \frac{1}{3} - \frac{4}{\pi^2} \sum_{n=1}^{\infty} \frac{(-1)^n}{n^2} \cos n\pi x$.

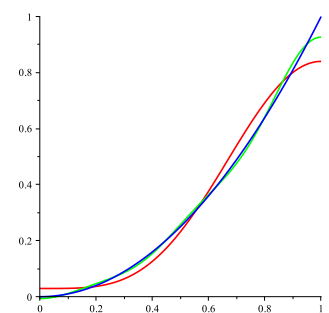
La de senos convergerá mal cerca de 1 y la de cosenos tiende a x^2 en todo $[0, 1]$.



a) sen, $m=2, 5, 20$



b) sen, $n=2, 5, 50$



b) cos, $n=2, 5$

19. $f(x) = \begin{cases} 1, & 0 \leq x \leq 1 \\ 0, & 1 < x \leq 2 \end{cases}$ Si $f(x) = \frac{a_0}{2} + \sum_{n=1}^{\infty} a_n \cos \frac{n\pi x}{L}$, es $a_n = \frac{2}{L} \int_0^L f(x) \cos \frac{n\pi x}{L} dx$. En este caso:

$$a_0 = \frac{2}{2} \int_0^1 dx = 1, a_n = \int_0^1 \cos \frac{n\pi x}{2} dx = \frac{2}{n\pi} \sin \frac{n\pi}{2} = \begin{cases} 0, & n \text{ par} \\ \frac{2(-1)^m}{(2m+1)\pi}, & n = 2m+1 \end{cases} \rightarrow \frac{1}{2} + \frac{2}{\pi} \sum_{m=0}^{\infty} \frac{(-1)^m}{2m+1} \cos \frac{(2m+1)\pi x}{2}$$

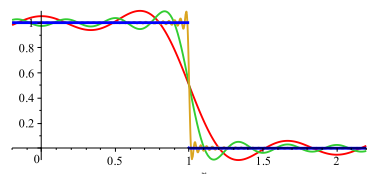
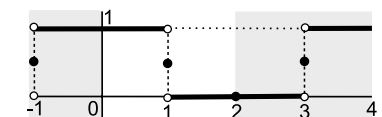
- i) En $x=1$ es f discontinua y la serie tenderá hacia $\frac{1}{2}[f(1^-) + f(1^+)] = \frac{1}{2}$:

$$\frac{1}{2} + \frac{2}{\pi} \sum_{m=0}^{\infty} \frac{(-1)^m}{2m+1} \cos \frac{(2m+1)\pi}{2} = \frac{1}{2} \text{ [los cosenos se anulan].}$$

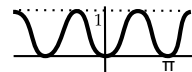
- ii) Como tiende en todo $\mathbb{R}$ hacia la extensión par y 4-periódica de f , en $x=2$ ha de tender hacia $f(2)=0$. Sustituyendo:

$$\frac{1}{2} + \frac{2}{\pi} \sum_{m=0}^{\infty} \frac{(-1)^m}{2m+1} \cos(2m+1)\pi = \frac{1}{2} - \frac{2}{\pi} \sum_{m=0}^{\infty} \frac{(-1)^m}{2m+1} = 0,$$

ya que la última serie $1 - \frac{1}{3} + \frac{1}{5} - \dots = \arctan 1 = \frac{\pi}{4}$.

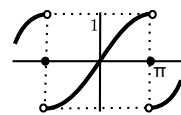


20. a) $f(x) = \sin^2 x = \frac{1}{2} - \frac{1}{2} \cos 2x$, ya desarrollada $[a_0 = \frac{1}{2}, a_2 = -\frac{1}{2}$ y resto de a_n y $b_n = 0]$.



b) $f(x) = \sin \frac{x}{2}$ Como es impar serán $a_n = 0$.

$$b_n = \frac{2}{\pi} \int_0^\pi \sin \frac{x}{2} \sin nx \, dx = \frac{1}{\pi} \int_0^\pi \left[\cos \frac{(1-2n)x}{2} - \cos \frac{(1+2n)x}{2} \right] dx = \frac{2}{\pi} \left[\frac{(-1)^n}{1-2n} - \frac{(-1)^n}{1+2n} \right] = \frac{8}{\pi} \frac{(-1)^n n}{1-4n^2}.$$



21. $\begin{cases} y'' + \lambda y = 0 \\ y(0) - 2y'(0) = y(1) - 2y'(1) = 0 \end{cases}$ Como $\beta \cdot \beta' > 0$ pueden existir $\lambda < 0$.

$$\lambda < 0, y = c_1 e^{px} + c_2 e^{-px} \rightarrow \begin{cases} c_1 + c_2 - 2p[c_1 - c_2] = 0 \\ c_1 e^p + c_2 e^{-p} - 2p[c_1 e^p - c_2 e^{-p}] = 0 \end{cases} \rightarrow \begin{vmatrix} 1-2p & 1+2p \\ [1-2p]e^p & [1+2p]e^{-p} \end{vmatrix} = [1-2p][1+2p][e^{-p} - e^p].$$

El determinante se anula si $p = \frac{1}{2}$ (y entonces es $c_2 = 0$). $\lambda_0 = -\frac{1}{4}$ y su autofunción es $y_0 = \{e^{x/2}\}$.

$\lambda = 0, y = c_1 + c_2 x \rightarrow \begin{cases} c_1 - 2c_2 = 0 \\ c_1 - c_2 = 0 \end{cases} \Rightarrow c_1 = c_2 = 0$. No es autovalor.

$\lambda > 0, y = c_1 \cos wx + c_2 \sin wx \rightarrow \begin{cases} c_1 - 2wc_2 = 0, c_1 = 2wc_2 \\ c_1 \cos w + c_2 \sin w - 2w[-c_1 \sin w + c_2 \cos w] = 0 \end{cases} \rightarrow c_1(1+4w^2) \sin w = 0.$

Por tanto, $w_n = n\pi$, $\lambda_n = n^2 \pi^2$, $y_n = \{2n\pi \cos n\pi x + \sin n\pi x\}$, $n = 1, 2, \dots$

$$\langle y_0, y_0 \rangle = \int_0^1 e^x dx = [e - 1]. \langle y_n, y_n \rangle = \int_0^1 [2n^2 \pi^2 (1 + \cos 2n\pi x) + 2n\pi \sin 2n\pi x + \frac{1}{2}(1 - \cos 2n\pi x)] dx = \left[\frac{1}{2} + 2n^2 \pi^2 \right].$$

22. $\begin{cases} y'' + \lambda y = 0 \\ y(0) = y(1) - y'(1) = 0 \end{cases}$ a) $\lambda = 0: y = c_1 + c_2 x, \begin{cases} c_1 = 0 \\ c_1 + c_2 - c_2 = 0 \end{cases} \rightarrow c_1 = 0, \forall c_2 \rightarrow \lambda_0 = 0, y_0 = \{x\}$.

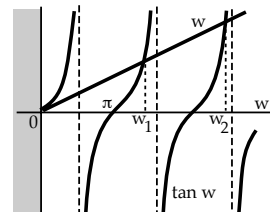
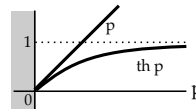
$\lambda > 0: y = c_1 \cos wx + c_2 \sin wx, w = \sqrt{\lambda} \rightarrow \begin{cases} c_1 = 0 \\ c_2(\sin w - w \cos w) = 0 \end{cases}$

c_2 cualquiera si $\sin w = w \cos w$. Infinitos $\lambda_n = w_n^2$ con $w_n = \tan w_n$, $y_n = \{\sin w_n x\}$.

Como $\beta\beta' < 0$ podría haber $\lambda < 0$. Para ver que no buscamos la tangente hiperbólica:

$$y = c_1 e^{px} + c_2 e^{-px} \rightarrow \begin{cases} c_1 = -c_2 \\ c_2(p[e^p + e^{-p}] - [e^p - e^{-p}]) = 0 \end{cases} \rightarrow y \equiv 0$$

[no existe $p > 0$ con $p = \text{th } p$, pues $(\text{th } p)'(0) = 1]$.



b) Los coeficientes del desarrollo $1 = \sum_{n=0}^{\infty} c_n y_n(x) = c_0 x + \sum_{n=1}^{\infty} \sin w_n x$ vienen dados por $c_n = \frac{\langle 1, y_n \rangle}{\langle y_n, y_n \rangle}$.

En particular, $\langle 1, x \rangle = \int_0^1 x dx = \frac{1}{2}$, $\langle x, x \rangle = \int_0^1 x^2 dx = \frac{1}{3}$. Por tanto, $c_0 = \frac{1/2}{1/3} = \frac{3}{2}$, $1 = \frac{3}{2}x + \dots$

23. $\begin{cases} y'' + 2y' + \lambda y = 0 \\ y(0) + y'(0) = y(1/2) = 0 \end{cases}$ En forma autoadjunta: $(y'e^{2x})' + \lambda e^{2x} y = 0$ [problema de S-L regular].

$\mu^2 + 2\mu + \lambda = 0 \rightarrow \mu = -1 \pm \sqrt{1 - \lambda}$. En principio, puede haber λ negativos.

$\lambda < 1, \sqrt{1 - \lambda} = p \rightarrow y = c_1 e^{(p-1)x} + c_2 e^{-(p+1)x} \rightarrow \begin{cases} y(0) + y'(0) = p(c_1 - c_2) = 0 \\ y(1/2) = (c_1 e^{p/2} + c_2 e^{-p/2}) e^{-1/2} = 0 \end{cases} \rightarrow c_1 = c_2 = 0.$

$\lambda = 1 \rightarrow y = (c_1 + c_2 x) e^{-x} \rightarrow \begin{cases} y(0) + y'(0) = c_2 = 0 \\ y(1/2) = (c_1 + \frac{1}{2}c_2) e^{-1/2} = 0 \end{cases} \rightarrow c_1 = c_2 = 0.$

$\lambda > 1, \sqrt{\lambda - 1} = w \rightarrow y = (c_1 \cos wx + c_2 \sin wx) e^{-x} \rightarrow y(0) + y'(0) = c_2 w = 0, c_2 = 0 \rightarrow y(1/2) = c_1 \cos \frac{w}{2} e^{-1/2} = 0$
 $\rightarrow w_n = (2n-1)\pi, \lambda_n = 1 + (2n-1)^2 \pi^2, y_n = \{e^{-x} \cos(2n-1)\pi x\}, n = 1, 2, \dots$

Por tanto: $1 = \sum_{n=1}^{\infty} \frac{\langle 1, y_n \rangle}{\langle y_n, y_n \rangle} y_n$, con $\langle y_n, y_n \rangle = \int_0^{1/2} \cos^2(2n-1)\pi x dx = \frac{1}{4}$, $\langle 1, y_n \rangle = \int_0^{1/2} e^x \cos(2n-1)\pi x dx$.

Como es $\int e^x \cos bx dx = \frac{(\cos bx + b \sin bx) e^x}{1+b^2}$, concluimos que: $1 = 4 \sum_{n=1}^{\infty} \frac{(-1)^{n+1} (2n-1) \pi e^{1/2-1}}{1+(2n-1)^2 \pi^2} e^{-x} \cos(2n-1)\pi x$.

24. $\begin{cases} x^2 y'' + 3xy' + y + \lambda y = 0 \\ y(1) = y(e) = 0 \end{cases}$ a) $y'' + \frac{3}{x} y' + \frac{1}{x^2} y + \frac{\lambda}{x^2} y = 0, e^{3/x} = x^3, (x^3 y')' + xy + \lambda xy = 0 \Rightarrow r(x) = x$ peso.

Ecuación de Euler: $\mu(\mu-1) + 3\mu + 1 + \lambda = 0, \mu = -1 \pm \sqrt{-\lambda}$.

b) Para $\lambda = 0, \mu = -1$ doble $\rightarrow y = (c_1 + c_2 \ln x) x^{-1} \rightarrow \begin{cases} y(1) = c_1 = 0 \\ y(e) = (c_1 + c_2) e^{-1} = 0 \end{cases} \Rightarrow c_1 = c_2 = 0$. No es autovalor.

c) Para $\lambda = \pi^2, \mu = -1 \pm i\pi$ la solución general es: $y = c_1 \frac{\cos(\pi \ln x)}{x} + c_2 \frac{\sin(\pi \ln x)}{x} \rightarrow \begin{cases} y(1) = c_1 = 0 \\ y(e) = -c_1 e^{-1} = 0 \end{cases}, \forall c_2$.

Es pues, $\lambda = \pi^2$ autovalor con **autofunción** asociada $y_1 = \left\{ \frac{\sin(\pi \ln x)}{x} \right\}$.

$$\langle y_1, y_1 \rangle = \int_1^e x \frac{1}{x^2} \sin^2(\pi \ln x) dx = [\ln x = s] = \int_0^1 \sin^2(\pi s) ds = \frac{1}{2} \int_0^1 [1 - \cos(2\pi s)] ds = \left[\frac{1}{2} \right].$$

25. $y'' + y' = x$ $\mu^2 + \mu = 0$, $\mu = 0, -1$. $y_p = Ax^2 + Bx \rightarrow 2A + 2Ax + B = x$, $A = \frac{1}{2}$, $B = -1$. $y = c_1 + c_2 e^{-x} + \frac{1}{2}x^2 - x$.

Imponiendo datos: $\begin{cases} c_1 + c_2 = -1 \\ -c_2 - 1 = -1 \end{cases}$, $c_1 = -1$, $c_2 = 0 \rightarrow y = \frac{1}{2}x^2 - x - 1$.

b) Mejor mirar el homogéneo: $y = c_1 + c_2 e^{-x} \xrightarrow{cc} \begin{cases} -c_2 = 0 \\ -c_2 e^{-1} = 0 \end{cases}$. Infinitas soluciones $y_h = \{1\}$ del homogéneo $\Rightarrow$ infinitas o ninguna del no homogéneo.

[Imponiendo los datos en la no homogénea o haciendo la integral del teorema se ve que no tiene solución].

26. $y'' + y' + \lambda y = 0$ Multiplicando por e^x : $(e^x y')' + \lambda e^x y = 0$ forma autoadjunta.
 $y'(0) = y'(1) = 0$ $\alpha\alpha' = \beta\beta' = q = 0 \Rightarrow$ los $\lambda \geq 0$ y $\lambda = -2$ no puede ser autovalor. O directamente:

$\mu^2 + \mu - 2 = 0$, $\mu = 1, -2$, $y = c_1 e^x + c_2 e^{-2x}$, $y' = c_1 e^x - 2c_2 e^{-2x}$, $\begin{cases} y'(0) = c_1 - 2c_2 = 0 \\ y'(1) = c_1 e - 2c_2 e^{-2} = 0 \end{cases}$, $\begin{vmatrix} 1 & -2 \\ e & -2e^{-2} \end{vmatrix} \neq 0 \rightarrow c_1 = c_2 = 0$.

$\mu^2 + \mu = 0$, $\mu = 0, -1$, $y = c_1 + c_2 e^{-x}$, $y' = -c_2 e^{-x}$, $\begin{cases} y'(0) = -c_2 = 0 \\ y'(1) = -c_2 e^{-1} = 0 \end{cases} \rightarrow c_2 = 0, \forall c_1$. $\lambda = 0$ autovalor [$y_0 = \{1\}$].

Como para $\lambda = -2$ el homogéneo tiene sólo la solución $y \equiv 0$, el no homogéneo tiene **solución única**.

27. $x^2 y'' + \lambda y = 0$ a) Euler, $\mu^2 - \mu - \lambda = 0$. $\lambda = 0 \rightarrow \mu = 0, 1$, $y = c_1 + c_2 x \xrightarrow{cc} \begin{cases} c_2 = 0 \\ c_2 = 0 \end{cases}$, $\forall c_1$. **Autovalor** con $y_0 = \{1\}$.
 $y'(1) = y'(2) = 0$

$\lambda = -2 \rightarrow \mu = 2, -1$, $y = c_1 x^2 + \frac{c_2}{x}$, $y' = 2c_1 x - \frac{c_2}{x^2} \xrightarrow{cc} \begin{cases} y'(1) = 2c_1 - c_2 = 0 \\ y'(2) = 2c_1 - \frac{1}{4}c_2 = 0 \end{cases} \Rightarrow c_1 = c_2 = 0$. **No es autovalor**.

[Como la ecuación es $[y']' + \lambda \frac{1}{x^2} y = 0$ en forma autoadjunta, al ser $q = \alpha\alpha' = \beta\beta' = 0$ sabemos que no había $\lambda < 0$].

b) Como el homogéneo tenía sólo la solución $y \equiv 0$, este no homogéneo tendrá seguro **solución única**.

[Imponiendo los datos en la solución de la no homogénea $y = c_1 x^2 + c_2 x^{-1} - 1$, se ve que esa única solución es $y = -1$].

28. $xy'' - y' = x^2 - a$ La homogénea se puede resolver como Euler: $\lambda(\lambda - 1) - \lambda = 0 \rightarrow y = c_1 + c_2 x^2$,
 $y'(2) = y'(4) = 0$ o haciendo $y' = v \rightarrow v' = \frac{v}{x} \rightarrow v = C e^{\ln x} = Cx \rightarrow y = c_1 + c_2 x^2$ [$y' = 2c_2 x$].

Imponiendo datos: $\begin{cases} y'(2) = 4c_2 = 0 \\ y'(4) = 8c_2 = 0 \end{cases} \rightarrow c_2 = 0$ y todo c_1 . El homogéneo tiene infinitas soluciones $y_h = \{1\}$.

El no homogéneo tendrá infinitas o ninguna. En forma autoadjunta: $y'' - \frac{1}{x} y' = x - \frac{a}{x} \xrightarrow{\times e^{-\log x}} (\frac{1}{x} y')' = 1 - \frac{a}{x^2}$.

$\int_2^4 1 \cdot (1 - \frac{a}{x^2}) dx = 2 + [\frac{a}{x}]_2^4 = 2 - \frac{a}{4} \rightarrow$ Si $a = 8$ tiene infinitas soluciones. [Si $a \neq 8$, ninguna].

[Se llega a lo mismo imponiendo los datos en la solución general de la no homogénea $\dots y = c_1 + c_2 x^2 + \frac{1}{3}x^3 + ax \dots$].

29. $y'' + \lambda y = \sin x$ La ecuación ya está en forma autoadjunta: $[y']' + \lambda y = \sin x$.
 $y(0) = y'(\frac{\pi}{2}) = 0$ Sabemos que para el homogéneo son: $\lambda_n = (2n-1)^2$, $y_n = \{\sin(2n-1)x\}$, $n = 1, 2, \dots$

Para cualquier $\lambda \neq (2n-1)^2$ el homogéneo sólo tiene la solución trivial y el no homogéneo **solución única**.

Si $\lambda = (2n-1)^2$ el no homogéneo tendrá infinitas o ninguna según sea $\neq 0$ la integral $I = \int_0^{\pi/2} \sin x \sin(2n-1)x dx$.

Por tanto, **no hay solución** sólo si $\lambda = 1$ [$\int_0^{\pi/2} \sin^2 x dx \neq 0$], ya que para los otros autovalores $\lambda = 9, 25, \dots$ la integral es cero [sin calcularla: $\sin x$ es ortogonal a las otras autofunciones] y hay **infinitas soluciones**.

[La integral no es difícil de calcular: $I = \frac{1}{2} \int_0^{\pi/2} [\cos 2(n-1)x - \cos 2nx] dx \stackrel{n \neq 1}{=} \frac{\sin(n-1)\pi}{4(n-1)} - \frac{\sin n\pi}{4n} = 0$, $n \neq 1$].

30. $y'' + \lambda y = \cos 3x$ $\alpha \cdot \alpha' = 0 \Rightarrow \lambda \geq 0$. $\lambda = 0$: $y = c_1 + c_2 x \rightarrow \begin{cases} y'(0) = c_2 = 0 \\ y'(\frac{\pi}{4}) + y(\frac{\pi}{4}) = c_1 = 0 \end{cases}$ $\lambda = 0$ no autovalor.
 $y'(0) = y'(\frac{\pi}{4}) + y(\frac{\pi}{4}) = 0$

$\lambda > 0$: $y = c_1 \cos wx + c_2 \sin wx$. $y'(0) = 0 \rightarrow c_2 = 0 \rightarrow y'(\frac{\pi}{4}) + y(\frac{\pi}{4}) = c_1 [\cos \frac{w\pi}{4} - w \sin \frac{w\pi}{4}] = 0$.

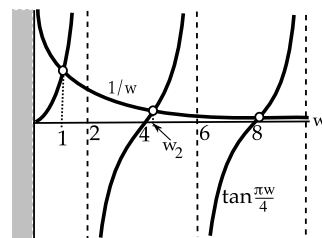
Si el corchete es cero, c_1 queda indeterminado. Infinitos w_n cumplen $\tan \frac{\pi w_n}{4} = \frac{1}{w_n}$. A cada $\lambda_n = w_n^2$, está asociada la $y_n = \{\cos w_n x\}$.

λ_1 se pueda hallar exactamente: $[\lambda_1 = 1 \rightarrow y_1 = \{\cos x\}]$ ($\tan \frac{\pi}{4} = 1$).

Si $\cos 3x = c_1 \cos x + \sum_{n=2}^{\infty} c_n \cos w_n x$, el primer coeficiente c_1 es:

$c_1 = \frac{\langle \cos 3x, \cos x \rangle}{\langle \cos x, \cos x \rangle} = \frac{\int_0^{\pi/4} \cos 3x \cos x dx}{\int_0^{\pi/4} \cos^2 x dx} = \frac{\int_0^{\pi/4} (\cos 4x + \cos 2x) dx}{\int_0^{\pi/4} (1 + \cos 2x) dx} = \frac{2}{\pi + 2}$.

Para i), por no ser $\lambda = 0$ autovalor hay solución única del no homogéneo. Para ii), hay infinitas del homogéneo $y_h = \{\cos x\}$ y el no homogéneo no tiene solución pues: $\int_0^{\pi/4} \cos 3x \cos x dx = \frac{1}{4} \neq 0$.



Soluciones de problemas 4 de MM(im) (2020)

1. i) $x^2 y'' - 2y = 0$ es de Euler. $\mu(\mu-1)-2=(\mu-2)(\mu+1)=0 \rightarrow y = c_1 x^2 + c_2 x^{-1}$, c_1, c_2 constantes.

ii) $x^2 u_x = 2u$ es EDP lineal de primer orden en x con una constante para cada y :

$$u(x, y) = p(y) e^{\int (2/y^2) dx} = p(y) e^{-2/x}, \text{ con } p \text{ función arbitraria.}$$

2. $\begin{cases} 2xu_y + u_x = 4xy \\ u(1, y) = 1 \end{cases} \quad \frac{dy}{dx} = 2x. \quad y = x^2 + C. \text{ Características: } y - x^2 = C. \quad [\text{También válidas } x^2 - y = C, \dots].$

$$\text{Haciendo } \begin{cases} \xi = y - x^2 \\ \eta = y \end{cases} \rightarrow \begin{cases} u_y = u_\xi + u_\eta \\ u_x = -2xu_\xi \end{cases} \rightarrow 2xu_\eta = 4xy, \quad u_\eta = 2y = 2\eta. \quad u = p(\xi) + \eta^2 = p(y - x^2) + y^2.$$

$$[\text{Más largo: } \begin{cases} \xi = y - x^2 \\ \eta = x \end{cases} \rightarrow \begin{cases} u_y = u_\xi \\ u_x = -2xu_\xi + u_\eta \end{cases} \rightarrow u_\eta = 4xy = 4\xi\eta - 4\eta^3. \quad u = q(\xi) + 2\xi\eta - \eta^3 = q(y - x^2) + 2x^2y - x^4.]$$

$$\text{Imponiendo el dato inicial: } u(1, y) = p(y-1) + y^2 = 1 \rightarrow p(v) = -2v - v^2, \quad u(x, y) = 2x^2y - 2y + 2x^2 - x^4. \\ [\text{O bien, } q(y-1) + 2y - 1 = 1, \quad q(v) = -2v - v^3.]$$

$$\text{Comprobamos: } u(1, y) = 1, \text{ y además: } 2xu_y + u_x = 2x(2x^2 - 2) + (4xy + 4x - 4x^3) = 4xy.$$

3. $u_y - 2u_x = (x+2y)u \quad \frac{dy}{dx} = -\frac{1}{2}. \quad \int 2 dy = -\int dx + C. \text{ Características: } x+2y = C.$

$$\text{Haciendo } \begin{cases} \xi = x+2y \\ \eta = y \end{cases} \rightarrow \begin{cases} u_y = 2u_\xi + u_\eta \\ u_x = u_\xi \end{cases} \rightarrow u_\eta = (x+2y)u = \xi u. \quad u = p(\xi) e^{\xi\eta} = p(x+2y) e^{xy+2y^2}.$$

$$[\text{O más largo: } \begin{cases} \xi = x+2y \\ \eta = x \end{cases} \rightarrow \begin{cases} u_y = 2u_\xi \\ u_x = u_\xi + u_\eta \end{cases} \rightarrow -2u_\eta = (x+2y)u, \quad u_\eta = -\frac{\xi}{2}u. \quad u = q(\xi) e^{-\xi\eta/2} = q(x+2y) e^{-x^2/2 - xy}.]$$

$$\text{Imponiendo el dato inicial: } u(x, 1) = p(x+2) e^{x+2} = 2, \quad p(v) = 2e^{-v}, \quad u(x, y) = 2 e^{xy-x-2y+2y^2} = 2 e^{(y-1)(x+2y)}.$$

$$\text{Comprobamos: } u(x, 1) = 1, \text{ y además: } u_y - 2u_x = (x-2+4y-2y+2) e^{\dots} = (x+2y) e^{\dots}.$$

[Como los datos se dan sobre una recta no característica, la solución debía ser única].

4. $u_y + u_x = u + x \quad \frac{dy}{dx} = \frac{1}{1} \rightarrow y - x = C. \text{ Más corto: } \begin{cases} \xi = y - x \\ \eta = x \end{cases} \rightarrow \begin{cases} u_y = u_\xi \\ u_x = -u_\xi + u_\eta \end{cases} \rightarrow u_\eta = u + x = u + \eta.$

$$u(\xi, \eta) = p(\xi) e^\eta + e^\eta \int e^{-\eta} \eta d\eta = p(\xi) e^\eta - \eta - 1 \quad [\text{o probando } u_p = A\eta + B]. \quad u(x, y) = p(y-x) e^x - x - 1.$$

$$[\text{Algo más largo: } \begin{cases} \xi = y - x \\ \eta = y \end{cases} \rightarrow \begin{cases} u_y = u_\xi + u_\eta \\ u_x = -u_\xi \end{cases} \rightarrow u_\eta = u + x = u + \eta - \xi. \quad u = p(\xi) e^\eta + e^\eta \int e^{-\eta} (\eta - \xi) d\eta = p(\xi) e^\eta + \xi - \eta - 1.]$$

$$\text{Imponiendo el dato inicial: } u(x, 0) = p(-x) e^x - x - 1 = -x \rightarrow p(-x) = e^{-x}, \quad p(v) = e^v, \quad u(x, y) = e^y - x - 1.$$

$$[\text{Comprobamos: } u(x, 0) = -x, \text{ y además: } u_y + u_x = e^y - 1 = e^y - x - 1 + x = u + x.]$$

5. $2xyu_y - u_x = 2xy \quad \frac{dy}{dx} = -2xy \text{ lineal. } y = Ce^{-\int 2x} = Ce^{-x^2}. \text{ Características: } y e^{x^2} = C.$

$$\text{Haciendo } \begin{cases} \xi = y e^{x^2} \\ \eta = y \end{cases} \rightarrow \begin{cases} u_y = e^{x^2} u_\xi + u_\eta \\ u_x = 2xy e^{x^2} u_\xi \end{cases} \rightarrow 2xyu_\eta = 2xy, \quad u_\eta = 1, \quad u = \eta + p(\xi) = y + p(y e^{x^2}).$$

$$[\text{Peor: } \begin{cases} \xi = y e^{x^2} \\ \eta = x \end{cases} \rightarrow \begin{cases} u_y = e^{x^2} u_\xi \\ u_x = 2xy e^{x^2} u_\xi + u_\eta \end{cases} \rightarrow -u_\eta = 2xy, \quad u_\eta = -2\xi\eta e^{-\eta^2}, \quad u = \xi e^{-\eta^2} + p(\xi) = y + p(y e^{x^2}) \text{ como antes}.]$$

$$\text{Imponiendo el dato inicial: } u(1, y) = y + p(ey) = 0 \rightarrow p(v) = -v/e, \quad u(x, y) = y - y e^{x^2-1}.$$

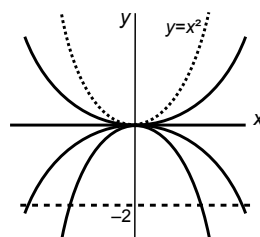
6. $2yu_y + xu_x = 4x^2y \quad \frac{dy}{dx} = \frac{2y}{x} \rightarrow y = C e^{\int (2/x)} = C e^{2\ln x} = Cx^2. \quad \begin{cases} \xi = yx^{-2} \\ \eta = y \end{cases} \rightarrow \begin{cases} u_y = x^{-2} u_\xi + u_\eta \\ u_x = -2yx^{-3} u_\xi \end{cases}$

$$2yu_\eta = 4x^2y, \quad u_\eta = 2x^2 = \frac{2\eta}{\xi} \rightarrow u = \frac{\eta^2}{\xi} + p(\xi). \quad u(x, y) = x^2y + p\left(\frac{y}{x^2}\right) \quad \text{solución general}$$

$$\text{O bien } \begin{cases} \xi = yx^{-2} \\ \eta = x \end{cases} \rightarrow xu_\eta = 4x^2y, \quad u_\eta = 4xy = 4\xi\eta^3 \rightarrow u = \xi\eta^4 + p(\xi) = x^2y + p\left(\frac{y}{x^2}\right).$$

$$\text{Imponiendo el dato: } 4y + p\left(\frac{y}{4}\right) = 3y, \quad p\left(\frac{y}{4}\right) = -y, \quad p(v) = -4v, \quad u(x, y) = x^2y - \frac{4y}{x^2}.$$

[Solución única por no ser tangente $x = -2$ a las características].



7. a) $\boxed{(2y-x)u_y + xu_x = 2y}$ $\frac{dy}{dx} = \frac{2y}{x} - 1 \rightarrow y = Cx^2 + x \rightarrow \begin{cases} \xi = \frac{y}{x^2} - \frac{1}{x} \\ \eta = x \end{cases} \rightarrow xu_\eta = 2y, u_\eta = 2 + 2\xi\eta \rightarrow$
 $u = 2\eta + \xi\eta^2 + p(\xi) = x + y + p(\frac{y}{x^2} - \frac{1}{x}) \rightarrow 1 + y + p(y-1) = 0 \rightarrow p(v) = -v - 2, \boxed{u = \frac{1}{x} - \frac{y}{x^2} + x + y - 2}.$

b) $\boxed{u_y + 3y^2u_x = \frac{2u}{y} + 6y^4x}$ $\begin{cases} \xi = x - y^3 \\ \eta = y \end{cases} \rightarrow u_\eta = \frac{2u}{\eta} + 6\eta^4(\xi + \eta^3), u = p(\xi)\eta^2 + \eta^2(\xi + \eta^3)^2 = p(x - y^3)y^2 + x^2y^2$
 O bien, $\begin{cases} \xi = x - y^3 \\ \eta = x \end{cases} \rightarrow u_\eta = \frac{2u}{3(\eta - \xi)} + 2(\eta - \xi)^{2/3}\eta, u = p(\xi)(\eta - \xi)^{2/3} + \eta^2(\eta - \xi)^{2/3}$
 $u(x, 1) = p(x - 1) + x^2 = x^2 \rightarrow p(v) = 0 \rightarrow \boxed{u = x^2y^2}$

c) $\boxed{u_y + 2yu_x = 3xu}$ $\frac{dy}{dx} = \frac{1}{2y} \cdot \int 2y dy = \int dx + C. \boxed{y^2 - x = C}.$ [También se podría haber puesto $x - y^2 = C, \dots$].

Haciendo $\begin{cases} \xi = y^2 - x \\ \eta = y \end{cases} \rightarrow \begin{cases} u_y = 2yu_\xi + u_\eta \\ u_x = -u_\xi \end{cases} \rightarrow u_\eta = 3xu = (3\eta^2 - 3\xi)u, u = p(\xi)e^{\eta^3 - 3\xi\eta} = \boxed{p(y^2 - x)e^{3xy - 2y^3}}.$

[Peor con: $\begin{cases} \xi = y^2 - x \\ \eta = x \end{cases} \rightarrow \begin{cases} u_y = 2yu_\xi \\ u_x = -u_\xi + u_\eta \end{cases} \rightarrow u_\eta = \frac{3x}{2y}u = \frac{3\eta}{2\sqrt{\xi + \eta}}u$, con la integral bastante más larga de calcular].

Imponiendo el dato inicial: $u(x, 0) = p(-x) = 2x \rightarrow p(v) = -2v, u(x, y) = \boxed{2(x - y^2)e^{3xy - 2y^3}}.$

d) $\boxed{2yu_y - xu_x = 2u}$ $\frac{dy}{dx} = -\frac{2y}{x}$ lineal. $y = C e^{\int (-2/x) dx} = C e^{\ln x^{-2}} = \frac{C}{x^2}.$ Características: $\boxed{x^2y = C}.$

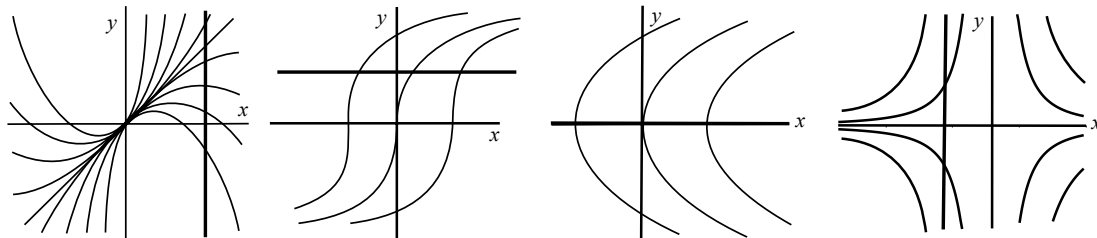
$\begin{cases} \xi = x^2y \\ \eta = y \end{cases} \rightarrow \begin{cases} u_y = x^2u_\xi + u_\eta \\ u_x = 2xyu_\xi \end{cases} \rightarrow 2yu_\eta = 2u, u_\eta = \frac{1}{\eta}u, u = p(\xi)e^{\ln \eta} = p(\xi)\eta. \boxed{u(x, y) = p(x^2y)y}.$

O bien: $\begin{cases} \xi = x^2y \\ \eta = x \end{cases} \rightarrow \begin{cases} u_y = x^2u_\xi \\ u_x = 2xyu_\xi + u_\eta \end{cases} \rightarrow -xu_\eta = 2u, u_\eta = -\frac{2}{\eta}u, u = q(\xi)e^{-2\ln \eta} = \frac{q(\xi)}{\eta^2}. \boxed{u(x, y) = \frac{q(x^2y)}{x^2y}}.$

Imponiendo el dato inicial: $u(-1, y) = p(y)y = y^3, p(v) = v^2, u(x, y) = x^4y^2y = \boxed{x^4y^3}.$

O bien: $u(-1, y) = q(y) = y^3, q(v) = v^3, u(x, y) = x^6y^3/x^2 = \boxed{x^4y^3}.$

[Comprobamos: $u(-1, y) = y^3$, y además: $2yu_y - xu_x = 6x^4y^3 - 4x^4y^3 = 2u$. La solución debía ser única].



8. a) $\boxed{\frac{dy}{dx} = \frac{2x-y}{x}}$ i) Como lineal: $\frac{dy}{dx} = -\frac{y}{x} + 2 \rightarrow y = \frac{C}{x} + \frac{1}{x} \int 2x dx = \frac{C}{x} + x, x(y-x) = C.$

ii) $(y-2x) + x \frac{dy}{dx} = 0, M_y = 1 \equiv N_x. \begin{cases} U_x = y - 2x \\ U_y = x \end{cases} \rightarrow \begin{cases} U = xy - x^2 + p(y) \\ U = xy + q(x) \end{cases} \rightarrow xy - x^2 = C \uparrow$

iii) $y = xz \rightarrow xz' + z = 2 - z. \int \frac{dz}{z-1} = -\int \frac{2dx}{x} + C, \ln(z-1) = C - 2\ln x, z-1 = \frac{y}{x} - 1 = Cx^{-2}.$

b) $\boxed{(2x-y)u_y + xu_x = yu}$ $\begin{cases} \xi = xy - x^2 \\ \eta = x \end{cases} \rightarrow u_\eta = \frac{y}{x}u = (1 + \frac{\xi}{\eta^2})u, u = p(\xi)e^{\eta - \frac{\xi}{\eta}} = p(xy - x^2)e^{2x-y}.$

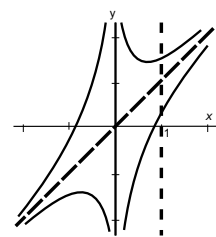
$u(1, y) = p(y-1)e^{2-y} = 1, p(v) = e^{v-1}, \boxed{u(x, y) = e^{xy-x^2+2x-y-1} = e^{(x-1)(y-x+1)}} [Solución única por no ser $x=1$ tangente a las características].$

[Más largo sería: $\begin{cases} \xi = xy - x^2 \\ \eta = y \end{cases}, \eta^2 - 4\xi = (2x-y)^2 \rightarrow u_\eta = \frac{y}{2x-y}u = \frac{\eta}{\sqrt{\eta^2 - 4\xi}}u \rightarrow u = p(\xi)e^{\sqrt{\eta^2 - 4\xi}} = p(xy - x^2)e^{2x-y}.$

Comprobando: $u(1, y) = e^0 = 1. u_y = (x-1)e^{\dots}, u_x = (y-2x+2)e^{\dots}, (2x^2 - 2x - xy + y + xy - 2x^2 + 2x)e^{\dots} = ye^{\dots}.$

c) $u(x, x) = p(0)e^x = e^x \rightarrow p(0) = 1.$ Para cada p derivable que lo cumpla tenemos una solución. Hay infinitas.

[Dato sobre una característica]. Por ejemplo, tomando $p(v) \equiv 1$ y $p(v) = e^v$ es: $u = e^{2x-y}$ y $u = e^{xy-x^2+2x-y}.$



9. $\boxed{u_{xx} + 4u_{xy} - 5u_{yy} + 6u_x + 3u_y = 9u}$ Hiperbólica $\begin{cases} \xi = x - \frac{y}{5} \\ \eta = x + y \end{cases}$ ó $\begin{cases} \xi = 5x - y \\ \eta = x + y \end{cases} \rightarrow 4u_{\xi\eta} + 3u_{\xi\xi} + u_{\eta\eta} = u$ no resoluble.

$\boxed{u_{yy} + 2u_{xy} + 2u_{xx} = 0}$ $B^2 - 4AC = -4$ elíptica $\begin{cases} \xi = x - y \\ \eta = y \end{cases} \begin{cases} u_{xx} = u_{\xi\xi} \\ u_{xy} = -u_{\xi\xi} + u_{\xi\eta} \\ u_{yy} = u_{\xi\xi} - 2u_{\xi\eta} + u_{\eta\eta} \end{cases} \rightarrow u_{\xi\xi} + u_{\eta\eta} = 0$ no resoluble

$\boxed{u_{xx} - 3u_x + 2u = y}$ Parabólica en forma normal. $\lambda^2 - 3\mu + 2 = 0 \rightarrow \mu = 1, 2 \rightarrow \boxed{u = p(y)e^x + q(y)e^{2x} + \frac{y}{2}}$

10. a) $y''+y=3$. $\mu^2+1=0$, $\mu=\pm i$. $y_p=3$ a ojo ($y_p=A \rightarrow A=3$). Solución $y=c_1 \cos x+c_2 \sin x+3$.

b) $u_{yy}-2u_{xy}+u_{xx}+u=x+y$ $B^2-4AC=0$ parabólica $\begin{cases} \xi=x+y \\ \eta=y \end{cases} \rightarrow \begin{cases} u_y=u_\xi+u_\eta \\ u_x=u_\xi \end{cases}, \begin{cases} u_{yy}=u_{\xi\xi}+2u_{\xi\eta}+u_{\eta\eta} \\ u_{xy}=u_{\xi\xi}+u_{\xi\eta} \\ u_{xx}=u_{\xi\xi} \end{cases} \rightarrow u_{\eta\eta}+u=x+y$
 $u_{\eta\eta}+u=\xi$ forma canónica. $\mu=\pm i$, $u_p=\xi \rightarrow u=p(\xi) \cos \eta+q(\xi) \sin \eta+\xi=p(x+y) \cos y+q(x+y) \sin y+x+y$ solución general

11. a) $u_{tt}-4u_{xx}+2u_t+4u_x=0$. $B^2-4AC=16$. Hiperbólica. En el formulario: $-\frac{B \pm \sqrt{B^2-4AC}}{2A} = \pm 2$, $\begin{cases} \xi=x+2t \\ \eta=x-2t \end{cases}$
(O directamente son las de la ecuación de ondas, pues sólo dependen de las derivadas de segundo orden).

Haciendo el cambio: $\begin{cases} u_t=2u_\xi-2u_\eta \\ u_x=u_\xi+u_\eta \end{cases}, \begin{cases} u_{tt}=4u_{\xi\xi}-8u_{\xi\eta}+4u_{\eta\eta} \\ u_{xx}=u_{\xi\xi}+2u_{\xi\eta}+u_{\eta\eta} \end{cases} \rightarrow -16u_{\xi\eta}+8u_\xi=0 \xrightarrow{u_\xi=v} v_{\eta\eta}=\frac{1}{2}v \rightarrow$

$v=p^*(\xi) e^{\eta/2} \rightarrow u=q(\eta)+p(\xi) e^{\eta/2} \rightarrow u=q(x-2t)+p(x+2t) e^{x/2-2t}$, q, p funciones C^2 arbitrarias.

$\begin{cases} q(x)+p(x)e^{x/2}=x, & q'(x)+[p'(x)+\frac{p(x)}{2}]e^{x/2}=1 \\ -2q'(x)+[2p'(x)-p(x)]e^{x/2}=0 \end{cases} \rightarrow 4p'(x)e^{x/2}=2, p(x)=-e^{-x/2}, q(x)=x+1, u=x-2t+1-e^{-2t}$

b) $u_{tt}+2u_{xt}=2$ Hiperbólica $\begin{cases} \xi=x-2t \\ \eta=x \end{cases} \rightarrow \begin{cases} u_{tt}=4u_{\xi\xi} \\ u_{tx}=-2u_{\xi\xi}-2u_{\xi\eta} \end{cases} \rightarrow u_{\xi\eta}=-\frac{1}{2}, u=p(\xi)+q(\eta)-\frac{\xi\eta}{2}$
 $u=p(x-2t)+q(x)-\frac{1}{2}x^2+xt$ $\begin{cases} 0=u(x,0)=p(x)+q(x)-\frac{1}{2}x^2 \\ 0=u_t(x,0)=-2p'(x)+x \end{cases} \rightarrow \begin{cases} q(x)=\frac{1}{4}x^2-C \\ p(x)=\frac{1}{4}x^2+C \end{cases} \rightarrow u=t^2$.
 $y=0$ no tangente a las características $\rightarrow$ solución única

c) $u_{yy}+2u_{xy}+u_{xx}-u=y$. $B^2-4AC=0$ parabólica, $\begin{cases} \xi=x-y \\ \eta=y \end{cases} \rightarrow \begin{cases} u_x=u_\xi \\ u_y=-u_\xi+u_\eta \end{cases}, \begin{cases} u_{xx}=u_{\xi\xi} \\ u_{xy}=-u_{\xi\xi}+u_{\xi\eta} \\ u_{yy}=u_{\xi\xi}-2u_{\xi\eta}+u_{\eta\eta} \end{cases}$
 $\rightarrow u_{\eta\eta}-u=\eta$ forma canónica $\mu=\pm 1$, $u_p=-\eta \rightarrow u=p(\xi) e^\eta+q(\xi) e^{-\eta}-\eta=p(x-y) e^y+q(x-y) e^{-y}-y$ solución general
 $\rightarrow u_y(x,y)=[p(x-y)-p'(x-y)] e^y-[q'(x-y)+q(x-y)] e^{-y}-1$.

$\begin{cases} p(x)+q(x)=1 \rightarrow p'(x)+q'(x)=0 \\ p(x)-q(x)-p'(x)-q'(x)=1, & p(x)-q(x)=1 \end{cases} \rightarrow p(x)=1, q(x)=0. u(x,y)=e^y-y$.

Comprobamos: $u(x,0)=1-0$, $u_y(x,0)=1-1$. Y además: $u_{yy}+2u_{xy}+u_{xx}-u=e^y+0+0-e^y+y=y$.

12. $u_t-u_{xx}=0, x \in (0, \pi), t > 0$
 $u(x,0)=0, u_x(0,t)=u_x(\pi,t)=1$ $v=x$ cumple las condiciones de contorno.
 $w=u-x \rightarrow \begin{cases} w_t-w_{xx}=0 \\ w(x,0)=-x, w_x(0,t)=w_x(\pi,t)=0 \end{cases} \rightarrow$
 $\begin{cases} X''+\lambda X=0 \\ X'(0)=X'(\pi)=0 \end{cases} \rightarrow \lambda_n=n^2, X_n=\{\cos nx\}, n=0,1,\dots$ y además $T'+\lambda T=0 \rightarrow T_n=\{e^{-n^2 t}\}$.

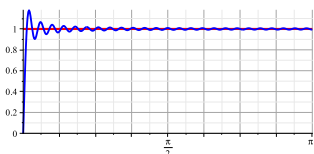
$w=\frac{a_0}{2}+\sum_{n=1}^{\infty} a_n e^{-n^2 t} \cos nx \rightarrow u(x,0)=\frac{a_0}{2}+\sum_{n=1}^{\infty} a_n \cos nx=-x$, con $a_n=-\frac{2}{\pi} \int_0^\pi x \cos nx dx$:

$a_0=-\frac{2}{\pi} \frac{\pi^2}{2}=-\pi$, $a_n=-\frac{2}{n\pi} x \sin nx \Big|_0^\pi + \frac{2}{n\pi} \int_0^\pi \sin nx dx = \frac{2}{n^2\pi} [1-\cos n\pi] = \begin{cases} 4/(n^2\pi), & n \text{ impar} \\ 0, & n \text{ par} \end{cases}$

$u(x,t)=x-\frac{\pi}{2}+\sum_{m=1}^{\infty} \frac{4}{\pi(2m-1)^2} e^{-(2m-1)^2 t} \cos(2m-1)x$

13. a) Sabemos que para $\begin{cases} X''+\lambda X=0 \\ X(0)=X'(\pi)=0 \end{cases}$ son $\lambda_n=\frac{(2n-1)^2}{4}$, $X_n=\{\sin \frac{(2n-1)x}{2}\}$ y $\langle X_n, X_n \rangle = \frac{L}{2} = \frac{\pi}{2}$. Como $r=1$,

es: $c_n = \frac{\langle 1, X_n \rangle}{\langle X_n, X_n \rangle} = \frac{2}{\pi} \int_0^\pi \sin \frac{(2n-1)x}{2} x dx = -\frac{4}{\pi(2n-1)} \cos \frac{(2n-1)x}{2} \Big|_0^\pi = \frac{4}{\pi(2n-1)}$. Por tanto: $1 = \frac{4}{\pi} \sum_{n=1}^{\infty} \frac{1}{2n-1} \sin \frac{(2n-1)x}{2}$.



f continua en $x=\frac{\pi}{2} \in (0, \pi) \rightarrow$ la suma de la serie será su valor en el punto: $1=f(\frac{\pi}{2})$.

[Con ordenador, sumando 50 términos de la serie para $x=\frac{\pi}{2}$ se obtiene 1.0090..., con 1000 sale 0.9995..., ... También se puede con un ordenador dibujar diferentes sumas parciales (a la izquierda S_{50}) y ver cómo se acercan en todo $(0, \pi]$ a la $f(x)=1$].

b) $u_t-4(1+2t)u_{xx}=0, x \in (0, \pi), t > 0$
 $u(x,0)=1, u(0,t)=u_x(\pi,t)=0$ $u=XT \rightarrow XT'=(4+8t)X''T, \frac{X''}{X}=\frac{T'}{(4+8t)T}=-\lambda \rightarrow$

$\begin{cases} X''+\lambda X=0 \\ X(0)=X'(\pi)=0 \end{cases}, \lambda_n=\frac{(2n-1)^2}{4}, X_n=\{\sin \frac{(2n-1)x}{2}\}$. Y además: $T'+4\lambda_n(1+2t)T=0, T_n=\{e^{-(2n-1)^2(t+t^2)}\}$.

Probamos: $u=\sum_{n=1}^{\infty} c_n e^{-(2n-1)^2(t+t^2)} \sin \frac{(2n-1)x}{2}$. Imponiendo el dato inicial $u(x,0)=\sum_{n=1}^{\infty} c_n \sin \frac{(2n-1)x}{2}=1$.

Los c_n son los calculados en a). Por tanto, la solución es: $u(x,t)=\frac{4}{\pi} \sum_{n=1}^{\infty} \frac{1}{2n-1} e^{-(2n-1)^2(t+t^2)} \sin \frac{(2n-1)x}{2}$.

14.
$$\begin{cases} u_t - 8tu_{xx} = 0, & x \in (0, \pi), t > 0 \\ u(x, 0) = f(x), & u_x(0, t) = u_x(\pi, t) = 0 \end{cases} \quad \text{con } \mathbf{a)} \ f(x) = \cos \frac{x}{2}, \ \mathbf{b)} \ f(x) = \frac{\pi}{4}.$$

$$u(x, t) = X(x) T(t), \quad \frac{T'}{8tT} = \frac{X''}{X} = -\lambda \rightarrow \begin{cases} X'' + \lambda X = 0 \\ X'(0) = X(\pi) = 0 \end{cases} \rightarrow \lambda_n = \frac{(2n-1)^2}{4}, \ n=1, 2, \dots, \ X_n = \left\{ \cos \frac{(2n-1)x}{2} \right\}.$$

$$\text{Y adem\'as } T' = -8\lambda t T = -2(2n-1)^2 t T \rightarrow T_n = \{e^{-(2n-1)^2 t^2}\}.$$

Probamos pues: $u(x, t) = \sum_{n=1}^{\infty} c_n e^{-(2n-1)^2 t^2} \cos \frac{(2n-1)x}{2}$, y s\'olo falta cumplirse $u(x, 0) = \sum_{n=1}^{\infty} c_n \cos \frac{(2n-1)x}{2} = f(x)$.

Para **a)** es claramente $c_1 = 1$ y los otros $c_n = 0 \rightarrow u(x, t) = e^{-t^2} \cos \frac{x}{2}$ es la soluci\'on (f\'acilmente comprobable).

En el caso **b)** debemos hallar el desarrollo en serie de autofunciones de $f(x) = \frac{\pi}{4}$:

$$c_n = \frac{2}{\pi} \int_0^{\pi} \frac{\pi}{4} \cos \frac{(2n-1)x}{2} dx = \frac{1}{2n-1} \sin \frac{(2n-1)x}{2} \Big|_0^{\pi} = \frac{(-1)^{n+1}}{2n-1}.$$
 La soluci\'on es ahora la serie:

$$u(x, t) = e^{-t^2} \cos \frac{x}{2} - \frac{1}{3} e^{-9t^2} \cos \frac{3x}{2} + \dots = \sum_{n=1}^{\infty} \frac{(-1)^{n+1}}{2n-1} e^{-(2n-1)^2 t^2} \cos \frac{(2n-1)x}{2}.$$

15. **a)** El formulario nos da las autofunciones $X_n = \{\sin nx\}$, $n=1, 2, \dots$, y la expresi\'on para los coeficientes:

$$c_n = \frac{2}{\pi} \int_0^{\pi} \left(\frac{\pi}{2} - x \right) \sin nx dx = \left[\frac{2x-\pi}{n\pi} \cos nx \right]_0^{\pi} + \frac{2}{\pi} \int_0^{\pi} \cos nx dx = \frac{(-1)^{n+1}}{n} \rightarrow \frac{\pi}{2} - x = \sum_{n=1}^{\infty} \frac{1}{n} \sin 2nx.$$

b) $u = X(x) T(t) \rightarrow \begin{cases} X'' + \lambda X = 0 \\ X(0) = X(\pi) = 0 \end{cases} \rightarrow X_n = \{\sin nx\}$ de antes. [Y adem\'as $T' + \lambda T = 0$ que ahora no usamos].

Llevamos la serie $u(x, t) = \sum_{n=1}^{\infty} T_n(t) \sin nx$ a la EDP no homog\'enea y al dato inicial para calcular los T_n :

$$\sum_{n=1}^{\infty} [T'_n + n^2 T_n] \sin nx = \sin x. \quad u(x, 0) = \sum_{n=1}^{\infty} T_n(0) \sin nx = \frac{\pi}{2} - x = \sum_{n=1}^{\infty} \frac{1}{n} \sin 2nx \rightarrow \begin{matrix} T_{2n-1}(0) = 0 \\ T_{2n}(0) = \frac{1}{n} \end{matrix}$$

Son no nulos: $\begin{cases} T'_1 + T_1 = 1 \\ T_1(0) = 0 \end{cases} \xrightarrow{T_p \text{ a ojo}} T_1 = C e^{-t} + 1 \xrightarrow{d.i.} C = -1 \quad \text{y} \quad \begin{cases} T'_n + 4n^2 T_n = 0 \\ T_{2n}(0) = \frac{1}{n} \end{cases} \xrightarrow{d.i.} T_{2n} = C e^{-4n^2 t} \xrightarrow{d.i.} C = \frac{1}{n}.$

$$\rightarrow u = (1 - e^{-t}) \sin x + \sum_{n=1}^{\infty} \frac{1}{n} e^{-4n^2 t} \sin 2nx = (1 - e^{-t}) \sin x + e^{-4t} \sin 2x + \frac{1}{2} e^{-16t} \sin 4x + \frac{1}{3} e^{-36t} \sin 6x + \dots$$

16. **a)**
$$\begin{cases} u_t - u_{xx} = 4 \sin 2x, & x \in (0, \pi), t > 0 \\ u(x, 0) = 2 \sin x, & u(0, t) = u(\pi, t) = 0 \end{cases} \quad u(x, t) = X(x) T(t) \rightarrow \begin{cases} X'' + \lambda X = 0 \\ X(0) = X(\pi) = 0 \end{cases} \rightarrow X_n = \{\sin nx\}, \ n=1, 2, \dots$$

Llevamos a la EDP: $u(x, t) = \sum_{n=1}^{\infty} T_n(t) \sin nx \rightarrow \sum_{n=1}^{\infty} [T'_n + n^2 T_n] \sin nx = 4 \sin 2x$ (ya desarrollada en senos).

Y del dato inicial se deduce: $u(x, 0) = \sum_{n=1}^{\infty} T_n(0) \sin nx = 2 \sin x \rightarrow T_1(0) = 2$ y el resto de $T_n(0) = 0$.

S\'olo nos falta resolver: $\begin{cases} T'_1 + T_1 = 0 \\ T_1(0) = 2 \end{cases}$ y $\begin{cases} T'_2 + 4T_2 = 4 \\ T_2(0) = 0 \end{cases}$, pues el resto de $T_{n \geq 3} \equiv 0$ claramente.

$$T_1 = C e^{-t} \xrightarrow{T_1(0)=2} T_1(t) = 2e^{-t}. \quad T_{2p} = 1 \text{ (a ojo)}, \ T_2 = C e^{-4t} + 1 \left[\text{\'o } T_2 = C e^{-4t} + e^{-4t} \int 4e^{4t} dt \right] \xrightarrow{T_2(0)=0} T_2(t) = 1 - e^{-4t}.$$

La soluci\'on es entonces: $u(x, t) = 2e^{-t} \sin x + (1 - e^{-4t}) \sin 2x$.

b)
$$\begin{cases} u_t - u_{xx} + u = e^{-2t}, & x \in (0, \pi), t > 1 \\ u(x, 0) = \cos x, & u_x(0, t) = u_x(\pi, t) = 0 \end{cases} \quad u = XT \rightarrow XT' + XT = X''T, \quad \frac{X''}{X} = \frac{T'}{T} + 1 = -\lambda \rightarrow \begin{cases} X'' + \lambda X = 0 \\ X'(0) = X'(\pi) = 0 \end{cases}$$

Llevamos a la ecuaci\'on: $u = T_0(t) + \sum_{n=1}^{\infty} T_n(t) \cos nx \rightarrow T_0 + T'_0 + \sum_{n=1}^{\infty} [T'_n + T_n + n^2 T_n] \cos nx = e^{-2t}.$

$[e^{-2t}$ (constante en x) ya est\'a desarrollada en las autofunciones $\{\cos nx\}$, pues la primera de ellas es $\{1\}]$.

Imponiendo el dato inicial $u(x, 0) = \sum_{n=1}^{\infty} T_n(0) \cos nx = \cos x$, deducimos: $T_1(0) = 1$ y $T_n(0) = 0$ si $n \neq 1$.

S\'olo son no nulas: $\begin{cases} T'_0 + T_0 = e^{-2t} \\ T_0(0) = 0 \end{cases} \rightarrow T_0(t) = C e^{-t} + e^{-t} \int e^t e^{-2t} dt = e^{-t} - e^{-2t} \xrightarrow{T_0(0)=0} C = 1,$
[La T_p se pod\'ia hallar tanteando: $T_p = A e^{-2t} \rightarrow A = -1 \rightarrow T_p = -e^{-2t}$].

$$\begin{cases} T'_1 + 2T_1 = 0 \\ T_1(0) = 1 \end{cases} \rightarrow T_1 = C e^{-2t} \xrightarrow{T_1(0)=1} T_1(t) = e^{-2t}. \quad \text{La soluci\'on es: } u(x, t) = e^{-t} - e^{-2t} + e^{-2t} \cos x.$$

[hasta aqu\'i el control]

17. a)
$$\boxed{u_{tt} - 4u_{xx} = e^{-t}, x, t \in \mathbf{R}} \quad u(x, 0) = x^2, u_t(x, 0) = -1$$

$$u = \frac{1}{2}[(x+2t)^2 + (x-2t)^2] + \frac{1}{4} \int_{x-2t}^{x+2t} ds + \frac{1}{4} \int_0^t \int_{x-2[t-\tau]}^{x+2[t-\tau]} e^{-\tau} ds d\tau$$

$$= x^2 + 4t^2 + e^{-t} - 1.$$

Una solución particular que sólo depende de t es: $v_{tt} = e^{-t} \rightarrow v = e^{-t}$. Con $w = u - e^{-t}$ se tiene:

$$\begin{cases} w_{tt} - w_{xx} = 0 \\ w(x, 0) = x^2 - 1, w_t(x, 0) = 0 \end{cases} \rightarrow w = \frac{1}{2}[(x+2t)^2 - 1 + (x-2t)^2 - 1] = x^2 + 4t^2 - 1, \text{ como antes.}$$

b)
$$\boxed{u_{tt} - 4u_{xx} = 4, x, t \in \mathbf{R}} \quad u(x, 0) = u_t(x, 0) = 0$$
 a) Las características de esta ecuación de ondas son:

$$\begin{cases} \xi = x+2t \\ \eta = x-2t \end{cases} \rightarrow \begin{cases} u_{xx} = u_{\xi\xi} + 2u_{\xi\eta} + u_{\eta\eta} \\ u_{tt} = 4u_{\xi\xi} - 8u_{\xi\eta} + 4u_{\eta\eta} \end{cases} \rightarrow \boxed{u_{\xi\eta} = -\frac{1}{4}} \text{ forma canónica.}$$

b) Para hallar la solución pedida podemos utilizar la fórmula de D'Alembert directamente:

$$u(x, t) = \frac{1}{4} \int_0^t \int_{x-2[t-\tau]}^{x+2[t-\tau]} 4 ds d\tau = \int_0^t 4[t-\tau] d\tau = [-2(t-\tau)^2]_0^t = \boxed{2t^2}.$$

O bien, por ser $v = 2t^2$ solución de la EDP, con $w = u - v$ se convertirá la más fácil:

$$\begin{cases} w_{tt} - w_{xx} = 0 \\ w(x, 0) = w_t(x, 0) = 0 \end{cases} \begin{matrix} \text{(casualidad que} \\ \text{sea tan sencilla)} \end{matrix} \rightarrow w = 0 \rightarrow u = w + v = 2t^2.$$

Mucho peor es calcularla a partir de la forma canónica de arriba. Resumiendo algo los cálculos:

$$u_{\xi} = -\frac{1}{4}\eta + p^*(\xi), u = -\frac{1}{4}\xi\eta + p(\xi) + q(\eta) = t^2 - \frac{1}{4}x^2 + p(x+2t) + q(x-2t), u_t = 2t + 2p'(x+2t) - 2q'(x-2t)$$

$$\xrightarrow{d.i.} \begin{cases} u(x, 0) = -x^2/4 + p(x) + q(x) = 0 \\ u_t(x, 0) = 2p'(x) - 2q'(x) = 0 \end{cases} \rightarrow p(x) = q(x) \nearrow p(x) = q(x) = \frac{x^2}{8} \uparrow$$

[Lo que no tiene sentido es usar separación de variables, pues es en todo $\mathbf{R}$ y no hay condiciones de contorno].

c)
$$\boxed{u_{tt} - u_{xx} = 2x, x, t \in \mathbf{R}} \quad u(x, 0) = u_t(x, 0) = 0$$
 Las características de la ecuación de ondas son:

$$\begin{cases} \xi = x+t \\ \eta = x-t \end{cases}, \begin{cases} u_{xx} = u_{\xi\xi} + 2u_{\xi\eta} + u_{\eta\eta} \\ u_{tt} = u_{\xi\xi} - 2u_{\xi\eta} + u_{\eta\eta} \end{cases}, -4u_{\xi\eta} = 2x \xrightarrow{\xi+\eta=2x} \boxed{u_{\xi\eta} = -\frac{1}{4}(\xi+\eta)}.$$

Para hallar la solución podemos utilizar D'Alembert directamente:

$$u(x, t) = \frac{1}{2} \int_0^t \int_{x-[t-\tau]}^{x+[t-\tau]} 2s ds d\tau = \int_0^t [s^2]_{x-[t-\tau]}^{x+[t-\tau]} d\tau = \int_0^t 2x[t-\tau] d\tau = x[-(t-\tau)^2]_0^t = \boxed{xt^2}.$$

O, por ser $v = -\frac{1}{3}x^3$ solución de la EDP, con $w = u - v$ se convertirá en una homogénea más fácil:

$$\begin{cases} w_{tt} - w_{xx} = 0 \\ w(x, 0) = x^3/3, w_t(x, 0) = 0 \end{cases} \rightarrow w = \frac{1}{6}[(x+t)^3 + (x-t)^3] = \frac{1}{3}x^3 + xt^2, u = w + v = xt^2.$$

Mucho peor es calcularla a partir de la forma canónica de arriba. Resumiendo algo los cálculos:

$$u_{\xi} = -\frac{1}{4}\xi\eta - \frac{1}{8}\eta^2 + p(\xi), u = -\frac{1}{8}\xi\eta(\xi+\eta) + p(\xi) + q(\eta) = \frac{1}{4}(xt^2 - x^3) + p(x+t) + q(x-t),$$

$$u_t = -\frac{xt}{2} + p'(x+t) - q'(x-t) \xrightarrow{d.i.} \begin{cases} u(x, 0) = -x^3/4 + p(x) + q(x) = 0 \\ u_t(x, 0) = p'(x) - q'(x) = 0 \end{cases} \rightarrow p(x) = q(x) \nearrow p(x) = q(x) = \frac{x^3}{8} \uparrow$$

[Lo que no tiene sentido es usar separación de variables, pues es en todo $\mathbf{R}$ y no hay condiciones de contorno].

18.
$$\boxed{u_{tt} - 4u_{xx} = 0, x \in [0, \pi], t \in \mathbf{R}} \quad u(x, 0) = 0, u_t(x, 0) = \sin 3x$$

$$u(0, t) = u(\pi, t) = 0$$
 a)
$$u(x, t) = X(x)T(t), \quad \frac{X''}{X} = \frac{T''}{4T} = -\lambda \rightarrow \begin{cases} X'' + \lambda X = 0 \\ X(0) = X(\pi) = 0 \end{cases}$$

$$\rightarrow \lambda_n = n^2, n = 1, 2, \dots, X_n = \{\sin nx\}.$$

Y además
$$\begin{cases} T'' + 4\lambda T = 0 \\ T(0) = 0 \end{cases}, \mu^2 + 4n^2 = 0 \rightarrow T = c_1 \cos 2nt + c_2 \sin 2nt \xrightarrow{d.i.} c_1 = 0, T_n = \{\sin 2nt\}.$$

Satisface la EDP, $u(x, 0) = 0$ y las condiciones de contorno la serie: $u(x, t) = \sum_{n=1}^{\infty} c_n \sin 2nt \sin nx.$

El dato inicial que falta nos da: $u_t(x, 0) = \sum_{n=1}^{\infty} 2nc_n \sin nx = \sin 3x \rightarrow 6c_3 = 1$ y el resto de los $c_n = 0$.

La solución es por tanto:
$$\boxed{u(x, t) = \frac{1}{6} \sin 6t \sin 3x}$$
 (fácil de comprobar).

b) Como $g(x) = \sin 3x$ es impar y de periodo 2π , su extensión $g^*(x)$ es ella misma. Y D'Alembert nos da:

$$u = \frac{1}{4} \int_{x-2t}^{x+2t} g^*(s) ds = \frac{1}{4} \int_{x-2t}^{x+2t} \sin 3s ds = -\frac{1}{12} \cos 3s \Big|_{x-2t}^{x+2t} = \frac{1}{12} [\cos(3x-6t) - \cos(3x+6t)] = \frac{1}{6} \sin 6t \sin 3x.$$

$cc+ss \qquad \qquad \qquad cc-ss$

19. a) $y'' + y = 2$ $\mu^2 + 1 = 0$, $\mu = \pm i$, $y_p = 2 \rightarrow y = c_1 \cos x + c_2 \sin x + 2$.

Imponiendo los datos: $\begin{cases} c_1 + 2 = 0 \\ c_2 = 0 \end{cases}$, se obtiene la solución única del problema: $y = 2 - 2 \cos x$.

b) $\begin{cases} u_{tt} - u_{xx} = 2 \sin x, x \in [0, \pi], t \in \mathbf{R} \\ u(x, 0) = u_t(x, 0) = u(0, t) = u(\pi, t) = 0 \end{cases}$ Al ser un problema no homogéneo, debemos probar una serie con las autofunciones X_n del homogéneo.

Sabemos que al separar variables aparece $\begin{cases} X'' + \lambda X = 0 \\ X(0) = X(\pi) = 0 \end{cases} \rightarrow X_n = \{\sin nx\}, n = 1, 2, \dots$

Llevamos, pues, a la ecuación $u(x, t) = \sum_{n=1}^{\infty} T_n(t) \sin nx$, obteniendo $\sum_{n=1}^{\infty} [T_n'' + n^2 T_n] \sin nx = 2 \sin x$.

Los datos iniciales: $u(x, 0) = \sum_{n=1}^{\infty} T_n(0) \sin nx = 0$, $u_t(x, 0) = \sum_{n=1}^{\infty} T_n'(0) \sin nx = 0 \Rightarrow T_n(0) = T_n'(0) = 0 \forall n$.

Todas las ecuaciones son homogéneas con datos nulos (y su solución es $T_n \equiv 0$) excepto $\begin{cases} T_1'' + T_1 = 2 \\ T_1(0) = T_1'(\pi) = 0 \end{cases}$,

que es lo resuelto (con otros nombres) en a). La solución es, entonces: $u(x, t) = 2(1 - \cos t) \sin x$.

[Se podría usar D'Alembert. Como $2 \sin x$ es impar y 2π -periódica, su extensión es ella misma, y entonces:

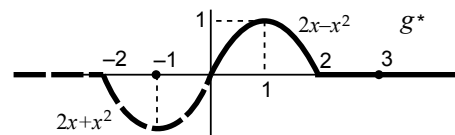
$$u(x, t) = \frac{1}{2} \int_0^t \int_{x-t}^{x+t} 2 \sin s \, ds \, d\tau = \int_0^t [\cos(x-t+\tau) - \cos(x+t-\tau)] \, d\tau \\ = [\sin(x-t+\tau) + \sin(x+t-\tau)]_0^t = 2 \sin x - 2 \sin x \cos t.$$

20. $\begin{cases} u_{tt} - u_{xx} = 0, x \geq 0, t \in \mathbf{R} \\ u(x, 0) = 0, u_t(x, 0) = \begin{cases} 2x - x^2, x \in [0, 2] \\ 0, x \in [2, \infty) \end{cases} \\ u(0, t) = 0 \end{cases}$

La solución es: $u(x, t) = \frac{1}{2} \int_{x-t}^{x+t} g^*(s) \, ds$, siendo g^* la extensión impar de g respecto del origen a todo $\mathbf{R}$.

En $[-2, 0]$ la expresión de g^* será: $-[2(-x) - (-x)^2]$.

Por tanto la expresión de g^* es: $g^*(x) = \begin{cases} 2x + x^2 & \text{si } x \in [-2, 0] \\ 2x - x^2 & \text{si } x \in [0, 2] \\ 0 & \text{en el resto de } \mathbf{R} \end{cases}$.



Entonces: i) $u(5, 2) = \frac{1}{2} \int_3^7 g^* = \frac{1}{2} \int_3^7 0 \, ds = 0$.

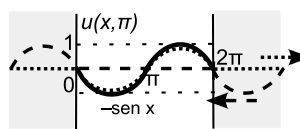
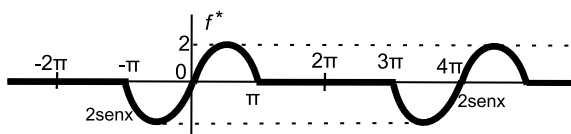
ii) $u(3, 2) = \frac{1}{2} \int_1^5 g^* = \frac{1}{2} \int_1^2 (2s - s^2) \, ds = \frac{1}{2} [s^2 - \frac{s^3}{3}]_1^2 = \frac{1}{3}$.

iii) $u(1, 2) = \frac{1}{2} \int_{-1}^3 g^* = \frac{1}{2} \int_{-1}^0 (2s + s^2) \, ds + \frac{1}{2} \int_0^2 (2s - s^2) \, ds \underset{g^* \text{ impar}}{=} \frac{1}{2} \int_1^2 (2s - s^2) \, ds = \frac{1}{3}$.

21. $\begin{cases} u_{tt} - u_{xx} = 0, x \in [0, 2\pi], t \in \mathbf{R} \\ u(x, 0) = \begin{cases} 2 \sin x, x \in [0, \pi] \\ 0, x \in [\pi, 2\pi] \end{cases}, u_t(x, 0) = 0 \\ u(0, t) = u(2\pi, t) = 0 \end{cases}$

La solución es $u(x, t) = \frac{1}{2} [f^*(x+t) + f^*(x-t)]$, con f^* extensión impar y 4π -periódica de la f inicial. Para dibujar $u(x, \pi)$ basta trasladar $\frac{1}{2}f(x)$ a la izquierda y a la derecha π unidades y sumar las gráficas en $[0, 2\pi]$.

[Sólo queda lo que va a la derecha con la mitad de altura].



Como si $x \in [0, 2\pi]$ siempre $x+\pi \in [\pi, 3\pi]$, $x-\pi \in [-\pi, \pi]$ y en todo este intervalo es $f^*(x) = 2 \sin x$ ($\sin x$ impar):

$$u(x, \pi) = \frac{1}{2} [f^*(x+\pi) + f^*(x-\pi)] = \frac{1}{2} [0 + 2 \sin(x-\pi)] = -\sin x \quad \forall x \in [0, 2\pi].$$

$$u = XT \rightarrow \begin{cases} X'' + \lambda X = 0 \\ X(0) = X(2\pi) = 0 \end{cases} \rightarrow \lambda_n = \frac{n^2}{4}, X_n = \{\sin \frac{nx}{2}\}, n = 1, 2, \dots \text{ y } \begin{cases} T' + \lambda T = 0 \\ T'(0) = 0 \end{cases} \rightarrow T_n = \{\cos \frac{nt}{2}\}.$$

$$u(x, t) = \sum_{n=1}^{\infty} c_n \cos \frac{nt}{2} \sin \frac{nx}{2} \rightarrow u(x, 0) = \sum_{n=1}^{\infty} c_n \sin \frac{nx}{2} = \begin{cases} 2 \sin x, x \in [0, \pi] \\ 0, x \in [\pi, 2\pi] \end{cases} \rightarrow$$

$$c_n = \frac{2}{2\pi} \int_0^\pi 2 \sin x \sin \frac{nx}{2} \, dx = \frac{1}{\pi} \int_0^\pi [\cos(\frac{n}{2}-1)x - \cos(\frac{n}{2}+1)x] \, dx = \frac{1}{\pi} \left[\frac{\sin(\frac{n}{2}-1)\pi}{\frac{n}{2}-1} - \frac{\sin(\frac{n}{2}+1)\pi}{\frac{n}{2}+1} \right] \\ = \frac{2}{\pi} \left[-\frac{\sin \frac{n\pi}{2}}{n-2} + \frac{\sin \frac{n\pi}{2}}{n+2} \right] = -\frac{8 \sin \frac{n\pi}{2}}{\pi(n^2-4)} = \begin{cases} 0, n=2m \\ \frac{8}{\pi} \frac{(-1)^m}{(2m-1)^2-4}, n=2m-1 \end{cases}. \text{ Además } c_2 = \frac{1}{\pi} \int_0^\pi [1 - \cos 2x] \, dx = 1.$$

$$u(x, t) = \cos t \sin x + \frac{8}{\pi} \sum_{m=1}^{\infty} \frac{(-1)^m}{(2m-1)^2-4} \cos \frac{(2m-1)t}{2} \sin \frac{(2m-1)x}{2} \rightarrow u(x, \pi) = -\sin x, \text{ pues } \cos \pi = -1, \cos \frac{(2m-1)\pi}{2} = 0.$$

22. $\begin{cases} u_{tt} + 4u_t - u_{xx} = 0, & x \in [0, \pi], t \in \mathbf{R} \\ u(x, 0) = \sin 2x, & u_t(x, 0) = u(0, t) = u(\pi, t) = 0 \end{cases}$ Como la ecuación es nueva empezamos separando variables:

$$u = XT \rightarrow X[T'' + 4T'] - X''T = 0 \rightarrow \frac{T'' + 4T'}{T} = \frac{X''}{X} = -\lambda \rightarrow$$

$$\begin{cases} X'' + \lambda X = 0 \\ X(0) = X(\pi) = 0 \end{cases} \rightarrow \lambda_n = n^2, X_n = \{\sin nx\}, n = 1, 2, \dots \rightarrow T'' + 4T' + n^2T = 0, \mu = -2 \pm \sqrt{4 - n^2} \rightarrow$$

$$T_1 = c_1 e^{(-2 + \sqrt{3})t} + c_2 e^{(-2 - \sqrt{3})t}, T_2 = (c_1 + c_2 t) e^{-2t}, T_{n \geq 3} = e^{-2t} (c_1 \cos \sqrt{4 - n^2} t + c_2 \sin \sqrt{4 - n^2} t).$$

Probamos, pues, $u(x, t) = \sum_{n=1}^{\infty} T_n(t) \sin nx$. Sólo faltan las condiciones iniciales:

$$u(x, 0) = \sum_{n=1}^{\infty} T_n(0) \sin nx = \sin 2x, u_t(x, 0) = \sum_{n=1}^{\infty} T'_n(0) \sin nx = 0 \rightarrow T_n(t) \equiv 0, \text{ si } n \neq 2 \quad [\text{ecuación homogénea con datos nulos}]$$

Sólo es no nula T_2 : $T_2(0) = c_1 = 1 \rightarrow u(x, t) = (1 + 2t) e^{-2t} \sin 2x$. [Tiende a pararse la cuerda con rozamiento].

23. $\begin{cases} u_{tt} - u_{xx} = 0, & x \in [0, \pi], t \in \mathbf{R} \\ u(x, 0) = 0, & u_t(x, 0) = x \\ u_x(0, t) = u_x(\pi, t) = 0 \end{cases}$ $u = X(x) T(t), \frac{X''}{X} = \frac{T''}{T} = -\lambda \rightarrow \begin{cases} X'' + \lambda X = 0 \\ X'(0) = X'(\pi) = 0 \end{cases}$

$$\rightarrow \lambda_n = n^2, n = 0, 1, \dots, X_n = \{\cos nx\} \quad [\text{con } X_0 = \{1\}].$$

Y además $\begin{cases} T'' + \lambda T = 0 \\ T(0) = 0 \end{cases}, \mu^2 + n^2 = 0 \rightarrow \begin{cases} T = c_1 + c_2 t \rightarrow T_0 = \{t\} \\ T = c_1 \cos nt + c_2 \sin nt \rightarrow T_n = \{\sin nt\}, n \geq 1 \end{cases}$

Satisface la EDP, $u(x, 0) = 0$ y las condiciones de contorno la serie: $u(x, t) = \frac{c_0}{2} t + \sum_{n=1}^{\infty} c_n \sin nt \cos nx$.

El dato inicial que falta nos da: $u_t(x, 0) = \frac{c_0}{2} + \sum_{n=1}^{\infty} n c_n \cos nx = x \rightarrow c_0 = \frac{2}{\pi} \int_0^{\pi} x dx = \pi$, y para $n \geq 1$:

$$c_n = \frac{2}{n\pi} \int_0^{\pi} x \cos nx dx = \frac{2x}{n^2\pi} \sin nx \Big|_0^{\pi} - \frac{2}{n^2\pi} \int_0^{\pi} \sin nx dx = \frac{2}{n^3\pi} [(-1)^n - 1] \quad (\text{se anula si } n \text{ par}).$$

La solución es por tanto:

$$u(x, t) = \frac{\pi}{2} t - \frac{4}{\pi} \left[\sin t \cos x + \frac{1}{27} \sin 3t \cos 3x + \dots \right] = \frac{4}{\pi} \sum_{n=1}^{\infty} \frac{1}{(2n-1)^3} \sin(2n-1)t \cos(2n-1)x.$$

[Esas condiciones de contorno significan que se da libertad a los extremos de la cuerda, y como esta está subiendo inicialmente, la altura de la cuerda lo hace indefinidamente].

24. a) $\begin{cases} \Delta u = 0, & (x, y) \in (0, \pi) \times (0, \pi) \\ u(\pi, y) = 5 + \cos y \\ u(0, y) = u_y(x, 0) = u_y(x, \pi) = 0 \end{cases}$ $Y'' + \lambda Y = 0, Y'(0) = Y'(\pi) = 0 \rightarrow \lambda_n = n^2, Y_n = \{\cos ny\}, n = 0, 1, \dots \rightarrow$

$$X'' - n^2 X = 0, X(0) = 0 \rightarrow X_0 = \{x\} \text{ y } X_n = \{\text{sh } nx\} \text{ si } n > 0.$$

$$u = c_0 x + \sum_{n=1}^{\infty} c_n \text{sh } nx \cos ny \rightarrow u(x, \pi) = c_0 \pi + \sum_{n=1}^{\infty} c_n \text{sh } n\pi \cos ny = 5 + \cos y \rightarrow u(x, y) = \frac{5x}{\pi} + \frac{\text{sh } x}{\text{sh } \pi} \cos y.$$

b) $\begin{cases} \Delta u = y \cos x, & (x, y) \in (0, \pi) \times (0, 1) \\ u_x(0, y) = u_x(\pi, y) = u_y(x, 0) = u_y(x, 1) = 0 \end{cases}$ $u = \sum_{n=0}^{\infty} Y_n(y) \cos nx \rightarrow \sum_{n=0}^{\infty} [Y''_n - n^2 Y_n] \cos nx = y \cos x \rightarrow$

$$Y''_1 - Y_1 = y \rightarrow Y_1 = c_1 e^y + c_2 e^{-y} - y \xrightarrow{\text{c.c.}} Y_1 = \frac{e^y - e^{-y}}{1 + e} - y. \text{ Los } Y_n \equiv 0 \text{ para } n \geq 2.$$

Como es de Neumann aparece (al resolver $Y''_0 = 0 + \text{c.c.}$) una C arbitraria: $u = C + \left[\frac{e^y - e^{-y}}{1 + e} - y \right] \cos x$.

c) $\begin{cases} u_{xx} + u_{yy} + 6u_x = 0 \text{ en } (0, \pi) \times (0, \pi) \\ u_y(x, 0) = 0, u_y(x, \pi) = 0 \\ u_x(0, y) = 0, u(\pi, y) = 2 \cos^2 2y \end{cases}$ $u = XY \rightarrow \frac{X'' + 6X'}{X} = -\frac{Y''}{Y} = \lambda \rightarrow \begin{cases} Y'' + \lambda Y = 0, Y'(0) = Y'(\pi) = 0 \\ X'' + 6X' - \lambda X = 0, X'(0) = 0 \end{cases}$

$$\rightarrow \lambda_n = n^2, Y_n = \{\cos ny\} \quad n = 0, 1, \dots \rightarrow X'' + 6X' - n^2 X = 0, X = c_1 e^{(\sqrt{9+n^2}-3)x} + c_2 e^{-(\sqrt{9+n^2}+3)x} \xrightarrow{X'(0)=0}$$

$$X_0 = \{1\}; X_n = \left\{ (\sqrt{9+n^2}+3) e^{(\sqrt{9+n^2}-3)x} + (\sqrt{9+n^2}-3) e^{-(\sqrt{9+n^2}+3)x} \right\}, n \geq 1.$$

$$u(x, y) = \sum_{n=0}^{\infty} c_n X_n(x) \cos ny \rightarrow u(\pi, y) = \sum_{n=0}^{\infty} c_n X_n(\pi) \cos ny = 1 + \cos 4y \rightarrow$$

$$c_0 = 1, c_4 = \frac{1}{x_4(\pi)} \text{ y los demás cero } \rightarrow u = 1 + \frac{4e^{2x} + e^{-8x}}{4e^{2\pi} + e^{-8\pi}} \cos 4y.$$

25. a)
$$\begin{cases} u_{rr} + \frac{1}{r}u_r + \frac{1}{r^2}u_{\theta\theta} = 0, & r < 2, 0 < \theta < \pi \\ u_r(2, \theta) = \sin 3\theta, & u(r, 0) = u(r, \pi) = 0 \end{cases} \quad u = R\Theta \rightarrow \begin{cases} \Theta'' + \lambda\Theta = 0 \\ \Theta(0) = \Theta(\pi) = 0 \end{cases}, \lambda_n = n^2, \{\sin n\theta\}, n = 1, 2, \dots$$

$$r^2 R'' + rR' - \lambda R = 0 \xrightarrow{\lambda=n^2} \mu = \pm n, R = c_1 r^2 + c_2 r^{-n} \xrightarrow{R \text{ acotado en } r=0} R_n = \{r^n\}, n = 1, 2, \dots$$

Probamos, pues, $u = \sum_{n=1}^{\infty} b_n r^n \sin n\theta$, y para calcular los b_n imponemos el dato que falta:

$$u_r(2, \theta) = \sum_{n=1}^{\infty} n 2^{n-1} b_n \sin n\theta = \sin 3\theta \rightarrow 12b_3 = 1 \text{ y el resto } 0. \quad \boxed{u(r, \theta) = \frac{1}{12} r^3 \sin 3\theta}.$$

Comprobando: $u_r = \frac{1}{4} r^2 \sin 3\theta, u_{rr} = \frac{1}{2} r \sin 3\theta, u_{\theta\theta} = -\frac{3}{4} r^2 \sin 3\theta \rightarrow \Delta u = r \sin 3\theta \left[\frac{1}{2} + \frac{1}{4} - \frac{3}{4} \right] = 0,$

$$u_r(2, \theta) = \frac{1}{4} \sin 3\theta, \text{ y es claro que se cumplen las condiciones de contorno.}$$



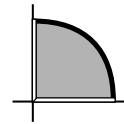
b)
$$\begin{cases} u_{rr} + \frac{1}{r}u_r + \frac{1}{r^2}u_{\theta\theta} = 0, & r < 1, 0 < \theta < \frac{\pi}{2} \\ u(1, \theta) = \cos 2\theta, & u_\theta(r, 0) = u_\theta(r, \frac{\pi}{2}) = 0 \end{cases} \quad u = R\Theta \rightarrow \begin{cases} \Theta'' + \lambda\Theta = 0 \\ \Theta'(0) = \Theta'(\frac{\pi}{2}) = 0 \end{cases}, \lambda_n = 4n^2, \Theta_n = \{\cos 2n\theta\}, n = 0, 1, 2, \dots$$

Y además $r^2 R'' + rR' - \lambda R = 0, \mu = \pm 2n, R \text{ acotado} \rightarrow R_n = \{r^{2n}\}. u = \frac{a_0}{2} + \sum_{n=1}^{\infty} a_{2n} r^{2n} \cos 2n\theta$

$$\rightarrow u(1, \theta) = \frac{a_0}{2} + \sum_{n=1}^{\infty} a_{2n} \cos 2n\theta = \cos 2\theta \rightarrow a_{2n} = 0 \forall n, \text{ menos } a_2 = 1. \quad \boxed{u = r^2 \cos 2\theta} \quad [=x^2 - y^2].$$

Comprobando: $u_r = 2r \cos 2\theta, u_{rr} = 2 \cos 2\theta, u_{\theta\theta} = -4r^2 \cos 2\theta \rightarrow \Delta u = [2 + 2 - 4] \cos 2\theta = 0,$

$$u(1, \theta) = \cos 2\theta, u_\theta = -2r^2 \sin 2\theta \text{ se anula si } \theta = 0 \text{ y si } \theta = \frac{\pi}{2}.$$



26.
$$\begin{cases} u_{rr} + \frac{1}{r}u_r + \frac{1}{r^2}u_{\theta\theta} = 0, & r < 1, 0 < \theta < \pi \\ u_r(1, \theta) = \frac{\theta}{2}, & u(r, 0) = u(r, \pi) = 0 \end{cases} \quad \text{Haciendo } u = R\Theta \text{ el formulario dice que sale:}$$

$$\begin{cases} \Theta'' + \lambda\Theta = 0 \\ \Theta(0) = \Theta(\pi) = 0 \end{cases}, \lambda_n = n^2, \Theta_n = \{\sin n\theta\}, n = 1, 2, \dots$$

[pues $u(r, 0) = R(r)\Theta(0) = 0, u(r, \pi) = R(r)\Theta(\pi) = 0$ y de $R(r) \equiv 0$ solo sale la solución cero],

y además: $r^2 R'' + rR' - \lambda R = 0 \xrightarrow{\lambda=n^2} R = c_2 r^n + c_2 r^{-n}. R \text{ debe ser acotada en } r=0 \rightarrow R_n = \{r^n\}.$

Probamos, pues, $u(r, \theta) = \sum_{n=1}^{\infty} c_n r^n \sin n\theta$, y para calcular los c_n imponemos el dato que falta:

$$u_r(1, \theta) = \sum_{n=1}^{\infty} n c_n \sin n\theta = \frac{\theta}{2} \rightarrow n c_n = \frac{2}{\pi} \int_0^\pi \frac{\theta}{2} \sin n\theta d\theta = -\frac{\theta}{n\pi} \cos n\theta \Big|_0^\pi + \frac{1}{n\pi} \int_0^\pi \cos n\theta d\theta = -\frac{\cos n\pi}{n} + 0 = \frac{(-1)^{n+1}}{n}.$$

La solución es, por tanto: $u(r, \theta) = \sum_{n=1}^{\infty} \frac{(-1)^{n+1}}{n^2} r^n \sin n\theta = \boxed{r \sin \theta - \frac{1}{4} r^2 \sin 2\theta + \dots}.$



27.
$$\begin{cases} u_{rr} + \frac{1}{r}u_r + \frac{1}{r^2}u_{\theta\theta} = 0, & r < 1, 0 < \theta < \frac{\pi}{2} \\ u(2, \theta) = \pi, & u(r, 0) = u_\theta(r, \frac{\pi}{2}) = 0 \end{cases} \quad \text{El formulario nos da las ecuaciones que aparecen separando variables en la EDP:}$$

$$u = R\Theta \rightarrow \begin{cases} \Theta'' + \lambda\Theta = 0 \\ \Theta(0) = \Theta'(\frac{\pi}{2}) = 0 \end{cases}, \lambda_n = (2n-1)^2, \{\sin(2n-1)\theta\}, n = 1, 2, \dots$$

$$r^2 R'' + rR' - \lambda R = 0 \rightarrow \mu^2 = (2n-1)^2, R = c_1 r^{2n-1} + c_2 r^{1-2n}. \text{ Como } R \text{ debe ser acotada: } R_n = \{r^{2n-1}\}.$$

Probamos entonces $u(r, \theta) = \sum_{n=1}^{\infty} c_n r^{2n-1} \sin(2n-1)\theta$. Sólo nos falta aplicar el último dato de contorno:

$$u(2, \theta) = \sum_{n=1}^{\infty} c_n 2^{2n-1} \sin(2n-1)\theta = \pi. \text{ Por tanto debe ser:}$$

$$c_n 2^{2n-1} = \frac{2}{\pi/2} \int_0^{\pi/2} \pi \sin(2n-1)\theta d\theta = \left[-\frac{4}{2n-1} \cos(2n-1)\theta \right]_0^{\pi/2} = \frac{4}{2n-1} \quad (\text{esos cosenos se anulan en } \frac{\pi}{2}).$$

Los 2 primeros coeficientes son: $2c_1 = 4, c_1 = 2$ y $8c_2 = \frac{4}{3}, c_2 = \frac{1}{6}.$

Y los 2 primeros de la serie solución son: $\boxed{u(r, \theta) = 2r \sin \theta + \frac{1}{6} r^3 \sin 3\theta + \dots} = \sum_{n=1}^{\infty} \frac{r^{2n-1}}{(2n-1) 2^{2n-3}} \sin(2n-1)\theta.$



28. a) $\Delta u = 0, 1 < r < 2$
 $u(1, \theta) = 0, u(2, \theta) = 1 + \sin \theta$ Sabemos: $\begin{cases} \Theta'' + \lambda \Theta = 0 \\ \Theta(2\pi) = \Theta(0) \end{cases} \rightarrow \lambda_n = n^2, \Theta_n = \{\cos n\theta, \sin n\theta\} \rightarrow$ 

$$r^2 R'' + rR' - n^2 R = 0 \rightarrow R_0 = c_1 + c_2 \ln r \xrightarrow{R(1)=0} R_0 = \{\ln r\}$$

$$R_n = c_1 r^n + c_2 r^{-n} \rightarrow R_n = \{r^n - r^{-n}\} \rightarrow u = a_0 \ln r + \sum_{n=1}^{\infty} (r^n - r^{-n}) [a_n \cos n\theta + b_n \sin n\theta]$$

$$u(2, \theta) = a_0 \ln 2 + \sum_{n=1}^{\infty} (2^n - \frac{1}{2^n}) [a_n \cos n\theta + b_n \sin n\theta] = 1 + \sin \theta \rightarrow a_0 = \frac{1}{\ln 2}, b_1 = \frac{2}{3}, \text{ resto } 0. \quad \boxed{u = \frac{\ln r}{\ln 2} + \frac{2}{3} (r - \frac{1}{r}) \sin \theta}.$$

b) $\Delta u = \cos \theta, r < 2$
 $u(2, \theta) = \sin 2\theta$ $u = a_0(r) + \sum_{n=1}^{\infty} [a_n(r) \cos n\theta + b_n(r) \sin n\theta] \rightarrow$

$$\frac{r a_0'' + a_0'}{r} + \sum_{n=1}^{\infty} \left[\frac{r^2 a_n'' + r a_n' - n^2 a_n}{r^2} \cos n\theta + \frac{r^2 b_n'' + r b_n' - n^2 b_n}{r^2} \sin n\theta \right] = \cos \theta.$$

Del dato de contorno: $a_n(2) = 0; b_n(2) = 0, n \neq 2; b_2(2) = 1$ y todas acotadas $\Rightarrow a_{n \neq 1}, b_{n \neq 2} \equiv 0$ (hay unicidad).

Y además: $\begin{cases} r^2 a_1'' + r a_1' - a_1 = r^2 \\ \text{acotada y } a_1(2) = 0 \end{cases} \xrightarrow{a_1 p = A r^2} a_1 = c_1 r + c_2 r^{-1} + \frac{r^2}{3} \rightarrow \frac{c_2}{2c_1 + \frac{4}{3}} = 0 \rightarrow a_1(r) = \frac{r^2}{3} - \frac{2r}{3};$

$\begin{cases} r^2 b_2'' + r b_2' - 4b_2 = 0 \\ \text{acotada y } b_2(2) = 1 \end{cases} \rightarrow b_2 = c_1 r^2 + c_2 r^{-2} \rightarrow \frac{c_2}{4c_1 + 1} = 0 \rightarrow b_2(r) = \frac{r^2}{4} \rightarrow \boxed{u(r, \theta) = \frac{1}{3} r(r-2) \cos \theta + \frac{1}{4} r^2 \sin 2\theta}.$

c) $\Delta u = 0, r < 4, 0 < \theta < \pi$
 $u(4, \theta) = 2 \sin \frac{\theta}{2}, u(r, 0) = u(r, \pi) = 0$ $\begin{cases} \Theta'' + \lambda \Theta = 0 \\ \Theta(0) = \Theta(\pi) = 0 \end{cases}, \lambda_n = \frac{(2n-1)^2}{4}, \Theta_n = \{\sin \frac{(2n-1)\theta}{2}\}, n = 1, \dots$ 

$$r^2 R'' + rR' - \lambda R = 0 \xrightarrow{\lambda = \frac{(2n-1)^2}{4}} \mu = \pm \frac{(2n-1)}{2}, R = c_1 r^{(2n-1)/2} + c_2 r^{-(2n-1)/2} \xrightarrow{R \text{ acotado en } r=0} R_n = \{r^{(2n-1)/2}\}, n = 1, 2, \dots$$

Probamos, pues, $u(r, \theta) = \sum_{n=1}^{\infty} c_n r^{(2n-1)/2} \sin \frac{(2n-1)\theta}{2}$, y para calcular los c_n imponemos el dato que falta:

$$u(4, \theta) = \sum_{n=1}^{\infty} 4^{(2n-1)/2} c_n \sin \frac{(2n-1)\theta}{2} = 2 \sin \frac{\theta}{2} \rightarrow 4^{1/2} c_1 = 2, c_1 = 1 \text{ y el resto } 0. \quad \boxed{u(r, \theta) = r^{1/2} \sin \frac{\theta}{2}}.$$

29. a) $r^2 R'' + rR' - \frac{1}{4} R = 7r^3$ Euler. $\mu(\mu-1) + \mu - \frac{1}{4} = 0 \rightarrow \mu = \pm \frac{1}{2}, R = c_1 r^{1/2} + c_2 r^{-1/2} + R_p.$

Con la fvc: $\begin{vmatrix} r^{1/2} & r^{-1/2} \\ r^{-1/2/2} & -r^{-3/2/2} \end{vmatrix} = -\frac{1}{r}, R_p = -r^{-1/2} \int \frac{r^{1/2} \cdot 7r}{1/r} + r^{1/2} \int \frac{r^{-1/2} \cdot 7r}{1/r} = -2r^3 + \frac{14}{5} r^3. \quad \boxed{R = c_1 r^{1/2} + c_2 r^{-1/2} + \frac{4}{5} r^3}.$

O mejor, $R_p = A r^3$ ($R_p = A e^{3s}$ en la de coeficientes constantes) $\rightarrow 6A + 3A - \frac{A}{4} = 7, A = \frac{4}{5} \nearrow$

b) $u_{rr} + \frac{1}{r} u_r + \frac{1}{r^2} u_{\theta\theta} = 7r \cos \frac{\theta}{2}, r < 1, 0 < \theta < \pi$
 $u(1, \theta) = u_{\theta}(r, 0) = u(r, \pi) = 0$ $u = R\Theta \rightarrow \begin{cases} \Theta'' + \lambda \Theta = 0 \\ \Theta'(0) = \Theta(\pi) = 0 \end{cases} \rightarrow$ 

$$\lambda_n = \frac{(2n-1)^2}{4}, \Theta_n = \{\cos \frac{(2n-1)\theta}{2}\}, n = 1, 2, \dots$$

Llevando $u(r, \theta) = \sum_{n=1}^{\infty} R_n(r) \cos \frac{(2n-1)\theta}{2}$, a la EDP queda: $\sum_{n=1}^{\infty} \left[R_n'' + \frac{1}{r} R_n' - \frac{(2n-1)^2}{4r^2} R_n \right] \cos \frac{(2n-1)\theta}{2} = 7r \cos \frac{\theta}{2}.$ (ya desarrollada)

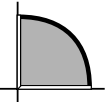
Del dato otro de contorno $u(1, \theta) = \sum_{n=1}^{\infty} R_n(1) \cos \frac{(2n-1)\theta}{2} = 0$ deducimos que $R_n(1) = 0 \forall n$.

Y además las soluciones han de estar acotadas en $r=0$. Sólo será no nula la R_1 , solución de:

$\begin{cases} r^2 R_1'' + r R_1' - \frac{1}{4} R_1 = 7r^3 \\ R_1 \text{ acotada, } R_1(1) = 0 \end{cases} \xrightarrow{a], ac.} R_1 = c_1 r^{1/2} + \frac{4}{5} r^3 \xrightarrow{R_1(1)=0} c_1 = -\frac{4}{5}. \text{ La solución es: } \boxed{u(r, \theta) = \frac{4}{5} [r^3 - r^{1/2}] \cos \frac{\theta}{2}}.$

30. $\begin{cases} \Delta u = 0, r < 2, \theta \in (0, \frac{\pi}{2}) \\ u_r(2, \theta) + k u(2, \theta) = 8 \cos 2\theta \\ u_{\theta}(r, 0) = u_{\theta}(r, \frac{\pi}{2}) = 0 \end{cases}$ $u = R\Theta \rightarrow \begin{cases} \Theta'' + \lambda \Theta = 0 \\ \Theta'(0) = \Theta'(\frac{\pi}{2}) = 0 \end{cases}, \lambda_n = 4n^2, \Theta_n = \{\cos 2n\theta\}, n = 0, 1, 2, \dots$

Y además: $r^2 R'' + rR' - \lambda R = 0, R \text{ acotado} \rightarrow R_n = \{r^{2n}\}.$

Imponiendo a $u = \frac{a_0}{2} + \sum_{n=1}^{\infty} a_{2n} r^{2n} \cos 2n\theta$ el último dato: $k \frac{a_0}{2} + \sum_{n=1}^{\infty} a_{2n} [2n 2^{2n-1} + k 2^{2n}] \cos 2n\theta = 8 \cos 2\theta.$ 

i) Si $k=1$, todos los $a_{2n}=0$, excepto $a_2[4+4]=8$, y la solución (única) es: $\boxed{u = r^2 \cos 2\theta}.$

ii) Si $k=0$ (Neumann), a_0 queda libre, $a_2[4+0]=8$, y demás $a_{2n}=0$. Infinitas soluciones: $\boxed{u = C + 2r^2 \cos 2\theta}.$